

Efficient interfacing of single photons and single molecules

Vahid Sandoghdar

*Max Planck Institute for the Science of Light
&
Friedrich-Alexander University
Erlangen, Germany*

www.mpl.mpg.de

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MAX PLANCK INSTITUTE
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Alexander von Humboldt
Stiftung/Foundation



Outline

Introductory discussion of light-matter interaction

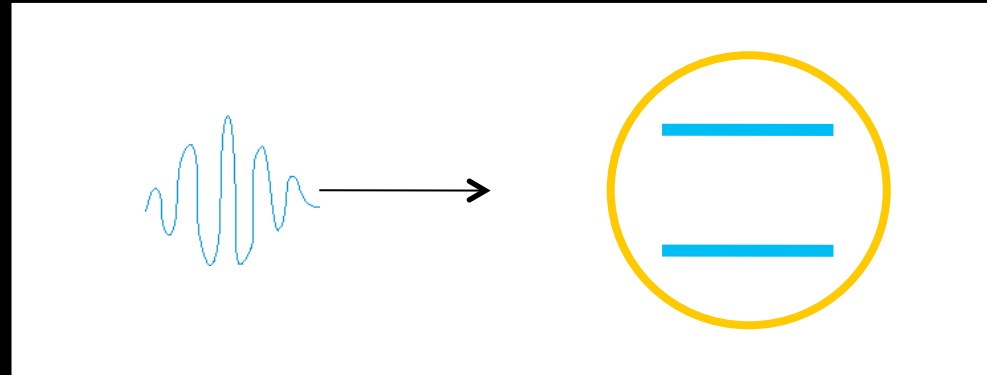
High resolution single molecule spectroscopy

Strong interaction of light and a single molecule

Generation of tunable narrow-band single photons

Tailoring light emission with optical antennas

A single photon meets a single atom



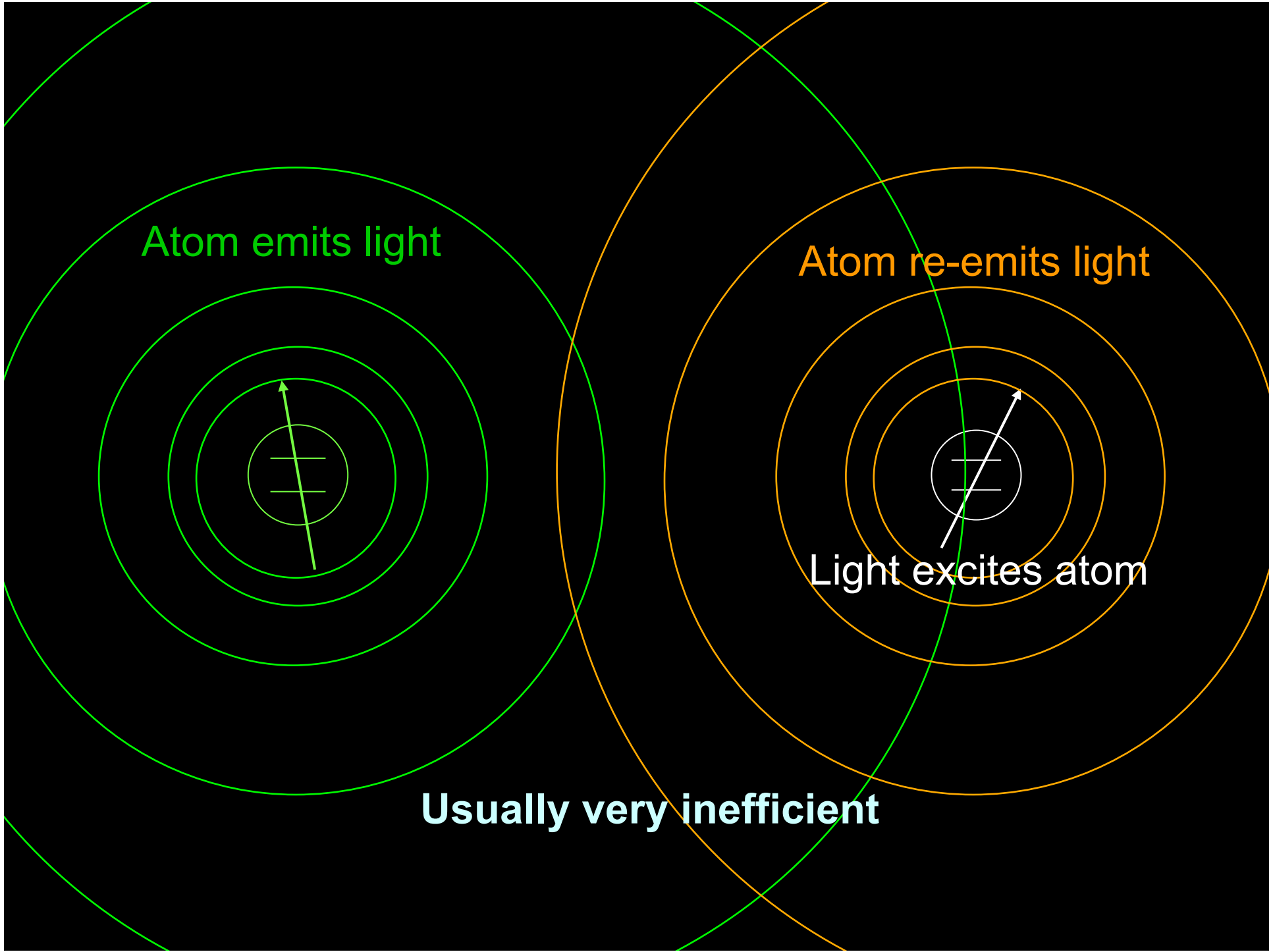
The probability that a photon meets the emitter?

Atom emits light

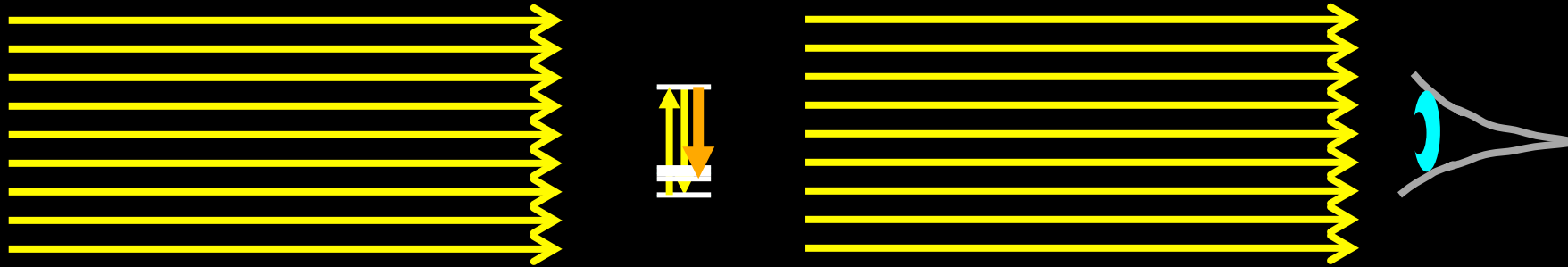
Atom re-emits light

Light excites atom

Usually very inefficient



Simple absorption measurement



$$I_{int} = I_0 (\sigma_{abs}/A)$$

Ideal case: $\sigma \sim (500 \text{ nm})^2$; typical $A \sim (10-100 \text{ } \mu\text{m})^2$

Experimental challenge:

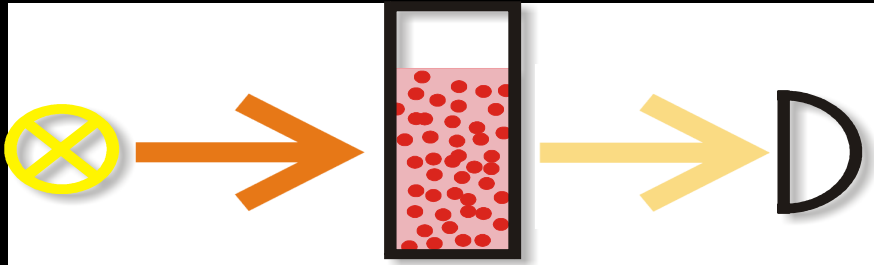
See the effect of one atom on the intensity of a laser beam

As a result, single molecule detection is commonly done in fluorescence

Improving the efficiency

Easiest solution:

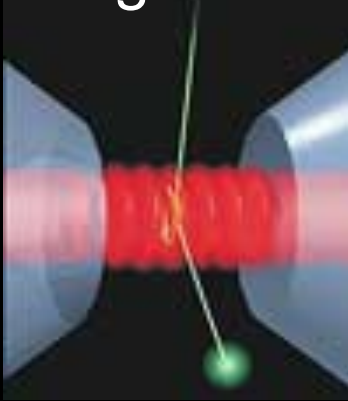
Strengthen the signal by using many atoms



$$I = I_0 \cdot \exp(-\sigma_{\text{abs}} \cdot N \cdot l)$$

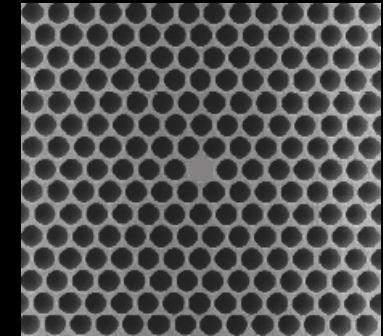
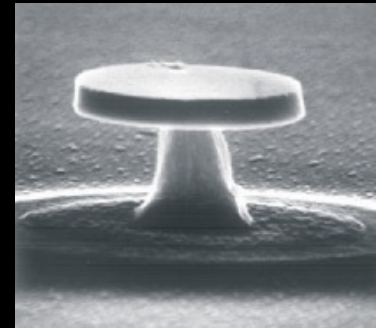
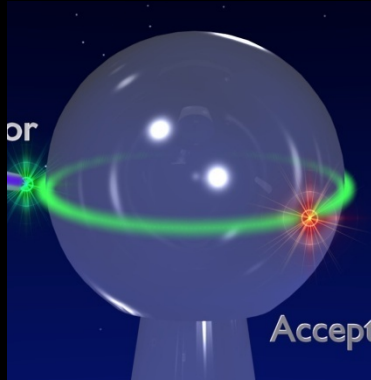
Another trick:

Lengthen the interaction time in a cavity



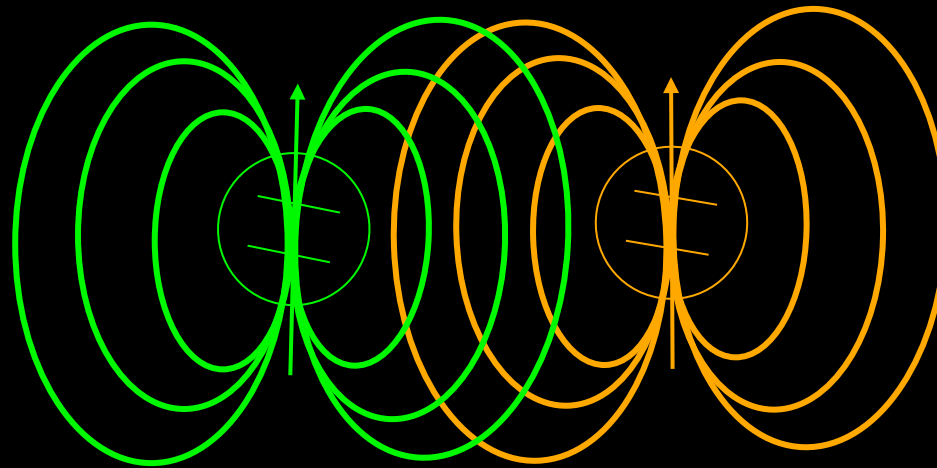
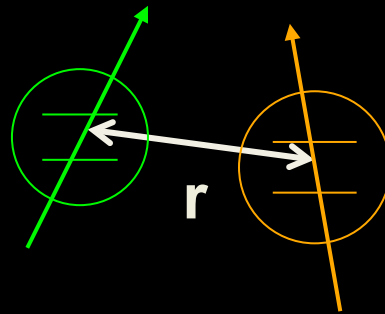
Make the Photon Stronger! Confinement in a Microcavity

$$E_{ph} \sim \sqrt{\frac{\hbar\omega}{\epsilon V}}$$



An atom also gets manipulated:
Change of the spontaneous emission rate
Change of the emission pattern

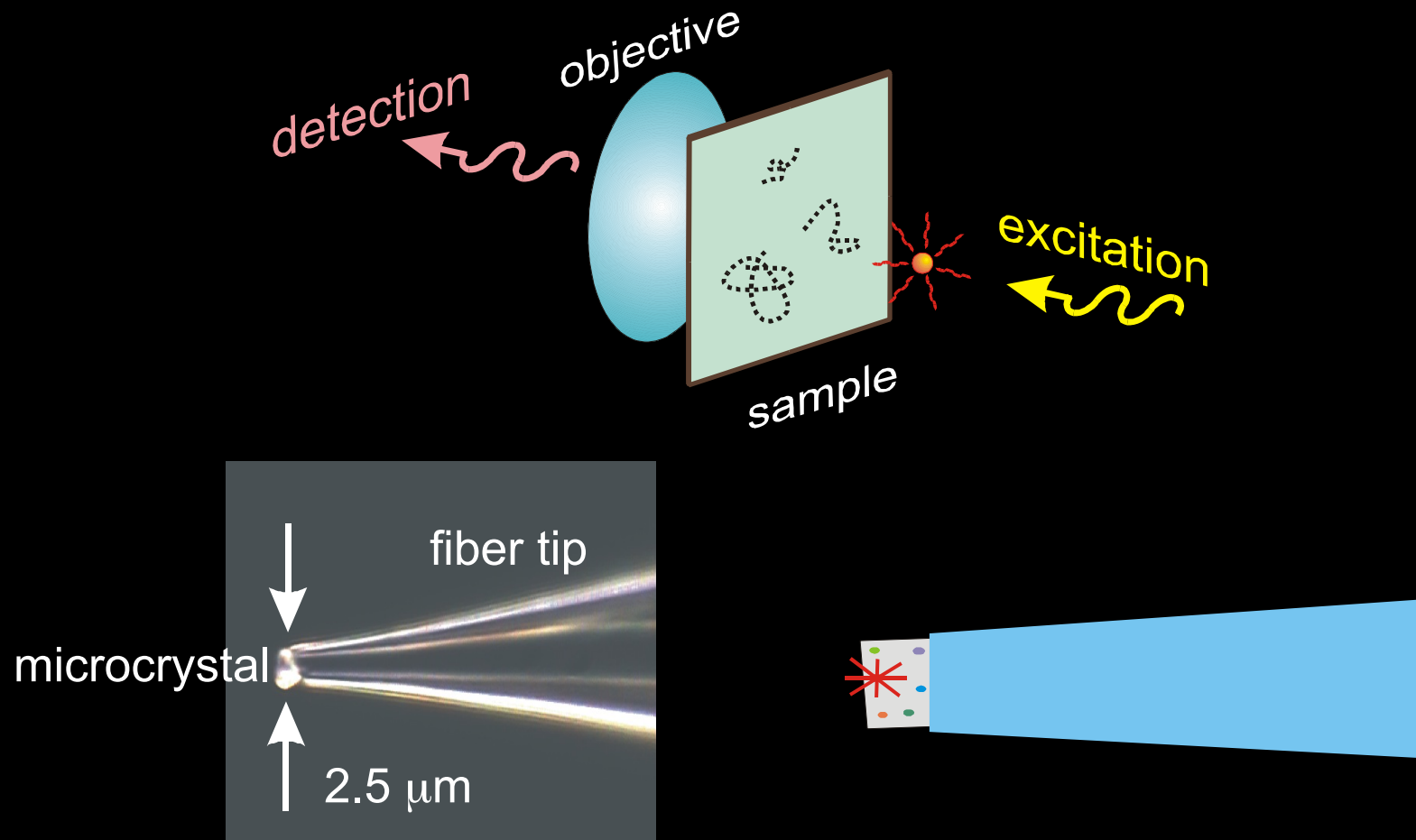
near-field dipole-dipole interaction



Variations:

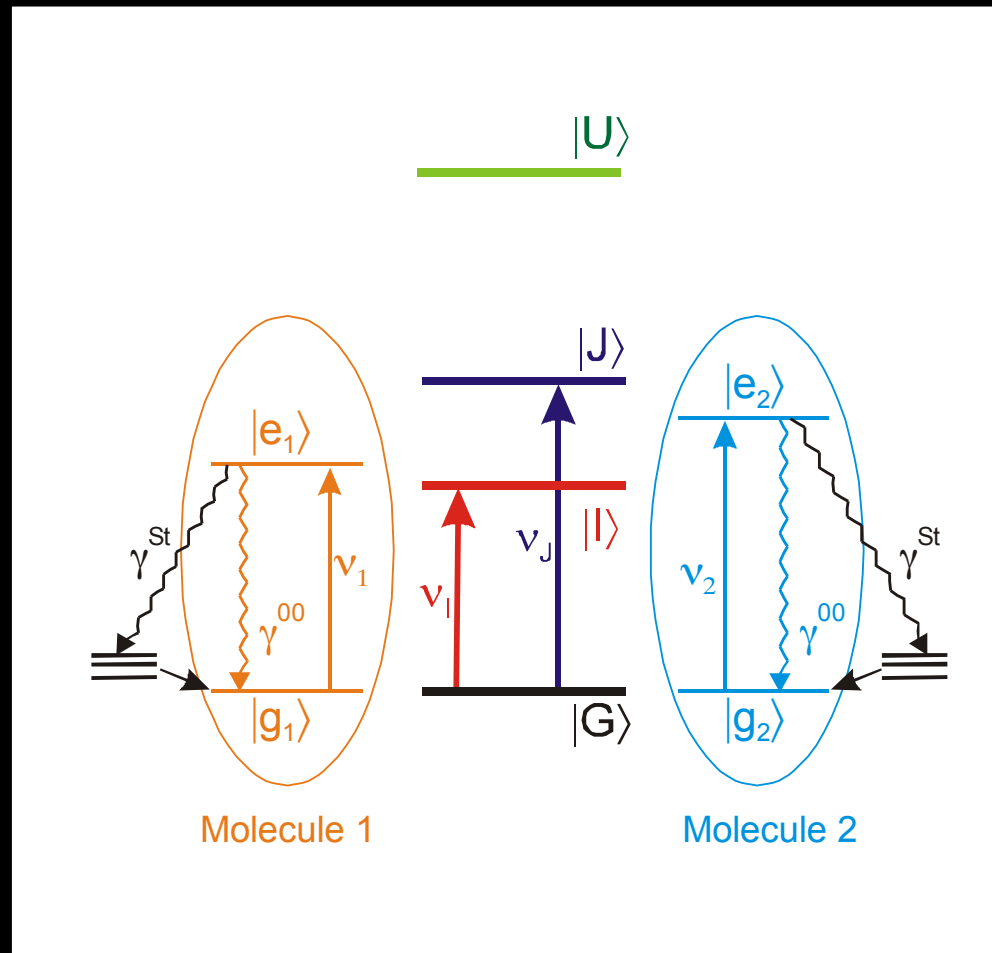
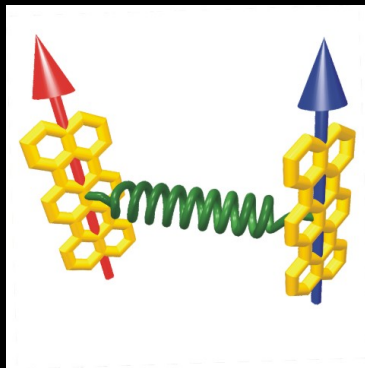
- dissipative \rightarrow Fluorescence (Förster) resonant energy transfer (FRET)
- coherent \rightarrow Dicke sub and superradiance

A positionable single-molecule probe: realization of a point-like light source



J. Michaelis, C. Hettich, J. Mlynek & V. Sandoghdar, *Nature* **405**, 325 (2000).

Coherent dipole-dipole coupling between two individual molecules

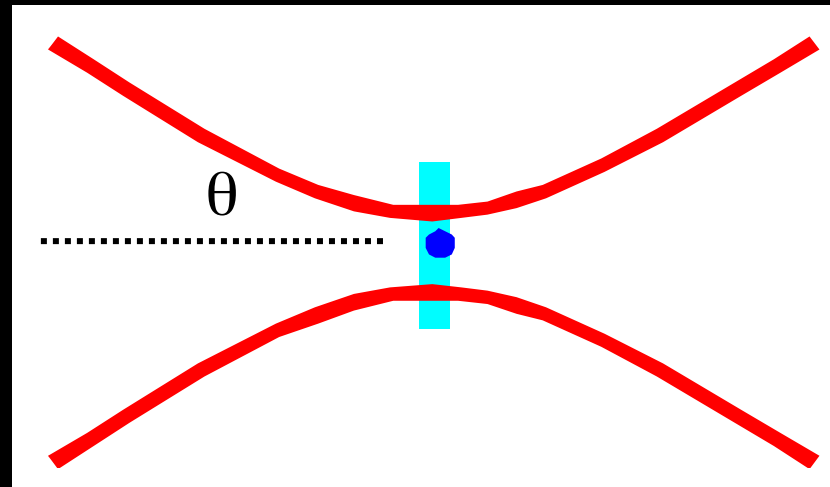


C. Hettich, C. Schmitt, J. Zitzmann, S. Kühn, I. Gerhardt & V. Sandoghdar, *Science* **298**, 385-389 (2002).

What if we simply confine a propagating wave?

Reduce A

$$I = I_0 (1 - \sigma_{\text{abs}}/A)$$



FWHM of a focused beam $\sim \lambda/2NA$; $NA = n \sin \theta$

$A \sim \lambda^2/4$ for $NA \sim 1$

$$\sigma_{\text{abs}} = 3\lambda^2/2\pi \sim \lambda^2/2$$

Scattered power = Incident intensity . cross section

Probability of being in the excited state Spontaneous emission rate (natural linewidth)

Scattered power = $\rho_{22} \cdot \gamma_0 \cdot h\nu$ — Energy per emitted photon

$$\rho_{22} = \frac{\frac{1}{4} \left(\frac{\gamma}{\gamma_0} \right) |\mathcal{V}|^2}{(\omega_0 - \omega)^2 + \frac{1}{2} \left(\frac{\gamma}{\gamma_0} \right) |\mathcal{V}|^2 + \gamma^2}$$

$$\gamma = \gamma_0 + \gamma_{coll}$$

On resonance and in the weak excitation limit,

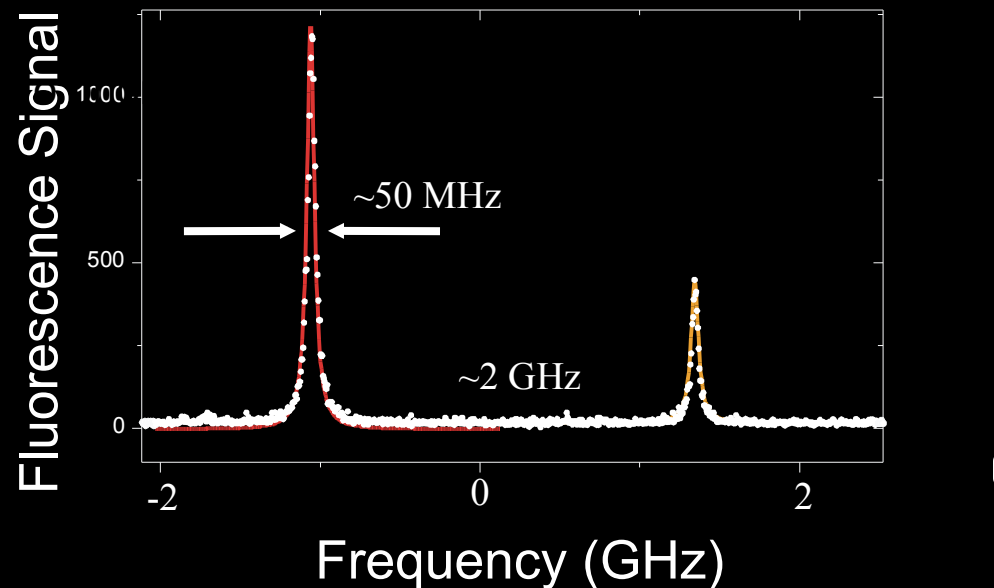
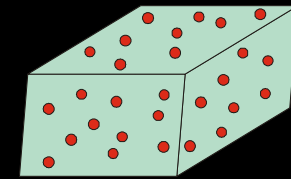
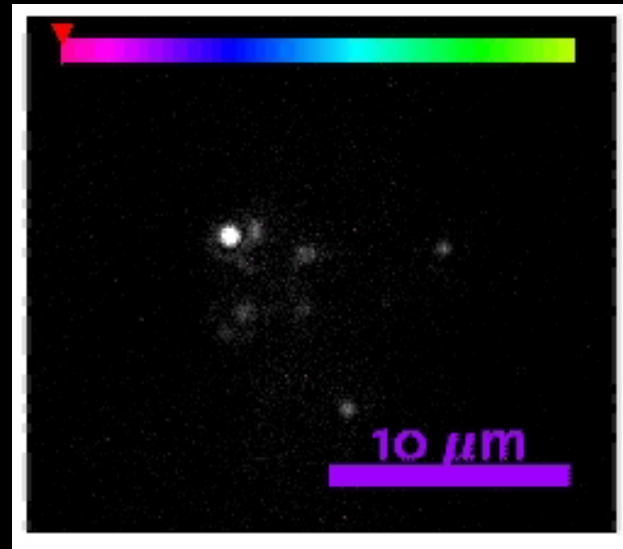
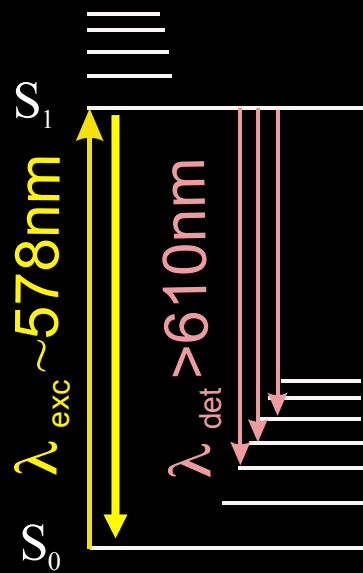
$$\text{Scattered power} \propto \gamma_0 \cdot h\nu \cdot \frac{|\mathcal{V}|^2}{\gamma_0 \gamma} = \left(\frac{\gamma_0}{\gamma} \right) h\nu \frac{|\mathcal{V}|^2}{\gamma_0}$$

$$|\mathcal{V}|^2 \propto |E|^2 |D_{eg}|^2$$

$$\gamma_0 \propto |D_{eg}|^2 \nu^3$$

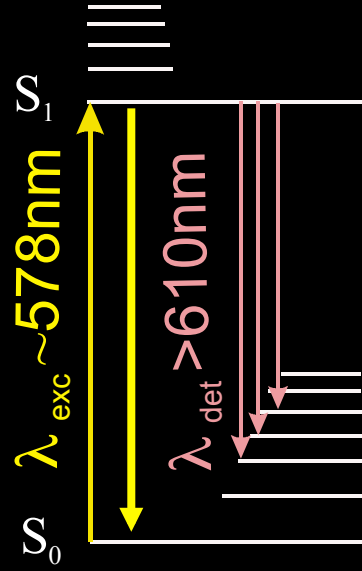
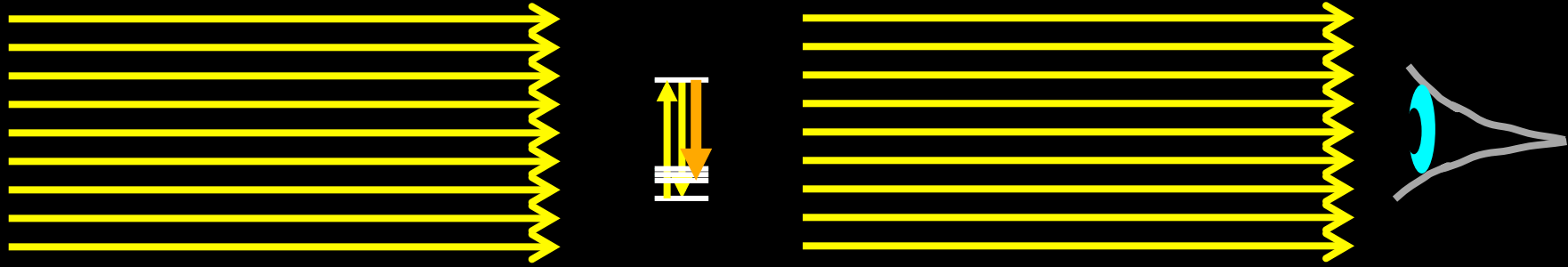
$$\sigma_{abs} = \frac{3\lambda^2}{2\pi} \frac{\gamma_0}{\gamma} \propto \begin{matrix} \rightarrow \text{natural linewidth} \\ \rightarrow \text{homogeneous linewidth} \end{matrix}$$

Single molecule spectroscopy in solid matrices

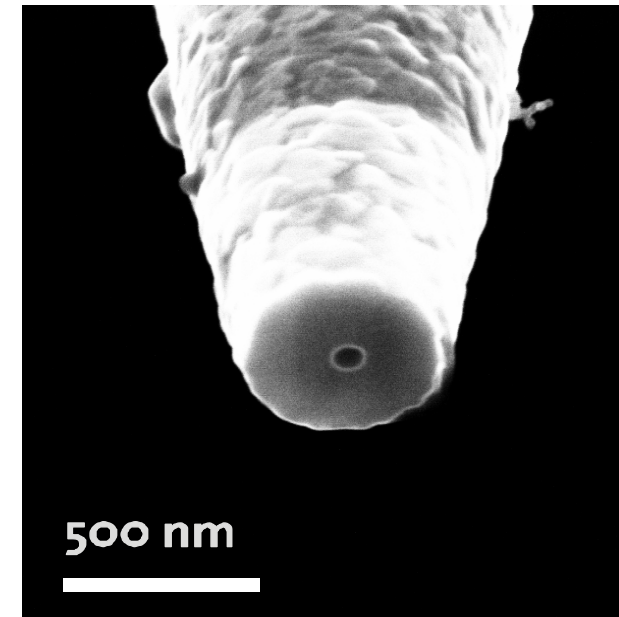
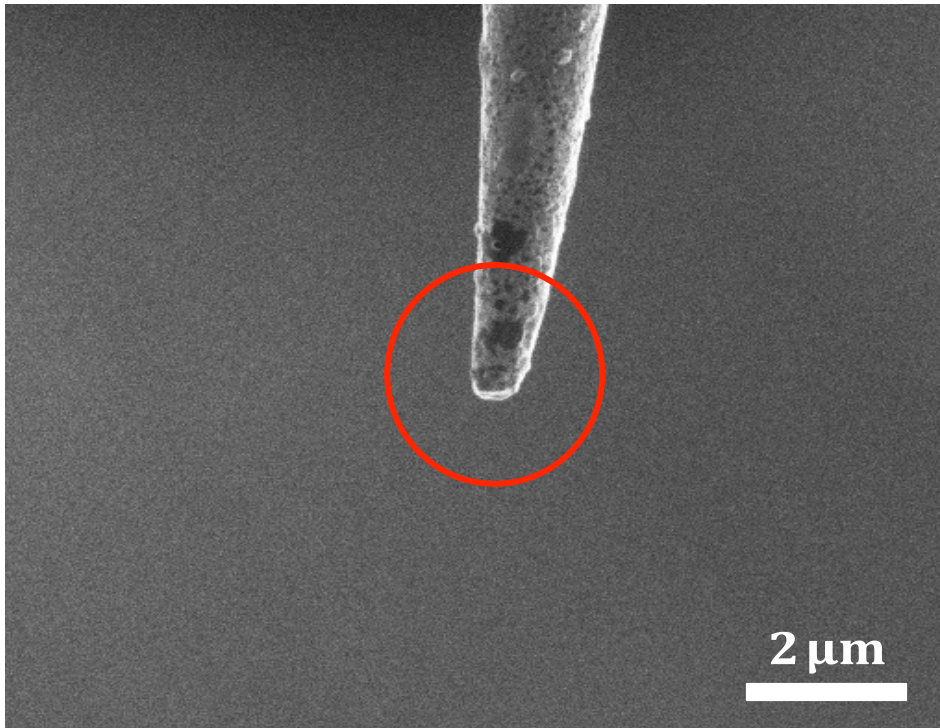


M. Orrit and J. Bernard, *Phys. Rev. Lett.* **65**, 2716 (1990).

Coherent detection of single molecules: Extinction spectroscopy

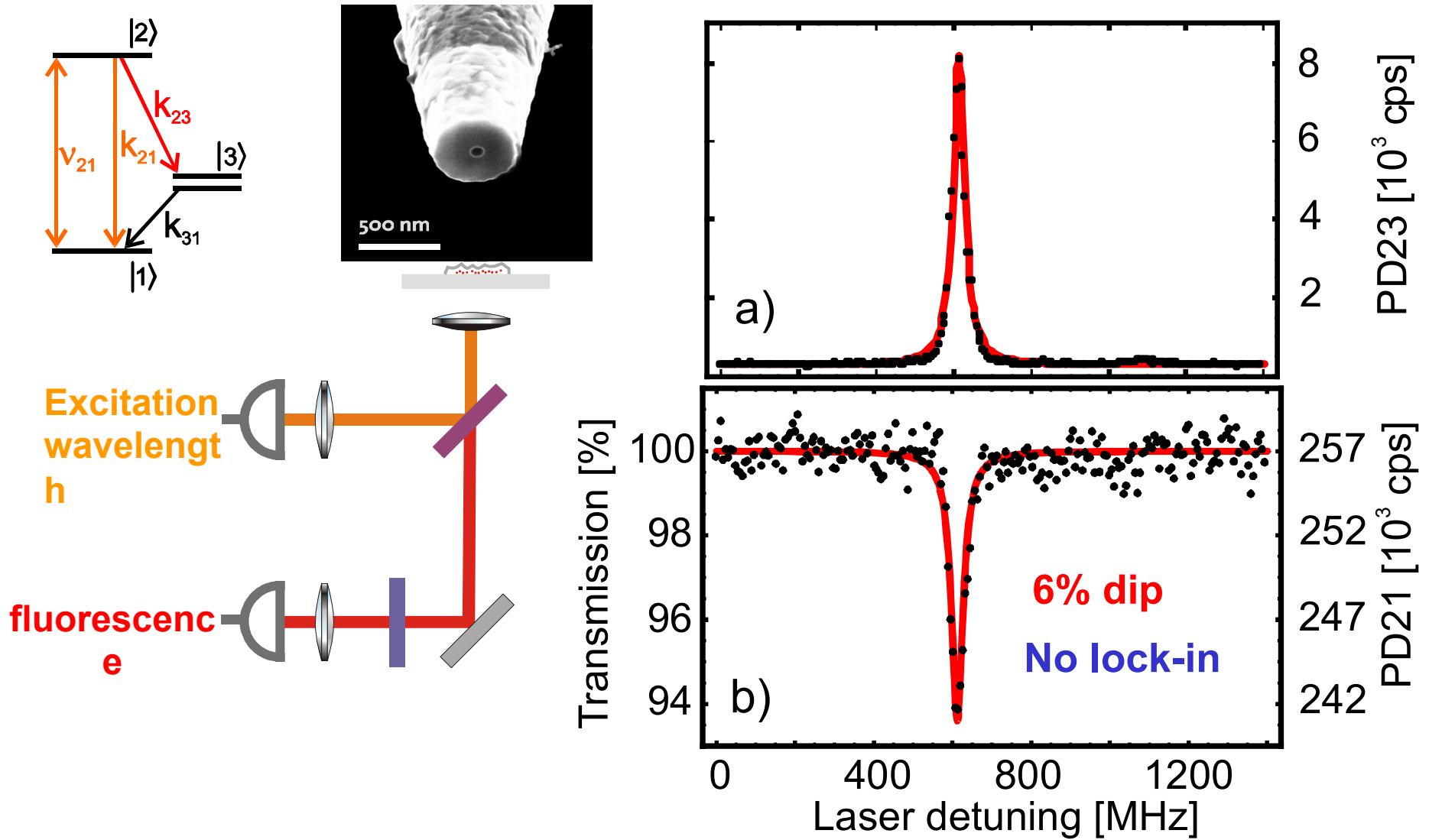


Confine light by a subwavelength aperture



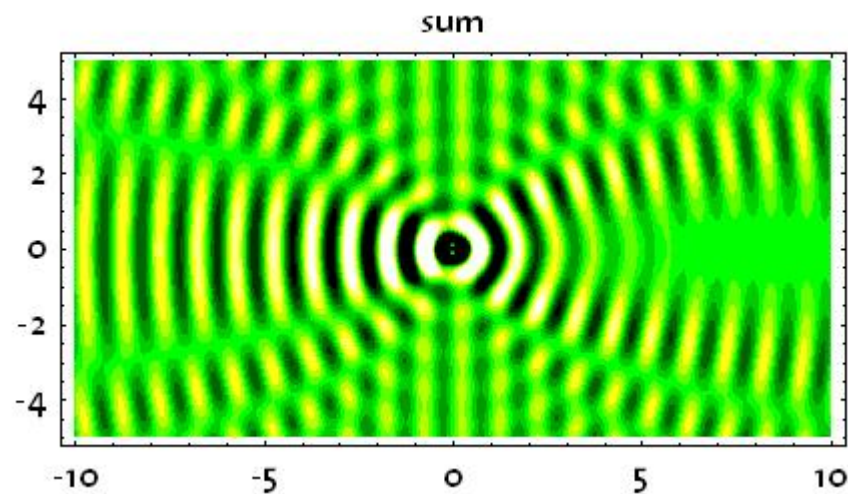
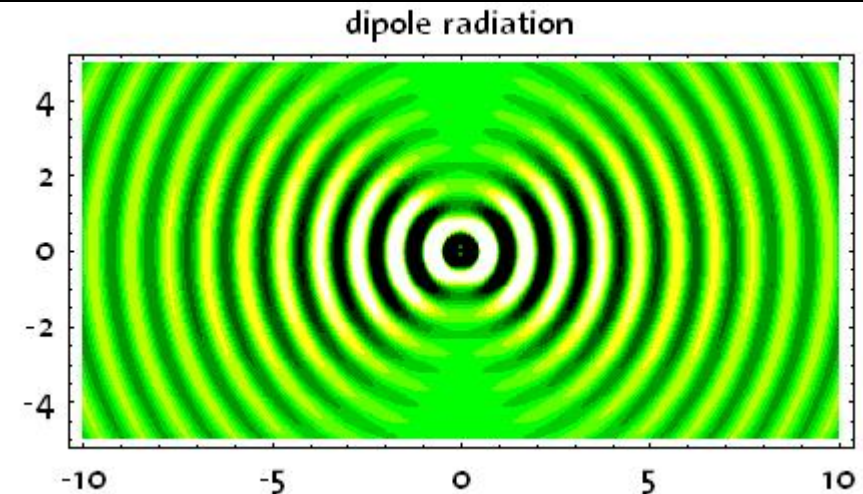
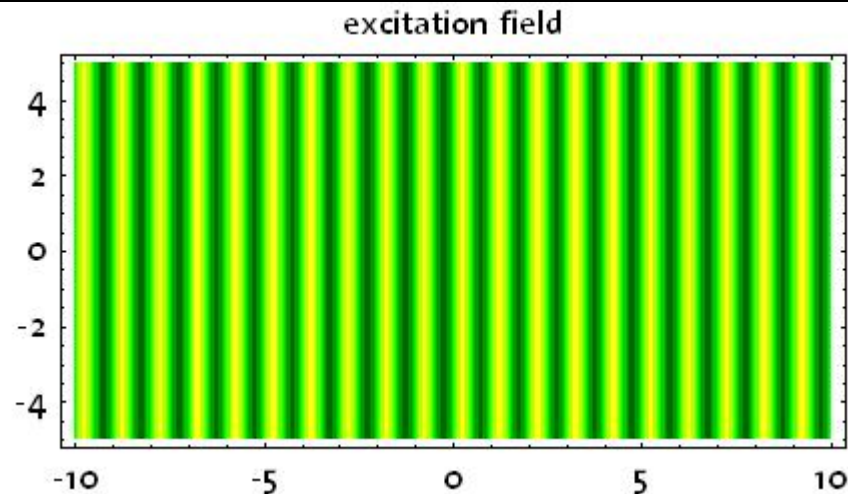
Focused ion beam (FIB) milled

Physical confinement of light to a subwavelength area



I. Gerhardt, G. Wrigge, P. Bushev, G. Zumofen, M. Agio, R. Pfab, V. Sandoghdar, *Phys. Rev. Lett.* **98**, 033601 (2007).

The extinction signal is due to interference



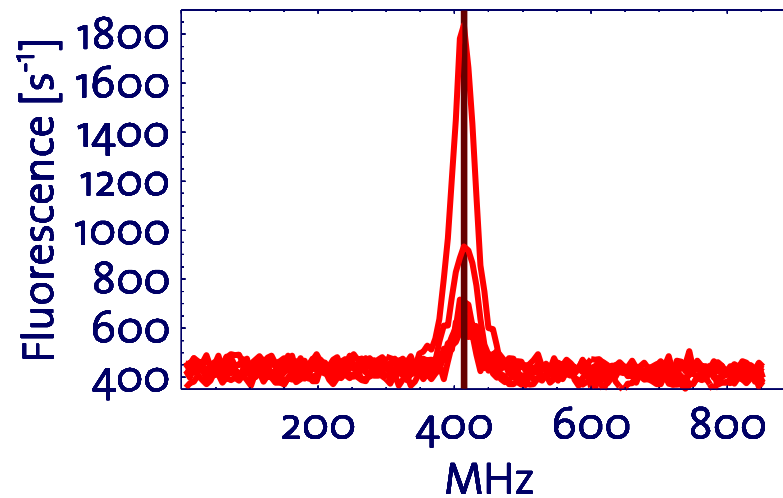
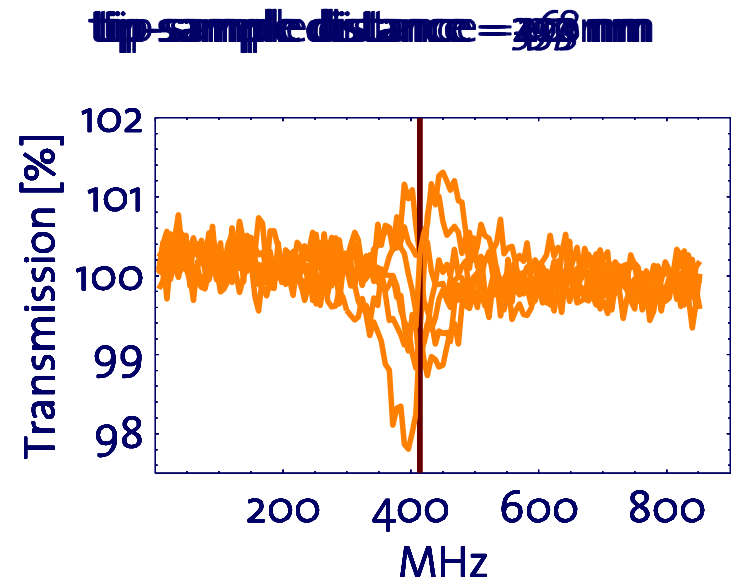
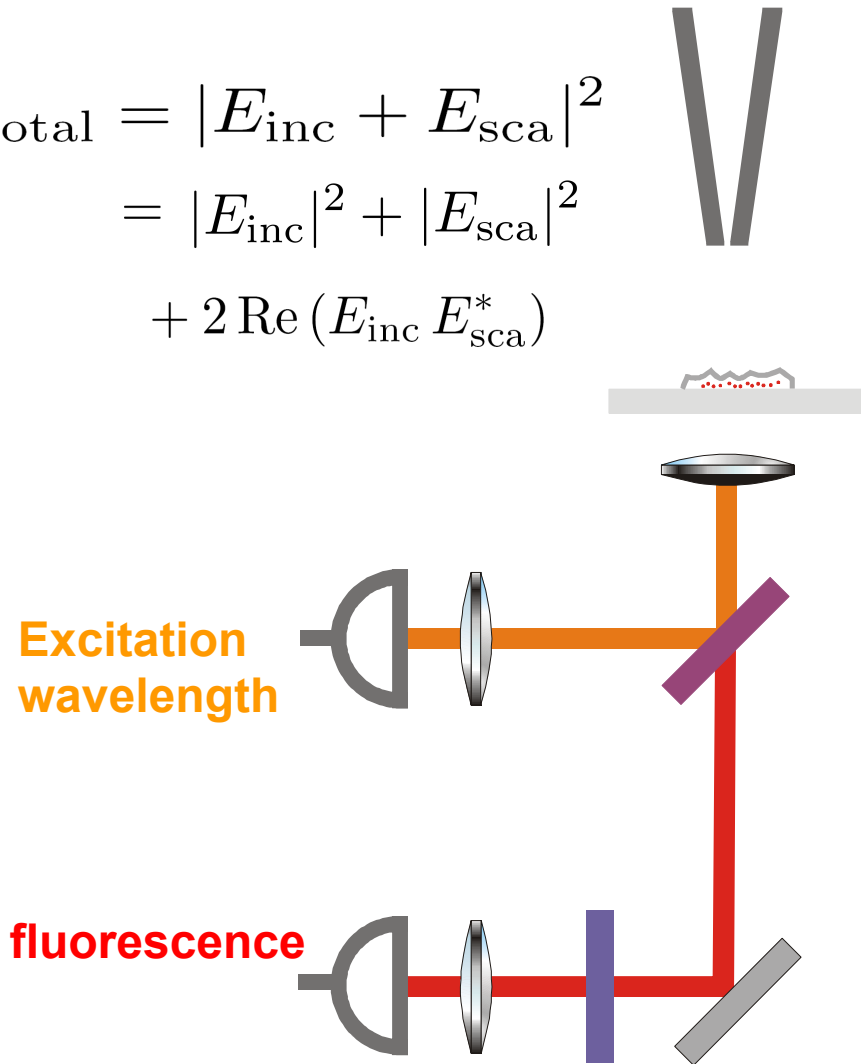
$$I_{total} = |E_{inc}|^2 + |E_{sca}|^2 + 2\text{Re}(E_{inc}E_{sca}^*)$$

↓
extinction

Complete extinction in the forward direction, hence the definition of the absorption cross section (optical theorem)

The signal originates from interference

$$\begin{aligned} I_{\text{total}} &= |E_{\text{inc}} + E_{\text{sca}}|^2 \\ &= |E_{\text{inc}}|^2 + |E_{\text{sca}}|^2 \\ &\quad + 2 \operatorname{Re}(E_{\text{inc}} E_{\text{sca}}^*) \end{aligned}$$



I. Gerhardt, G. Wrigge, M. Agio, P. Bushev, G. Zumofen, V. Sandoghdar, *Opt. Lett.* **32**, 1420 (2007).

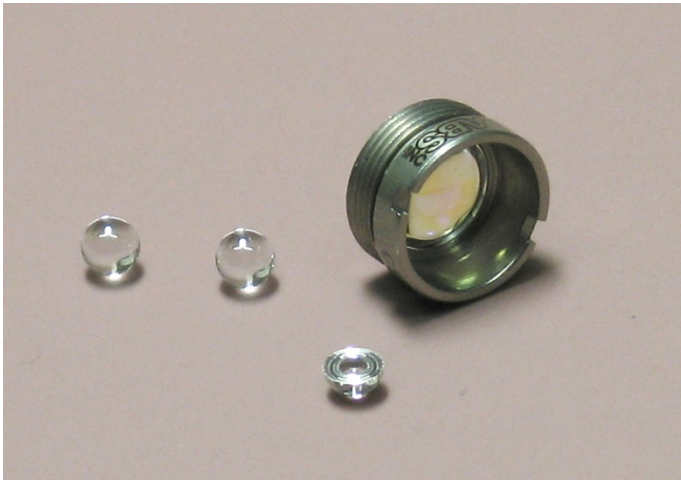
Disadvantages of near-field coupling

Very small through-put of about 10^{-4} - 10^{-5}

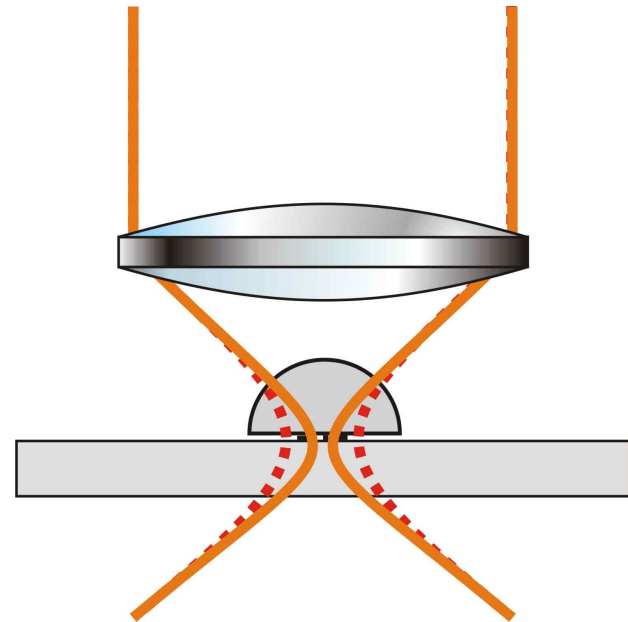
Fragile tips

Needs very thin samples

Solid immersion lenses

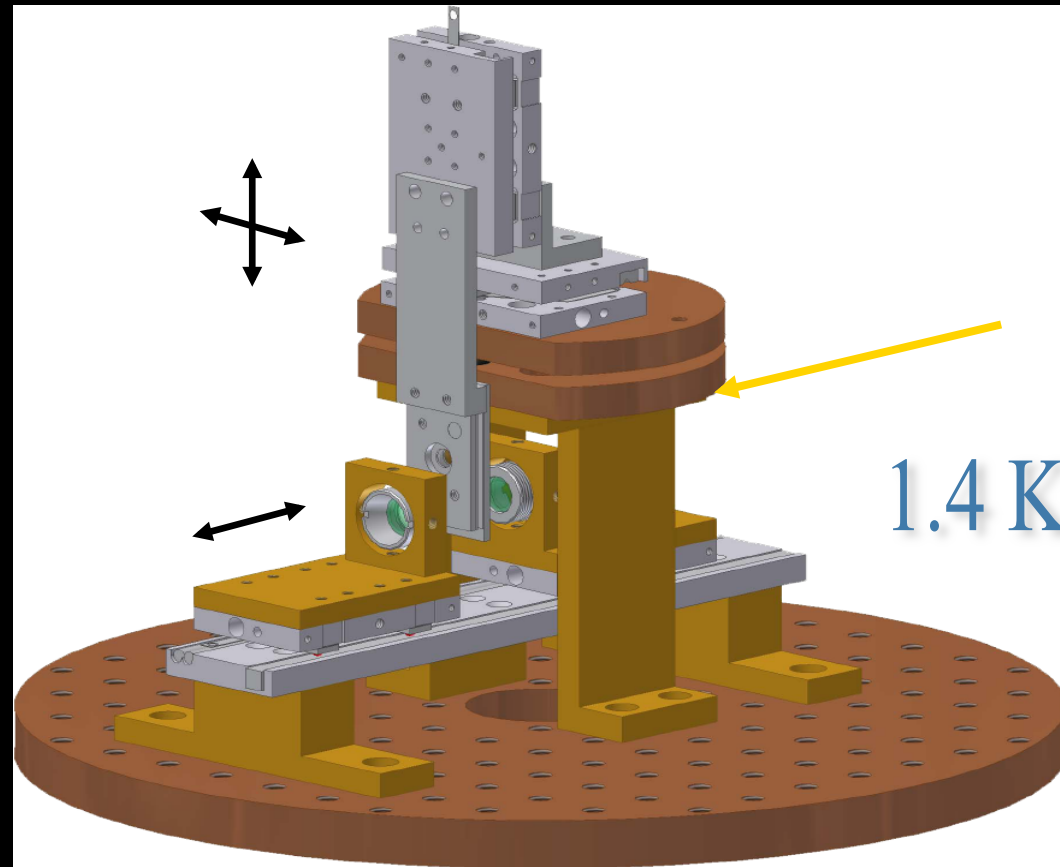


$$\lambda = \frac{\lambda_{\text{vac}}}{n}$$

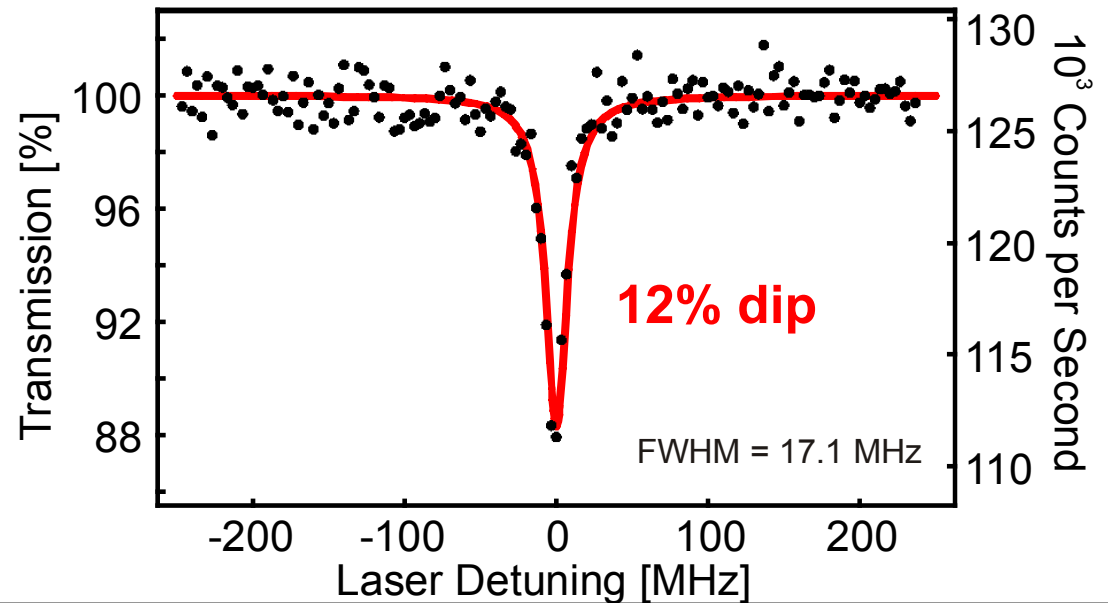
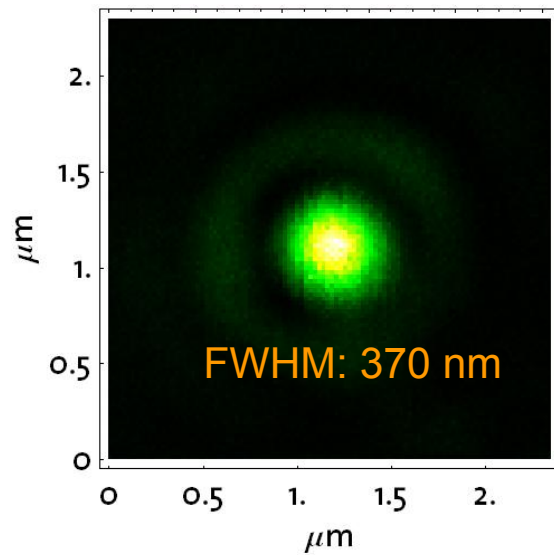
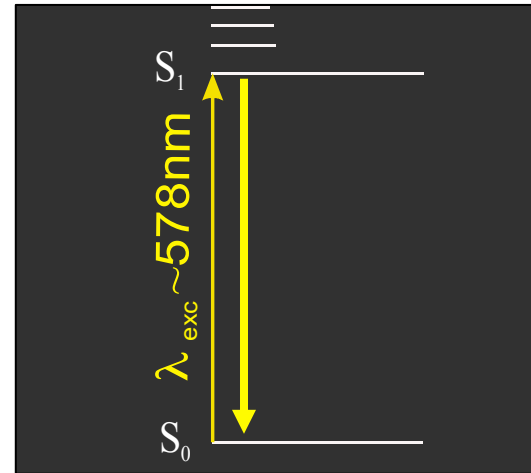
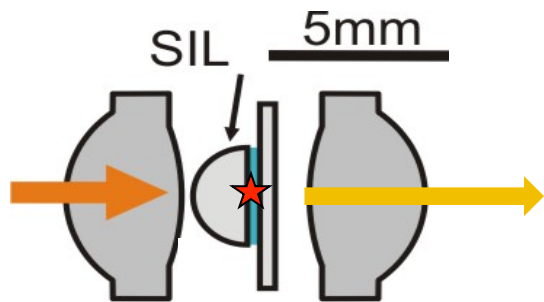


Mansfield & Kino, *Appl. Phys. Lett.* 57 (1990).

Inside the cryostat



Far-field extinction spectroscopy on a single molecule



G. Wrigge, I. Gerhardt, J. Hwang, G. Zumofen, V. Sandoghdar,
Nature Phys. 4, 60 (2008).

Absorption spectroscopy of a single molecule (1989)

VOLUME 62, NUMBER 21

PHYSICAL REVIEW LETTERS

22 MAY 1989

Optical Detection and Spectroscopy of Single Molecules in a Solid

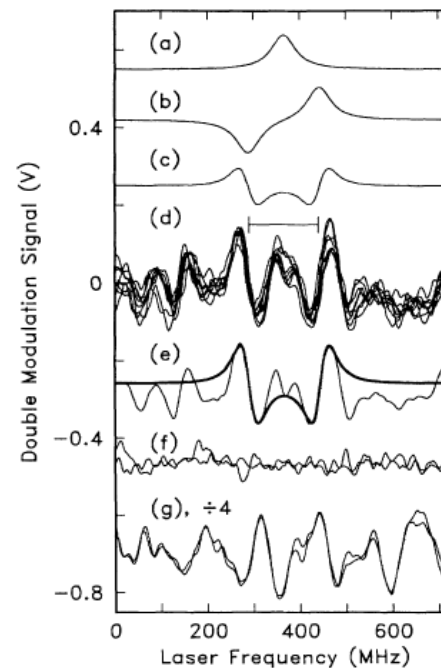


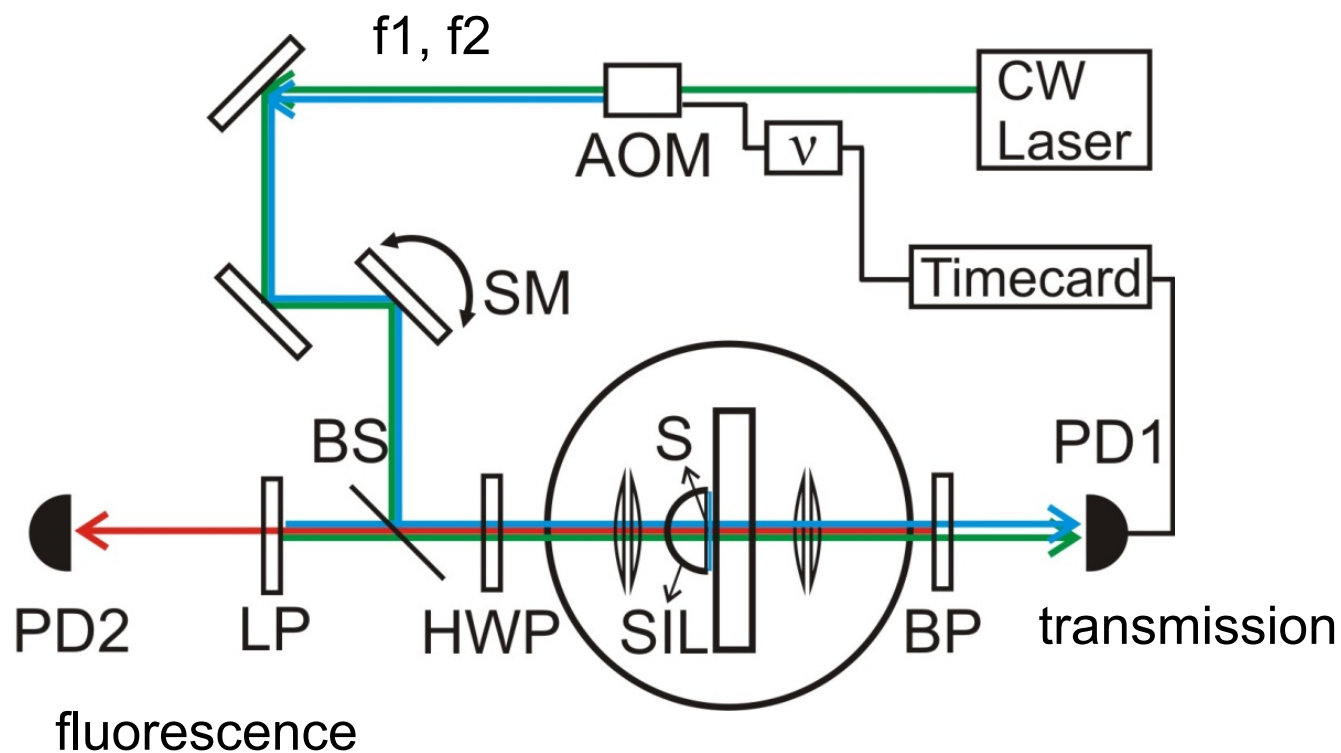
FIG. 1. Illustration of single-molecule spectra using FMS technique. (a) Simulation of absorption line, $\gamma=65$ MHz. (b) Simulation of FM spectrum for (a), $\nu_m=75$ MHz. (c) Simulation of FMS line shape. (d) SMD spectra at 592.423 nm, 512 averages, 8 traces overlaid, bar shows value of $2\nu_m=150$ MHz. (e) Average of traces in (d) (S_2 removed) with fit to the in-focus molecule (smooth curve). (f) Signal far off line at 597.514 nm. (g) Traces of SFS at the O_2 line center, 592.186 nm.

Double FM detection
 $\lambda \sim 3\mu\text{m}$

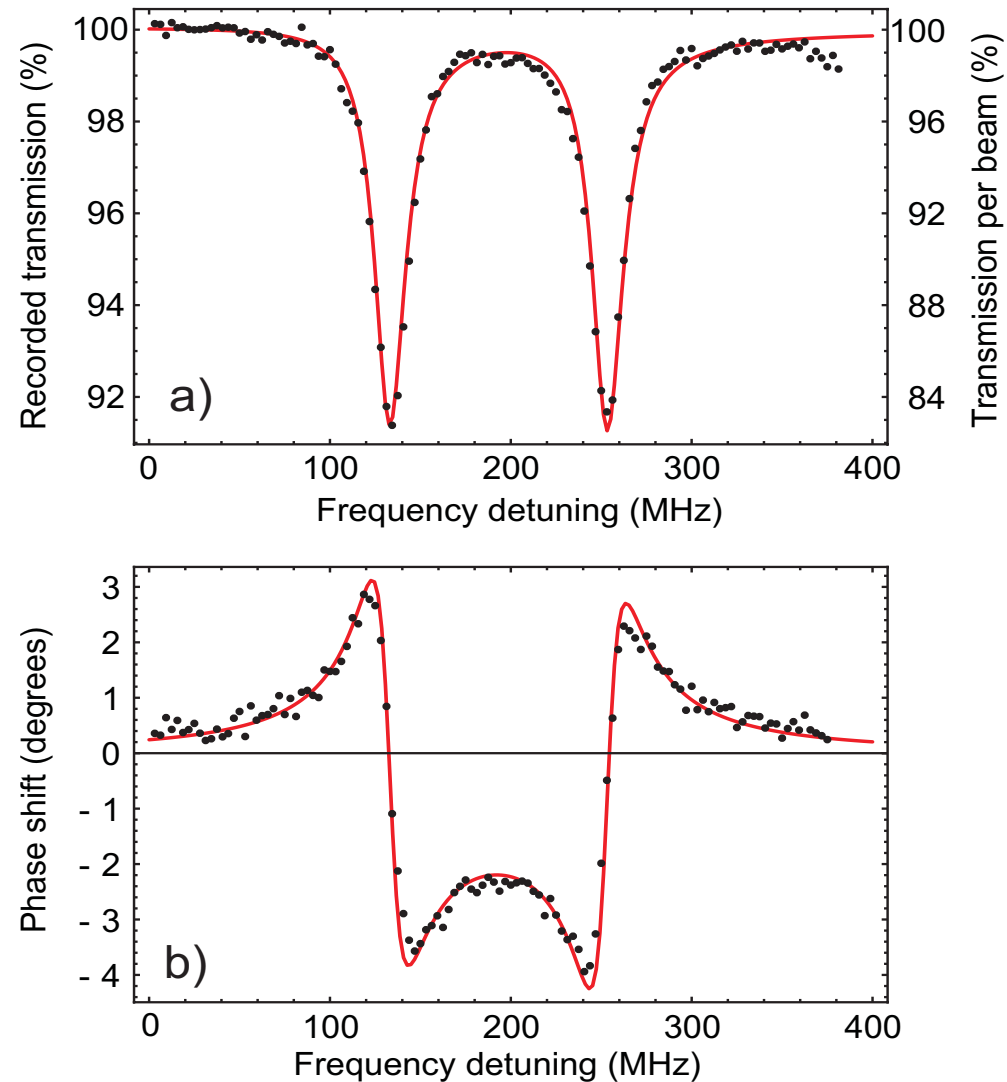
Phase shifting a laser beam by a single molecule

Need an interferometer to measure

A robust heterodyne interferometer

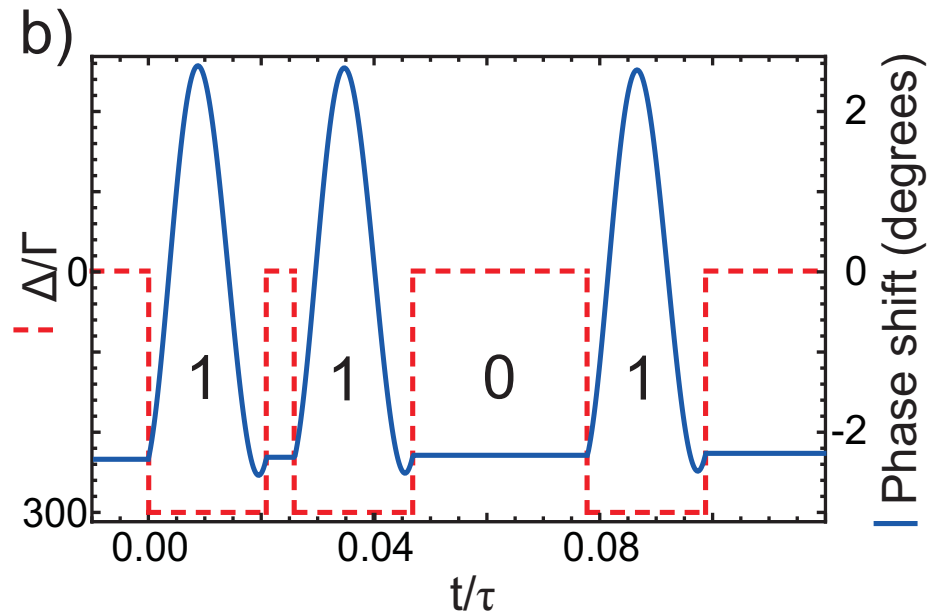
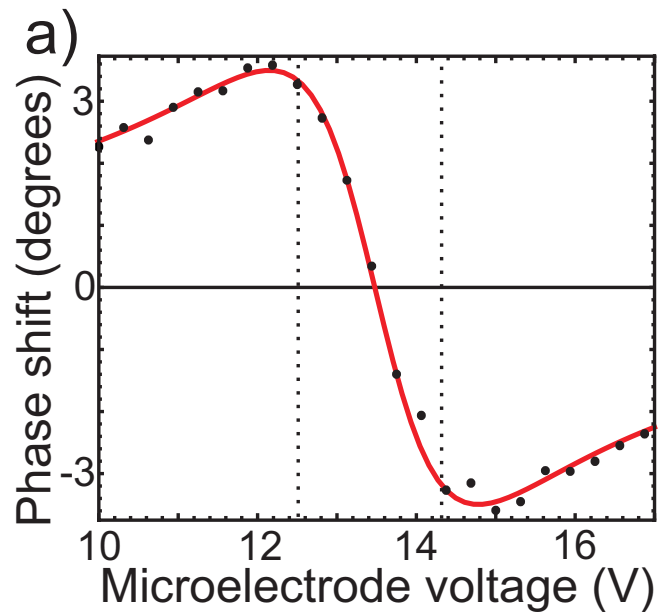
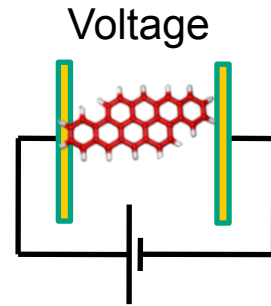


Phase shifting a laser beam by a single molecule



M. Pototschnig, Y. Chassagneux, J. Hwang, G. Zumofen, A. Renn, V. Sandoghdar, *Phys. Rev. Lett.* **107**, 063001 (2011).

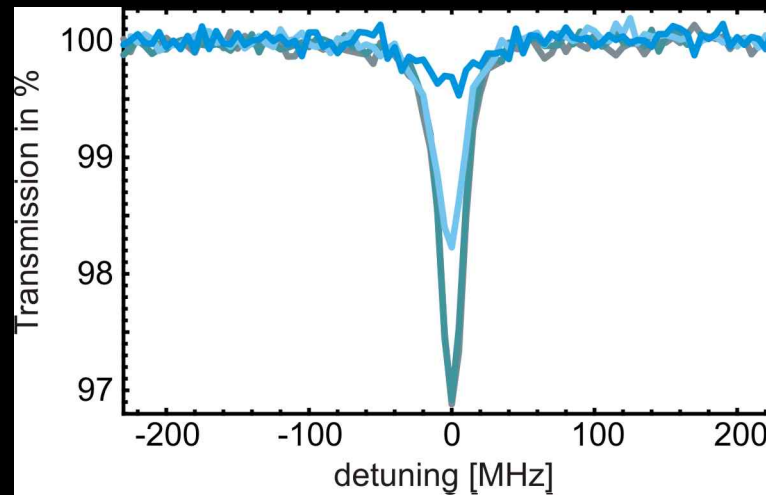
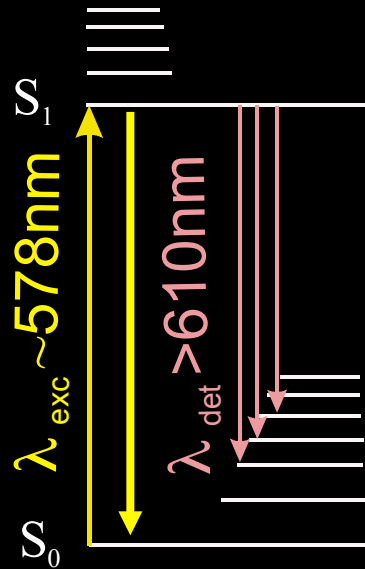
Electric-field controlled fast phase switch



Exploring nonlinearity

Intensity dependence of the extinction dip

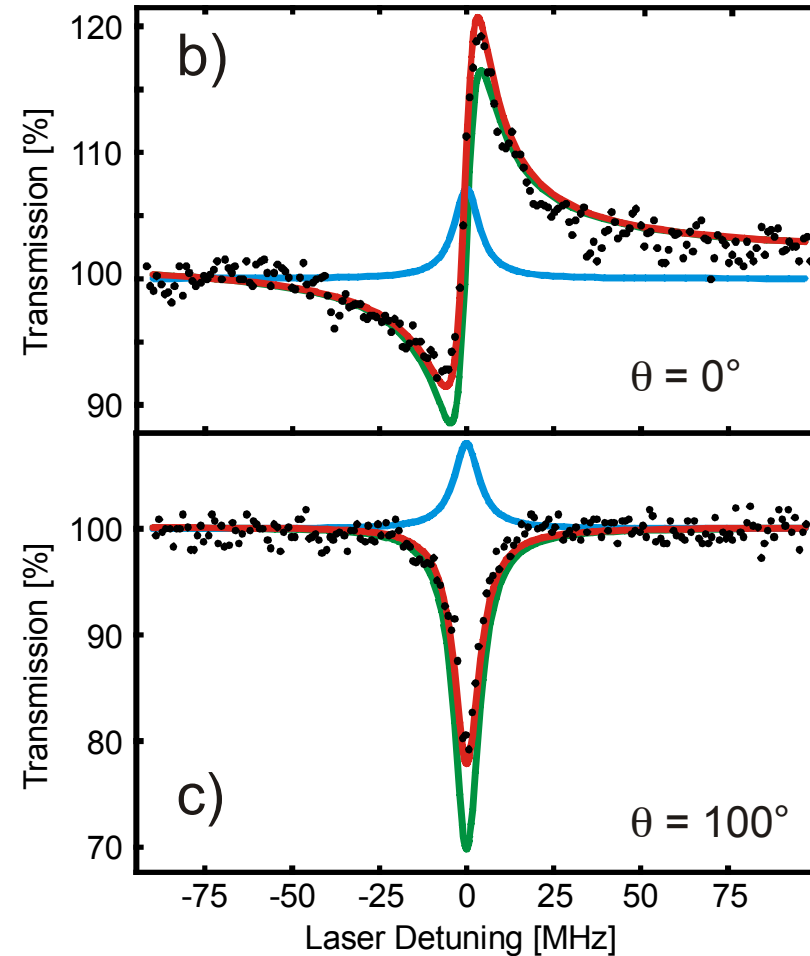
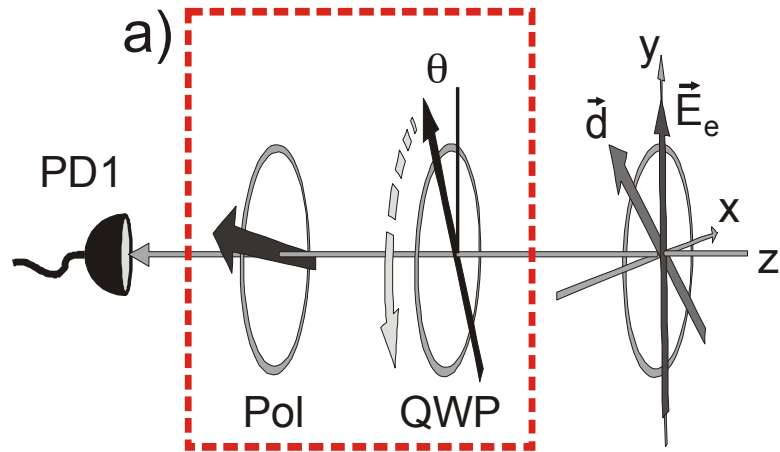
Transition from coherent to incoherent scattering



$2 \cdot 10^6$ cps
 $2 \cdot 10^7$ cps
 $2 \cdot 10^8$ cps
 $2 \cdot 10^9$ cps

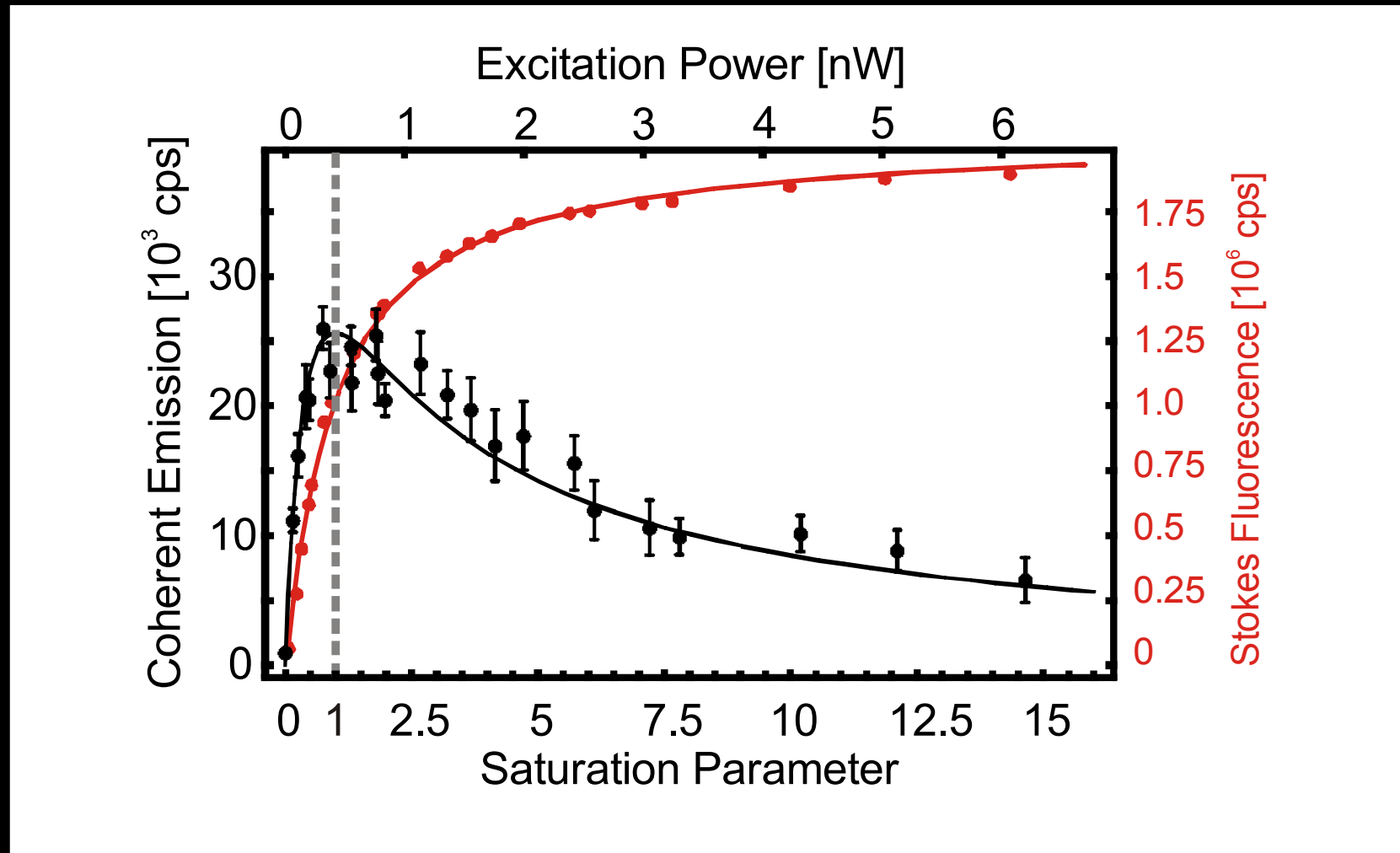
The molecule can maximally deal with one photon per two lifetimes (10 ns). At higher excitation intensity the extinction signal saturates.

Separating the coherent and incoherent parts



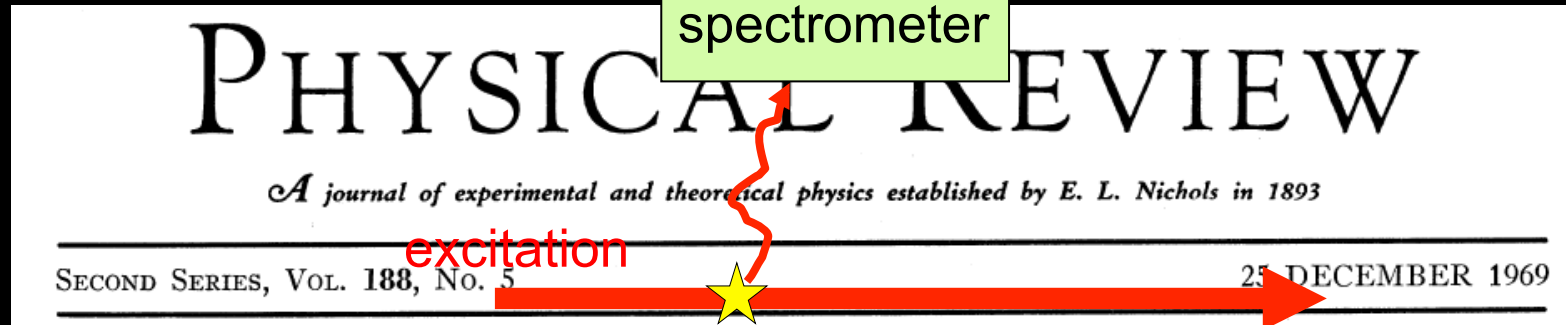
G. Wrigge, I. Gerhardt, J. Hwang, G. Zumofen, V. Sandoghdar,
Nature Physics 4, 60 (2008).

Transition from coherent to incoherent scattering



G. Wrigge, I. Gerhardt, J. Hwang, G. Zumofen, V. Sandoghdar,
Nature Physics **4**, 60 (2008).

Resonance fluorescence and Mollow Triplet



Power Spectrum of Light Scattered by Two-Level Systems

B. R. MOLLOW*

National Aeronautics and Space Administration, Lincoln Laboratory, Lexington, Massachusetts

The power spectrum of light scattered by a two-level system driven exactly on resonance. The power spectrum is shown as a function of the detuning $\nu - \omega_0$ of the scattered light from the resonance frequency ω_0 . The power spectrum is shown for three different values of the Rabi frequency Ω : $\Omega = K$, $\Omega = 5K$, and $\Omega = 3K$. The power spectrum is shown as a function of the detuning $\nu - \omega_0$ of the scattered light from the resonance frequency ω_0 . The power spectrum is shown for three different values of the Rabi frequency Ω : $\Omega = K$, $\Omega = 5K$, and $\Omega = 3K$.

chromatic
g field via
magnetic field
mic dipole
analogous

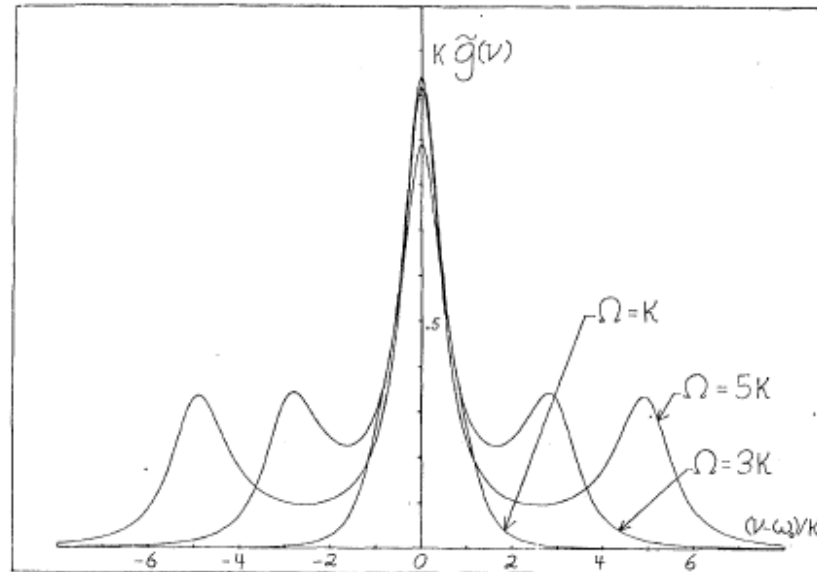
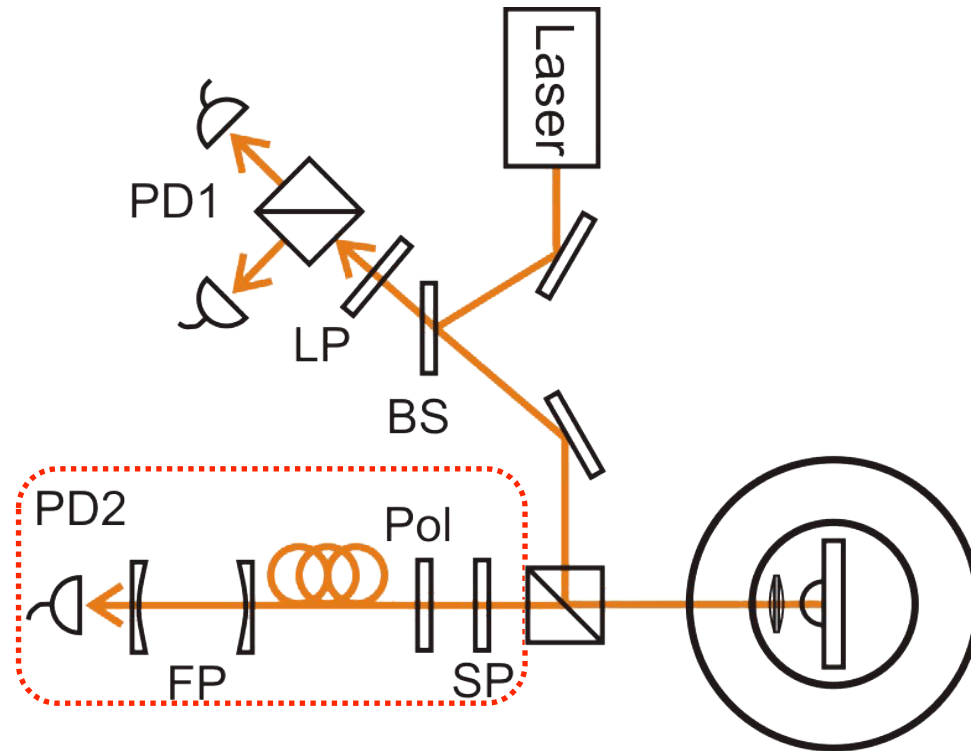


FIG. 1. Spectral density $\bar{g}(\nu)$ for a two-level atom driven exactly on resonance.

Direct detection of the Mollow triplet



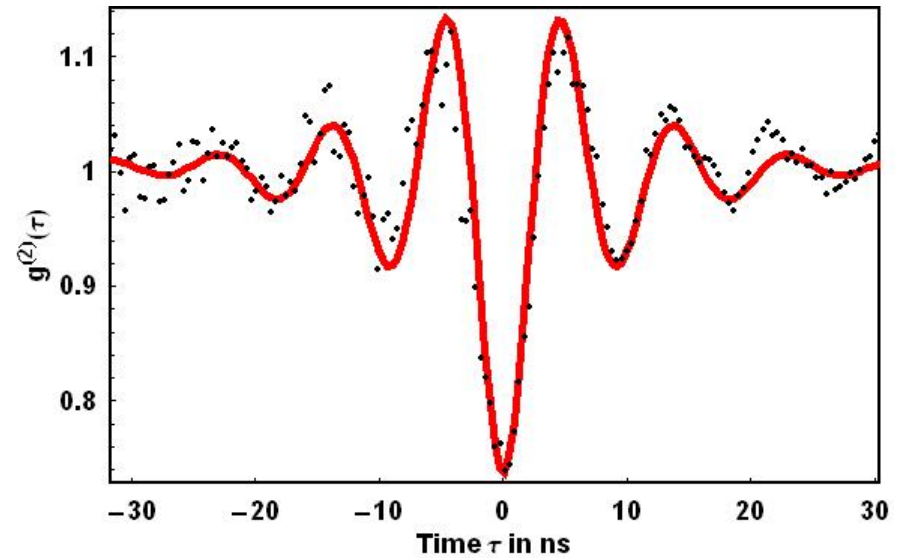
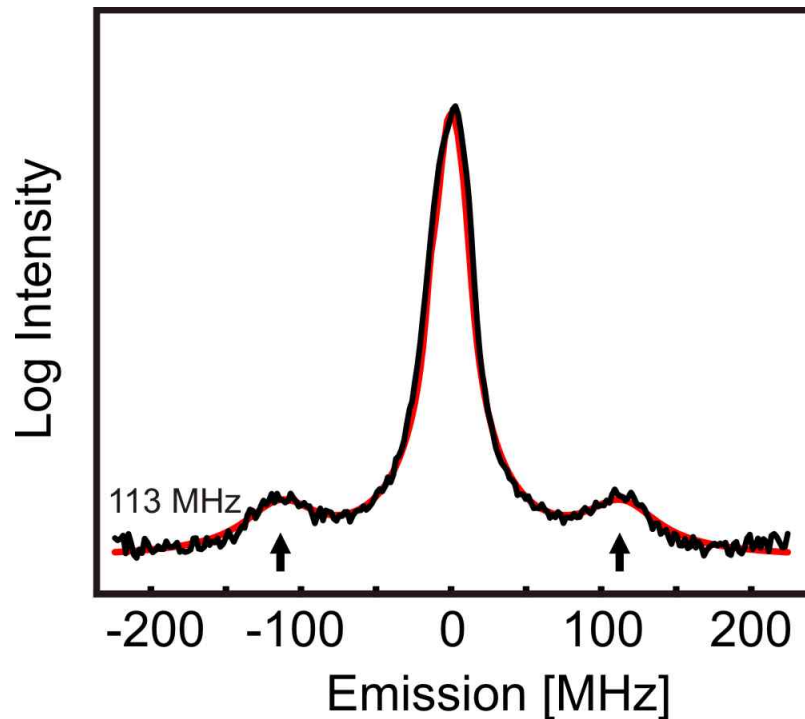
PD: photodetector

SP/LP: short/long pass

Pol: polarizer

FP: Fabry-Perot cavity

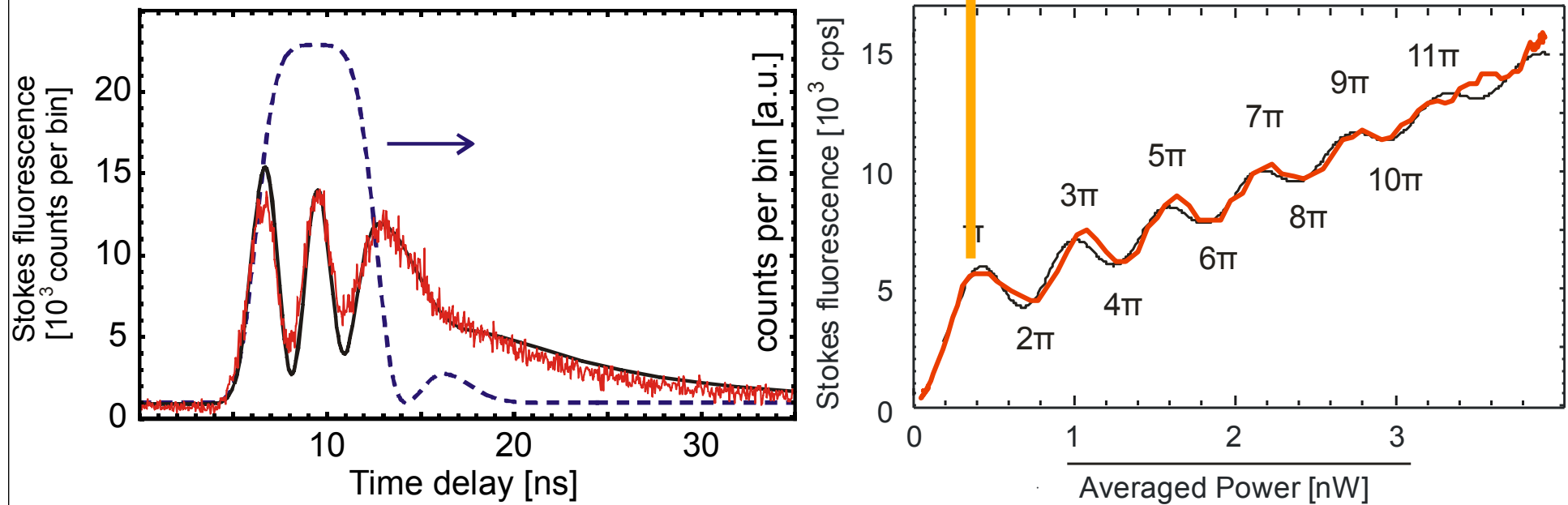
The first observation of resonance fluorescence and the Mollow triplet in the solid-state



G. Wrigge, I. Gerhardt, J. Hwang, G. Zumofen, V. Sandoghdar,
Nature Physics **4**, 60 (2008).

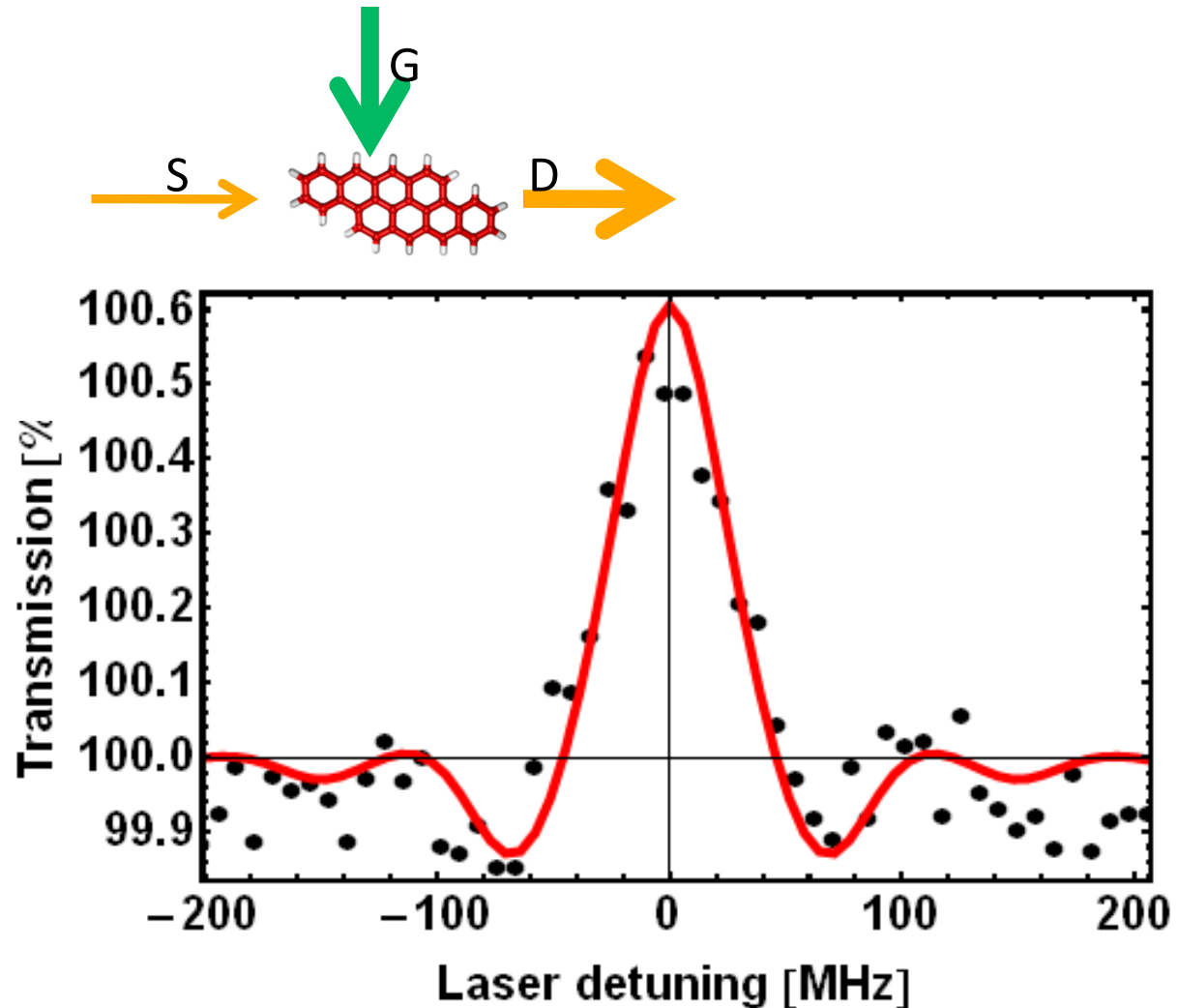
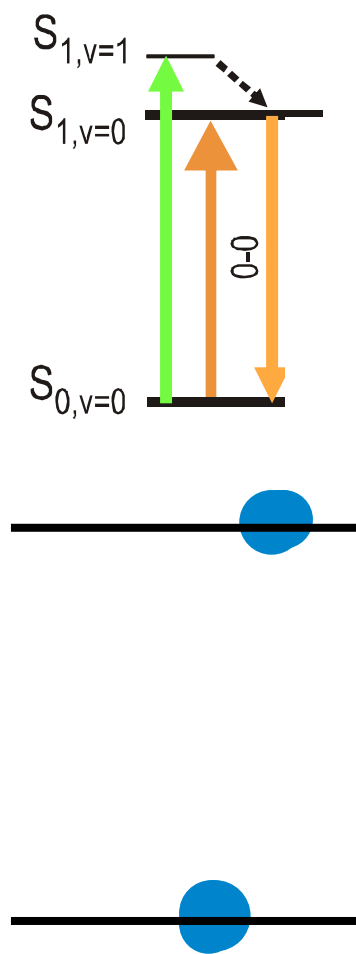
Pulsed excitation of single molecules and Rabi oscillations

Only 500 pump photons



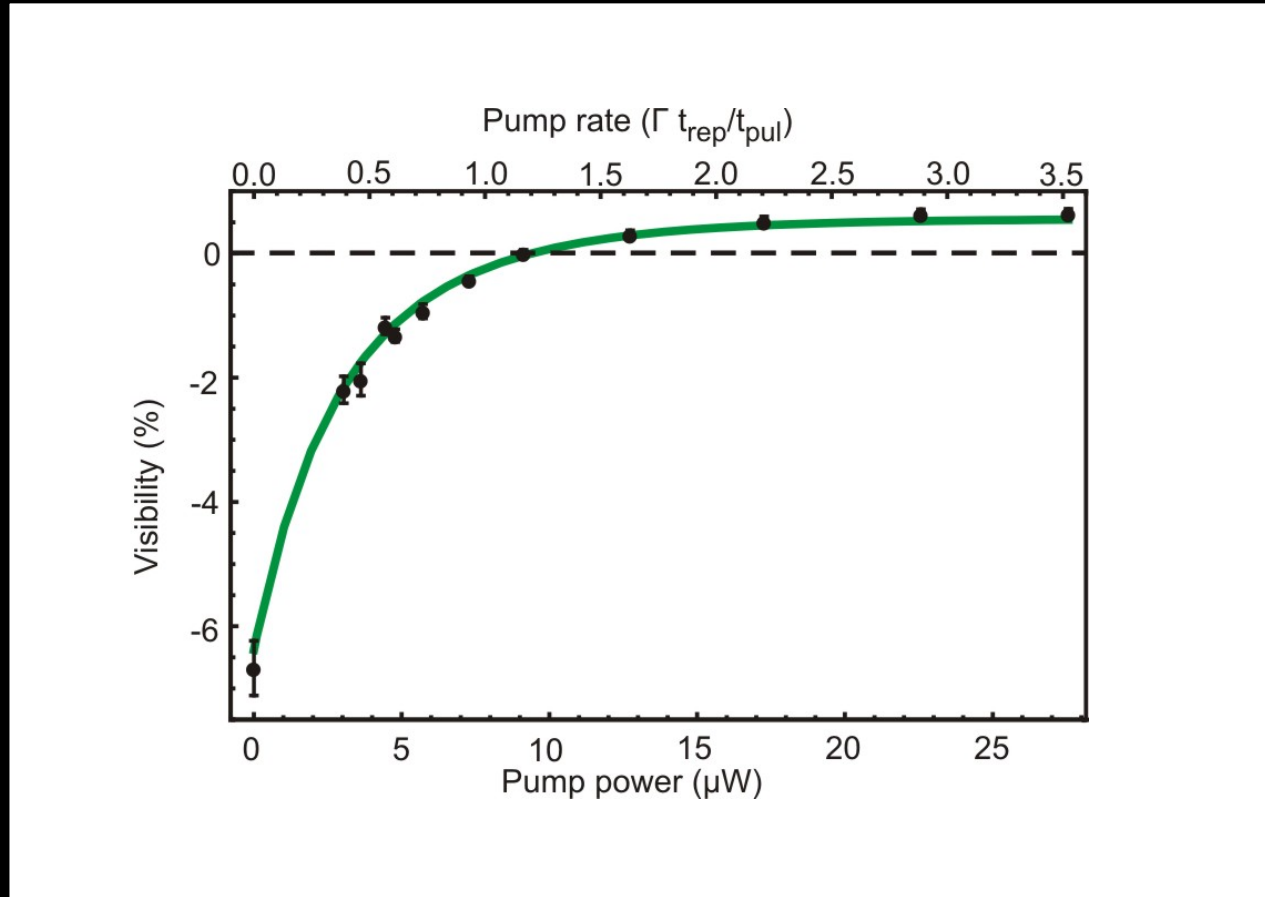
I. Gerhardt, G. Wrigge, G. Zumofen, J. Hwang, A. Renn, V. Sandoghdar,
Phys. Rev. A, **79**, 011402(R) (2009).

Amplification of light by a single molecule



J. Hwang, M. Pototschnig, R. Lettow, G. Zumofen, A. Renn, S. Götzinger, V. Sandoghdar, *Nature* 460, 76 (2009).

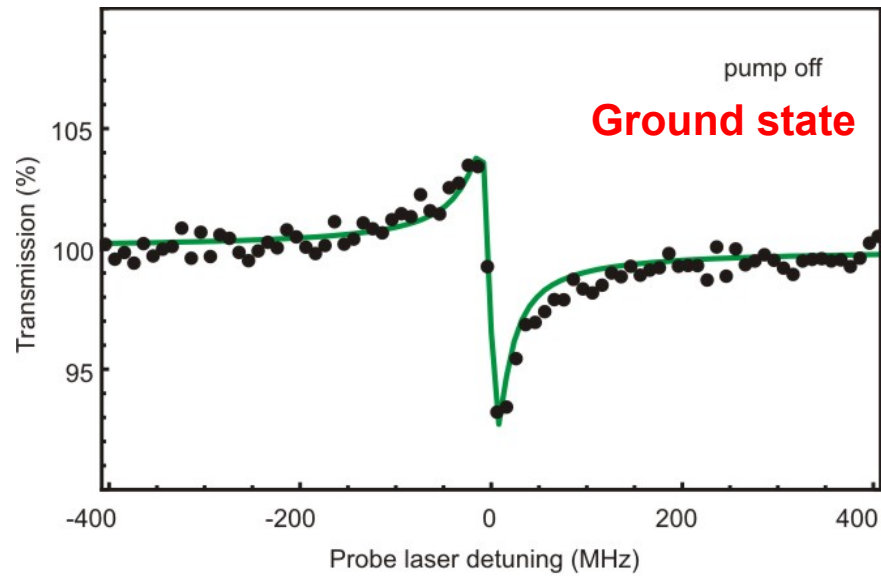
A single molecule as an optical transistor



J. Hwang, M. Pototschnig, R. Lettow, G. Zumofen, A. Renn, S. Götzinger, V. Sandoghdar, *Nature* **460**, 76 (2009).

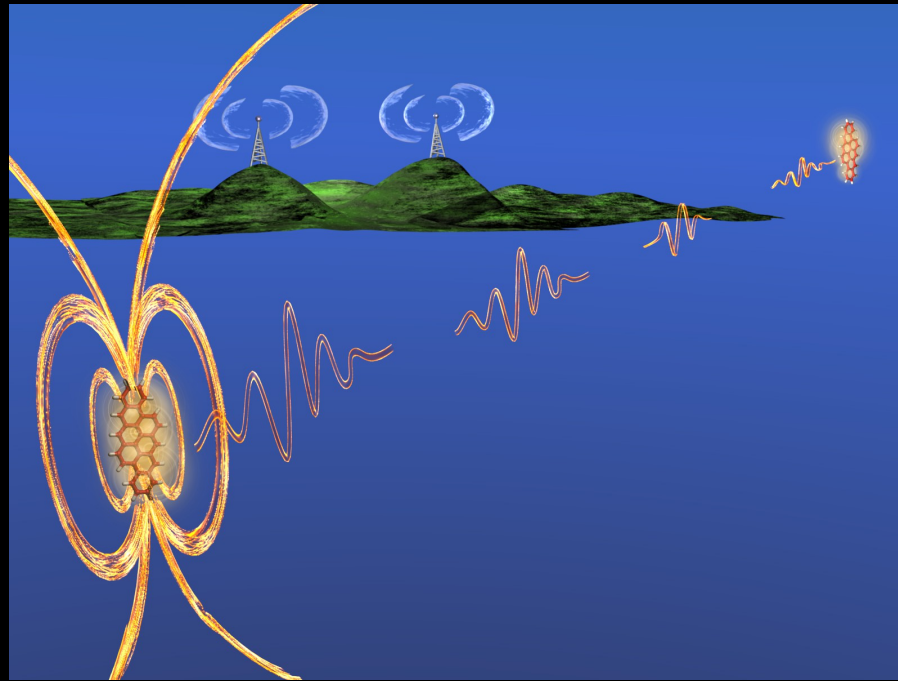
The coherent nature of stimulated emission

Adjust the phase to get a dispersive spectrum

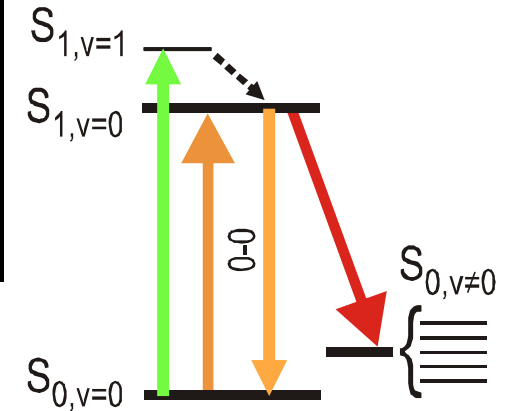


Far-field coupling of two single molecules

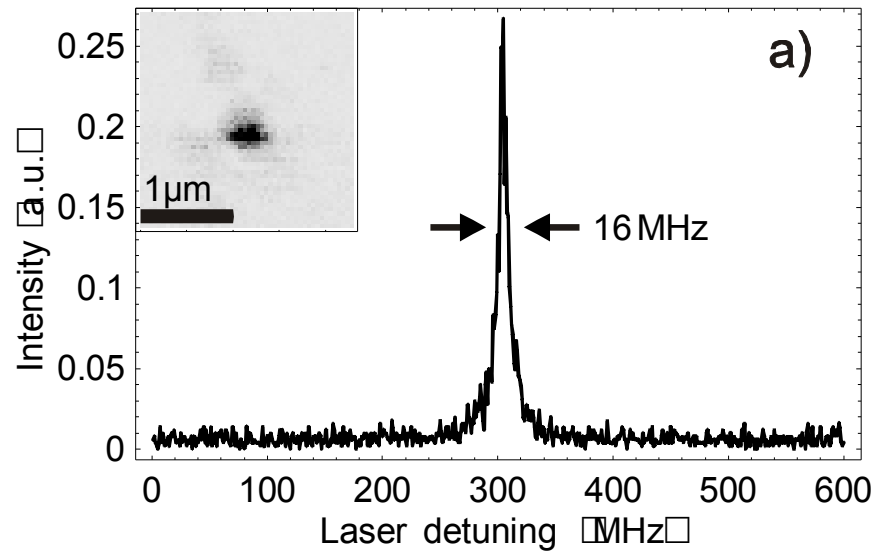
Spectroscopy with single photons



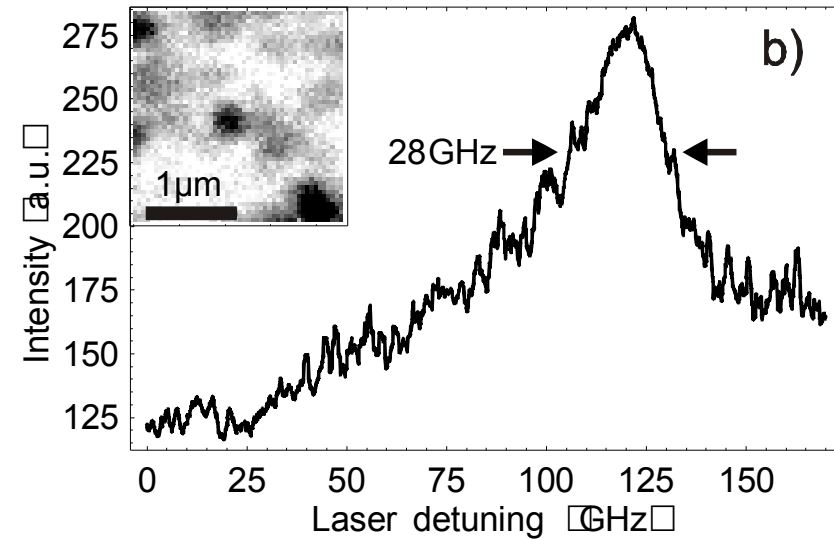
Collection of the lifetime-limited fluorescence



0-0 excitation



0-1 excitation

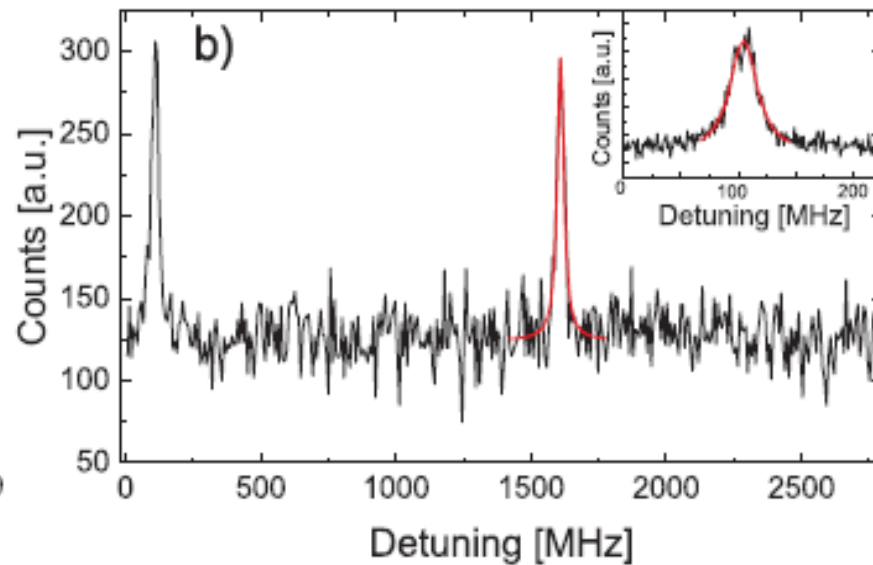


R. Lettow, et al., *Opt. Express* **15**, 15842 (2007).

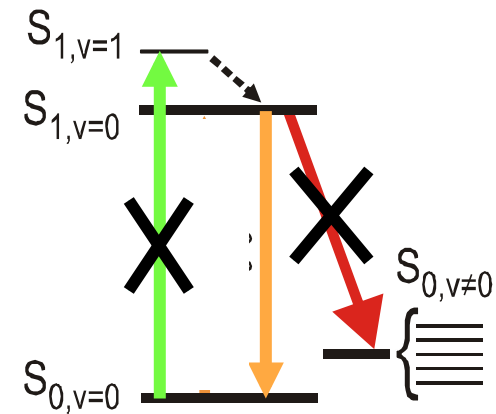
Single molecule as a single-photon source

Narrow-band single photons

Narrow emission

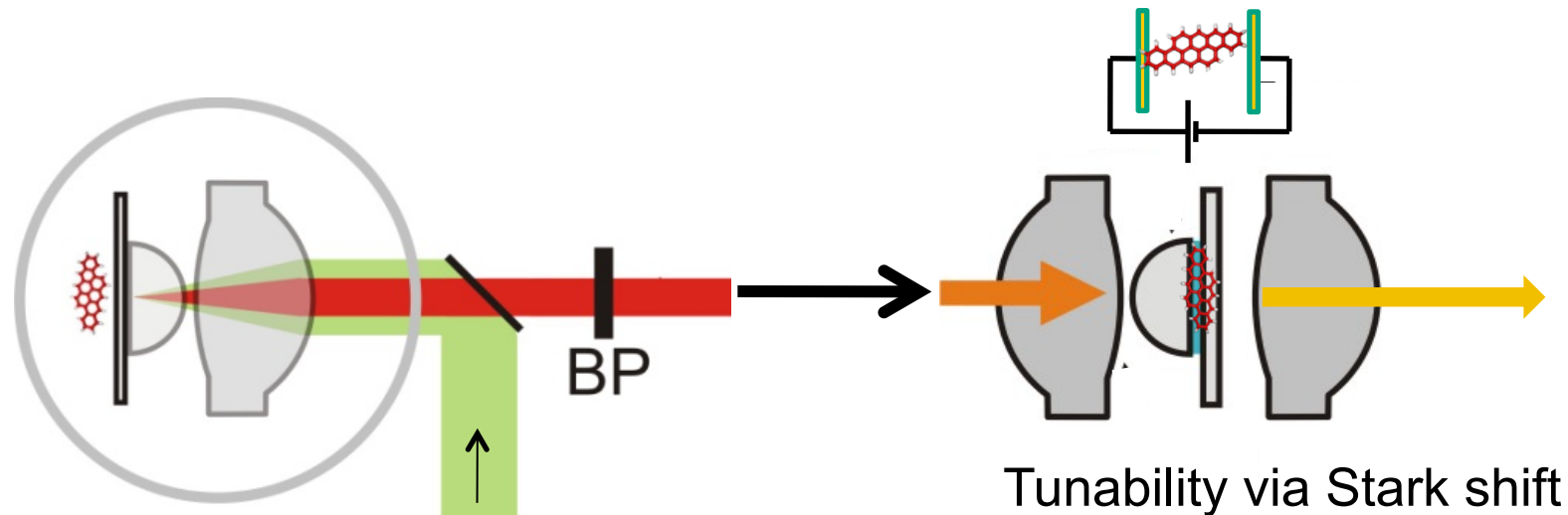


ZPL scanned with Fabry-Perot Cavity

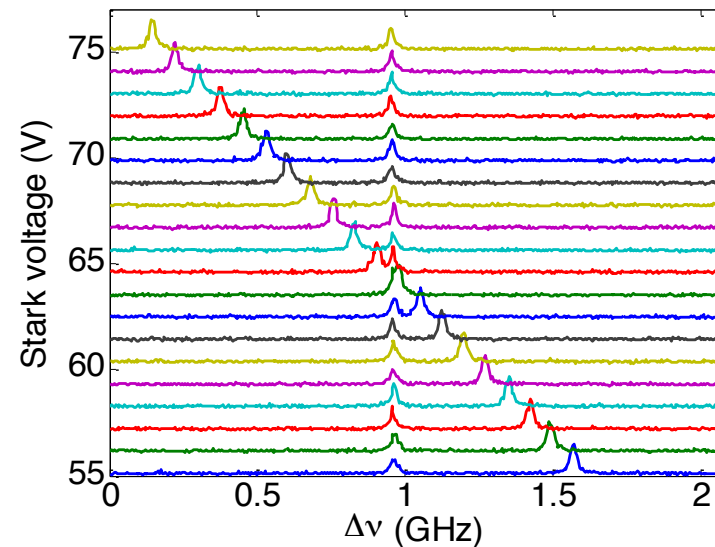


R. Lettow, V. Ahte, R. Pfab, A. Renn, E. Ikonen, S. Götzinger, V. Sandoghdar
Opt. Express **15**, 15842 (2007).

Exciting a single molecule with a single-photon source

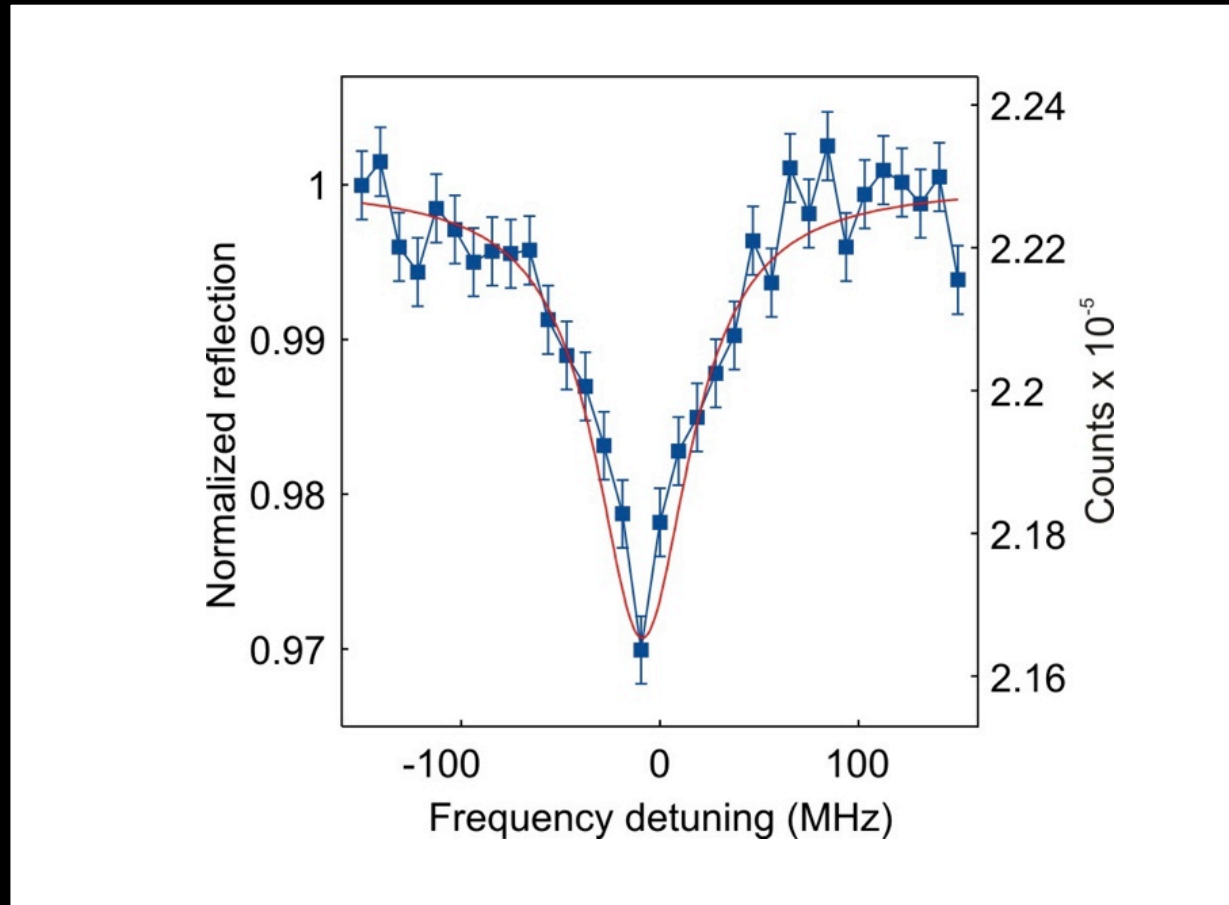


- Lifetime-limited photons
'gas-phase-like in the solid state'
- Cryogenic temperature \Rightarrow no bleaching
- High brightness, count rates > 1 Mcps



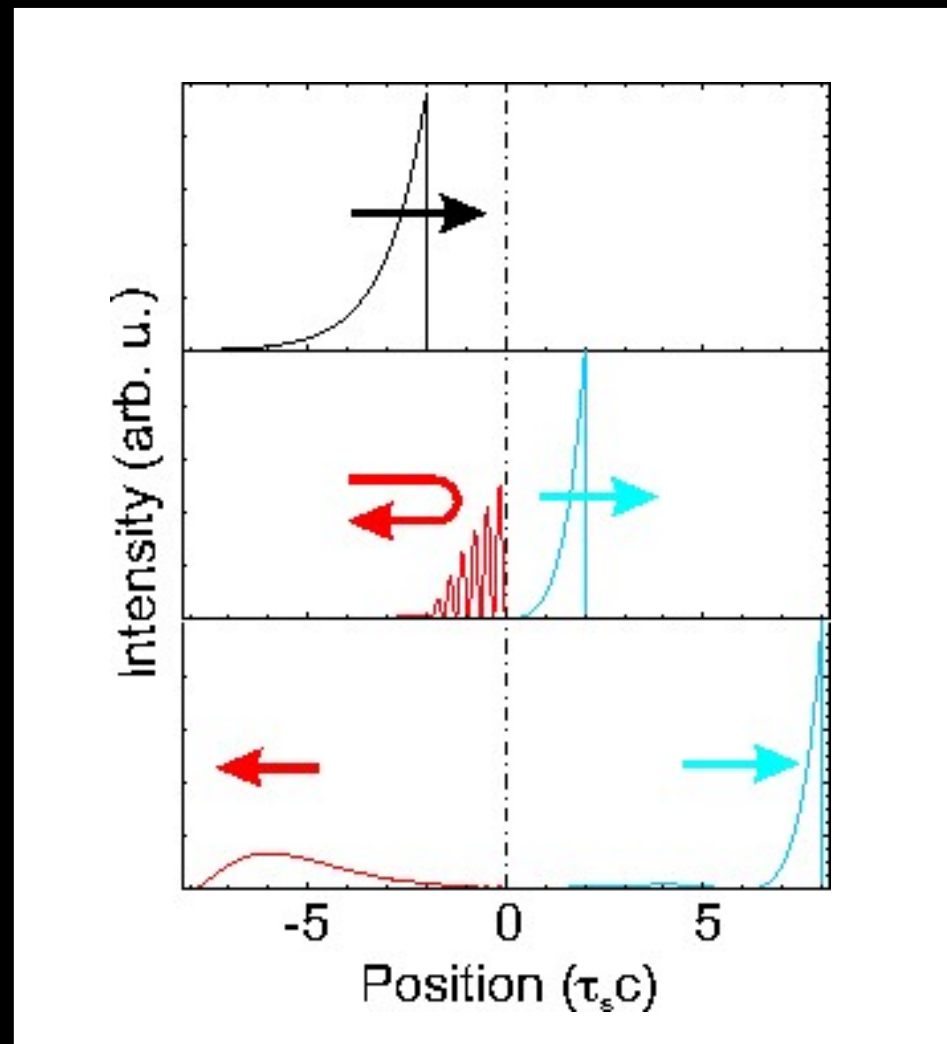
R. Lettow, V. Ahtee, R. Pfab, A. Renn, E. Ikonen, S. Göttinger, V. Sandoghdar, *Opt. Express* **15**, 15842 (2007).

Extinction spectrum of a single molecule using single photons from another molecule



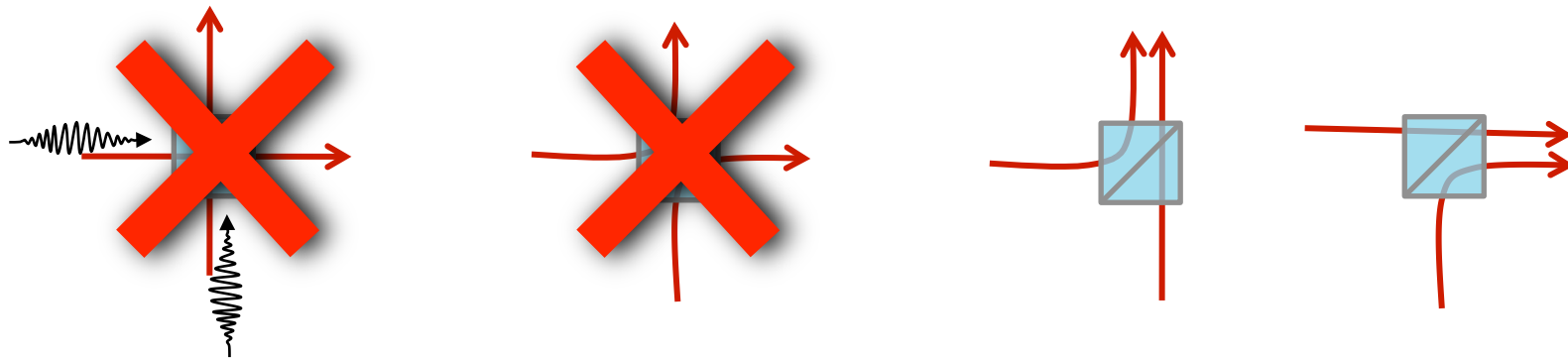
Y. Rezus, S. Walt, R. Lettow, A. Renn, G. Zumofen, S. Götzinger, V. Sandoghdar
Phys. Rev. Lett. **108**, 093601 (2012).

Bouncing of a photon from an atom



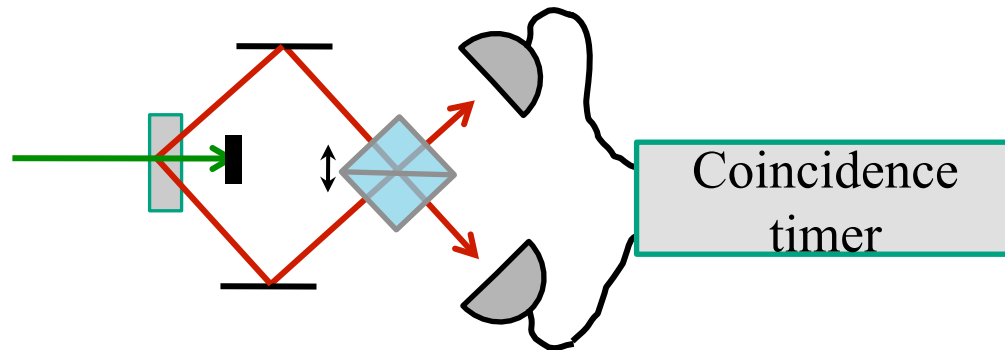
Y. Rezus, S. Walt, R. Lettow, A. Renn, G. Zumofen, S. Götzinger, V. Sandoghdar
Phys. Rev. Lett. **108**, 093601 (2012).

Two-photon interference

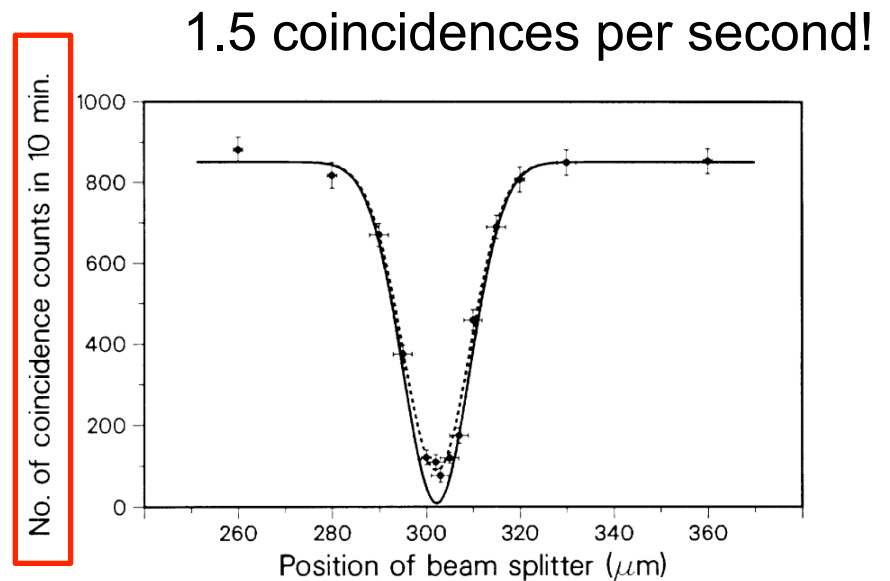


Original HOM experiment

- Originally the two photons were generated by Spontaneous Parametric Down Conversion (SPDC) in a nonlinear crystal



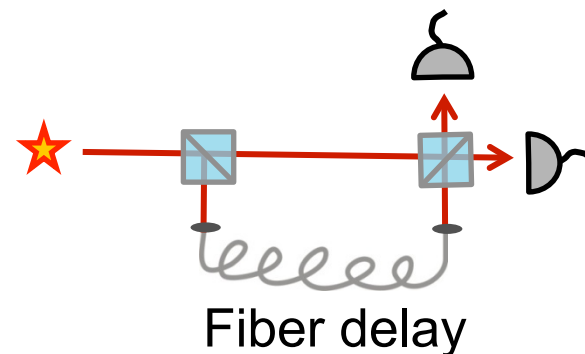
- Indistinguishable photons:
 - Frequency
 - Polarization
 - Coherence length
 - Spatial mode



TPI with true single-photon sources

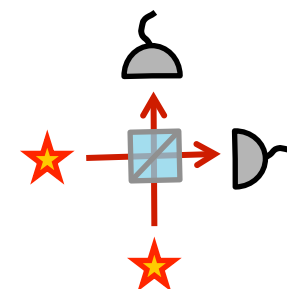
- Two-photon interference with one single-photon source

Molecules: Zumbusch
Gas-phase ions: Rempe

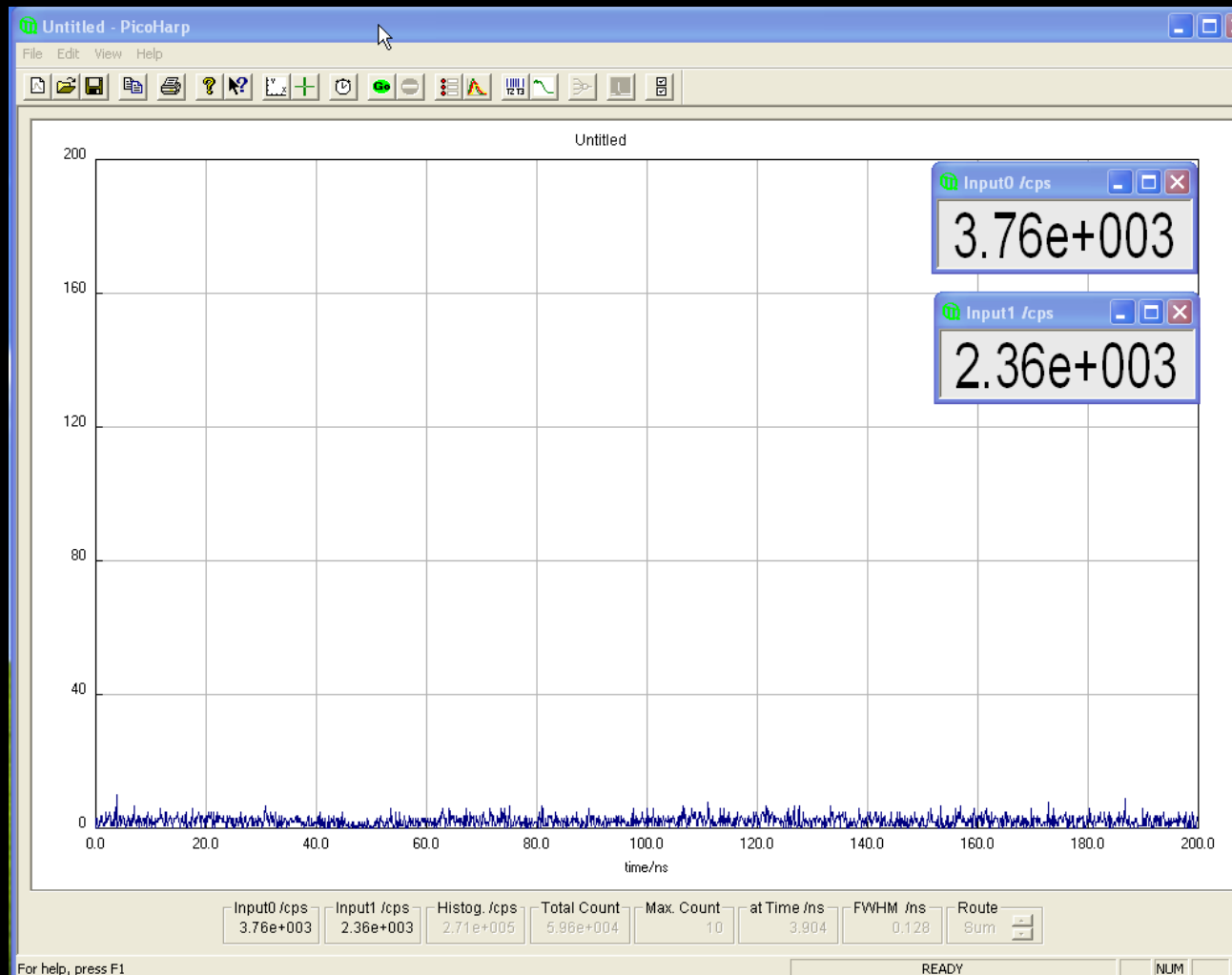


- Two-photon interference with independent single-photon sources

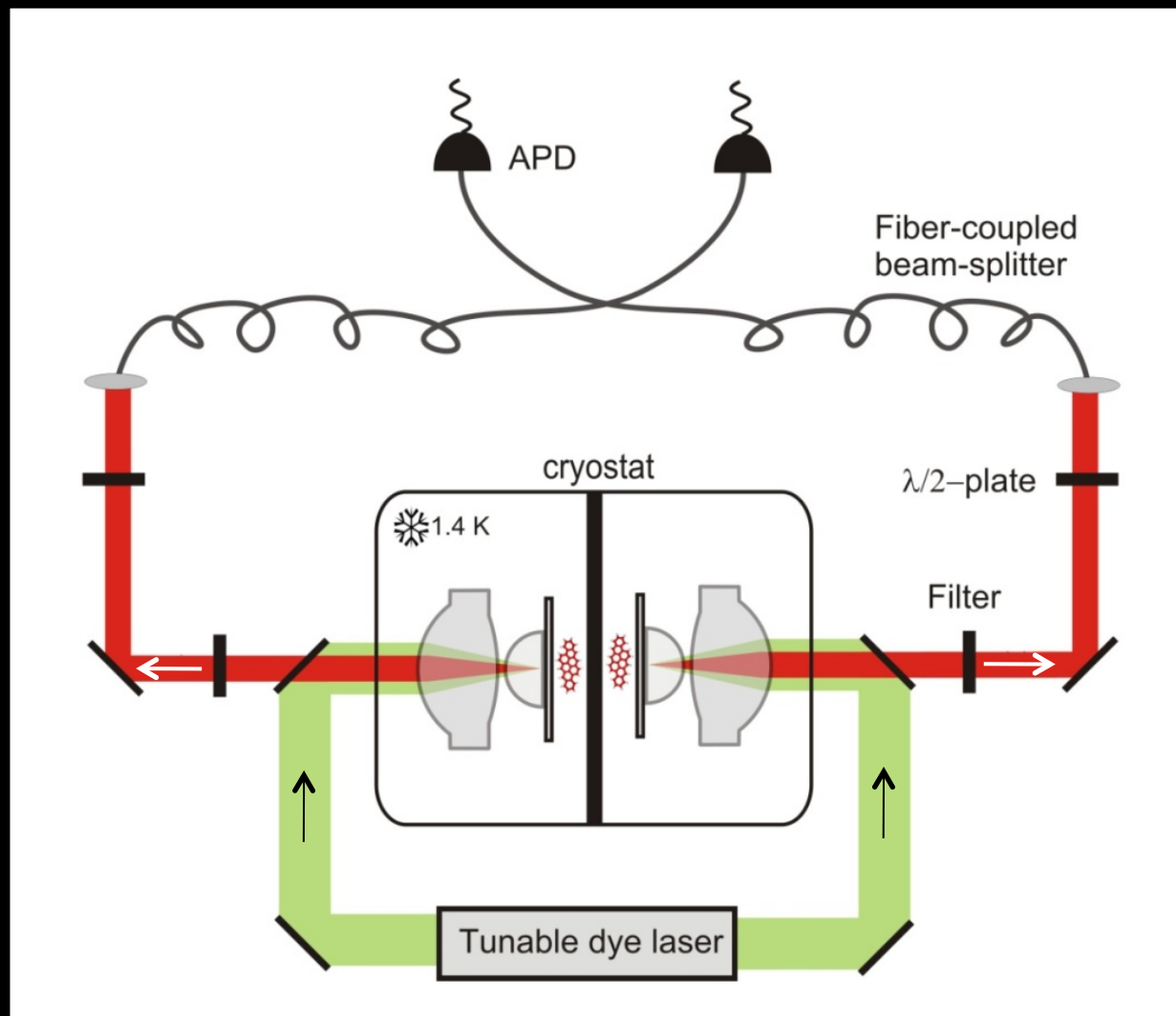
Gas-phase atoms and ions: Monroe, Grangier



Very bright source: Fast antibunching measurement



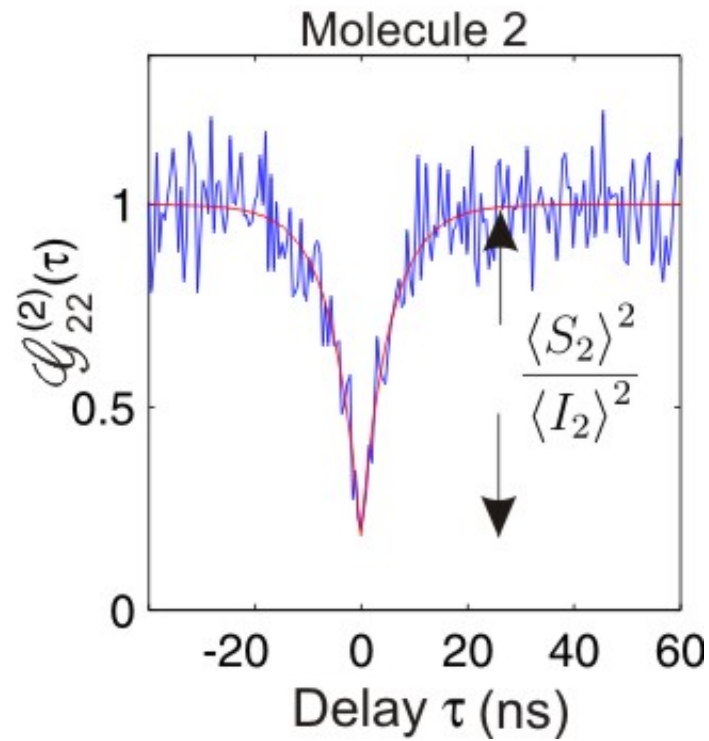
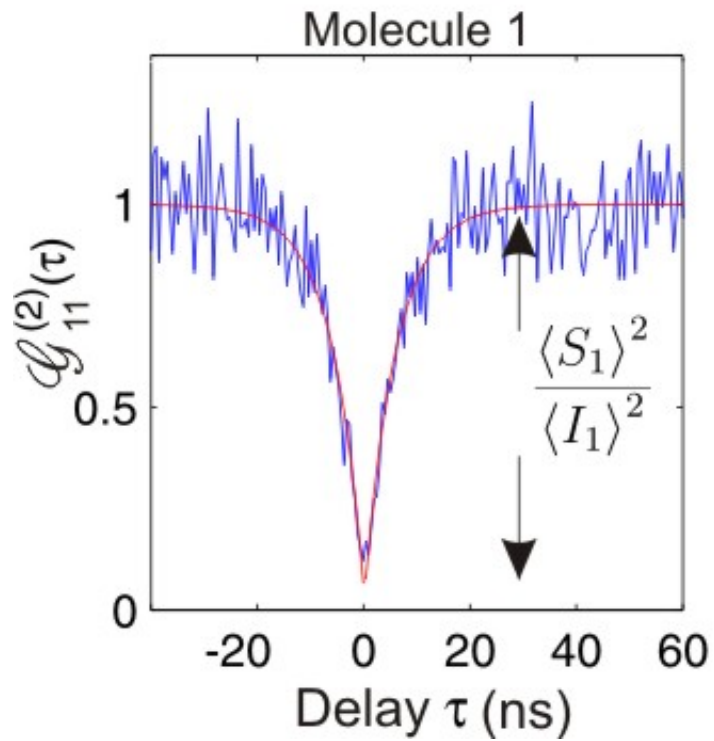
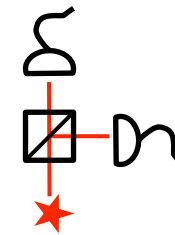
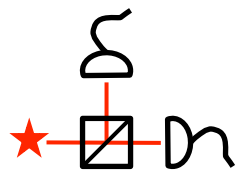
Hong-Ou-Mandel Two-Photon Interference



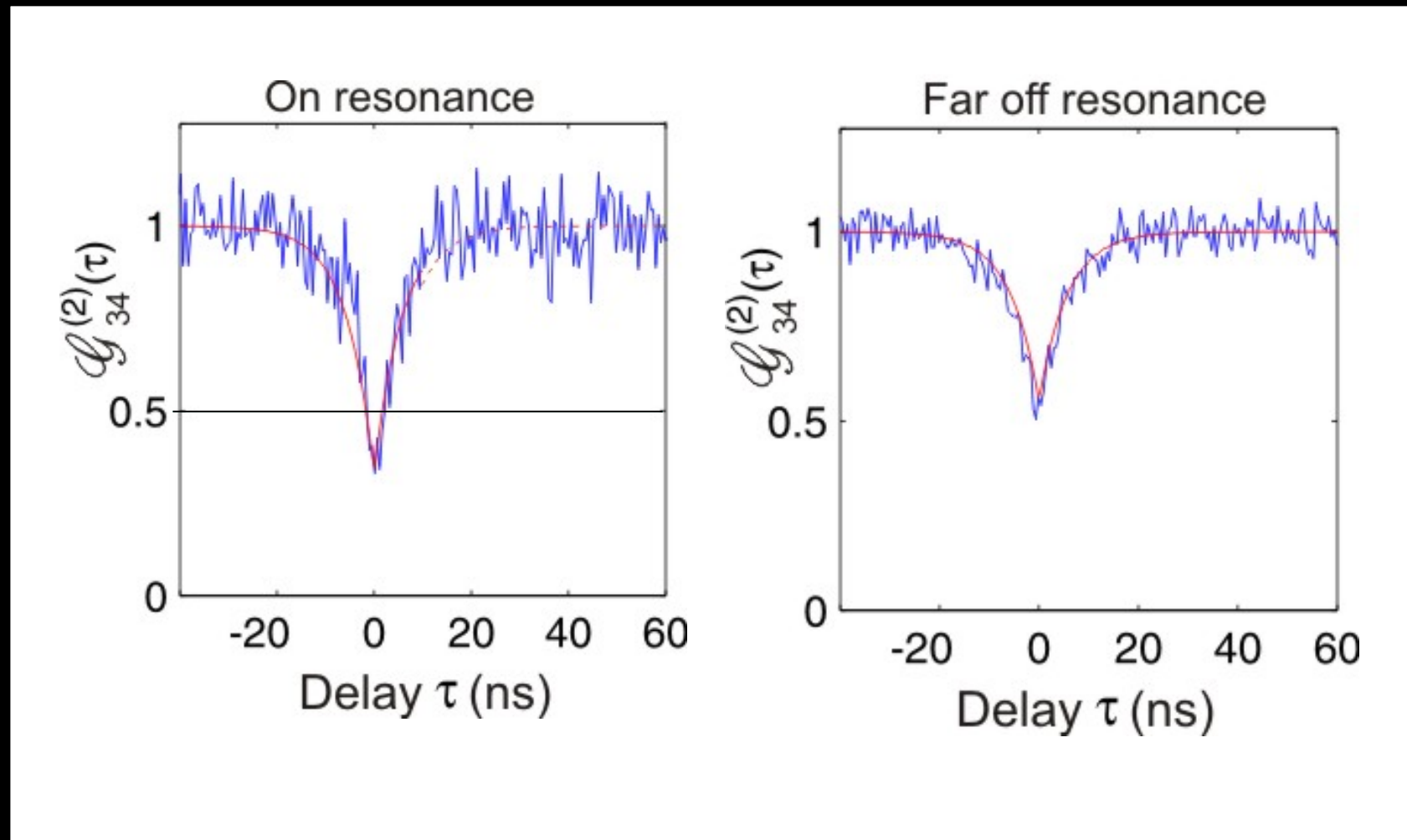
R. Lettow, V. Ahtee, R. Pfab, A. Renn, E. Ikonen, S. Götzinger, V. Sandoghdar
Opt. Express 15, 15842 (2007).

Checking that each arm is antibunched

Continuous-wave excitation

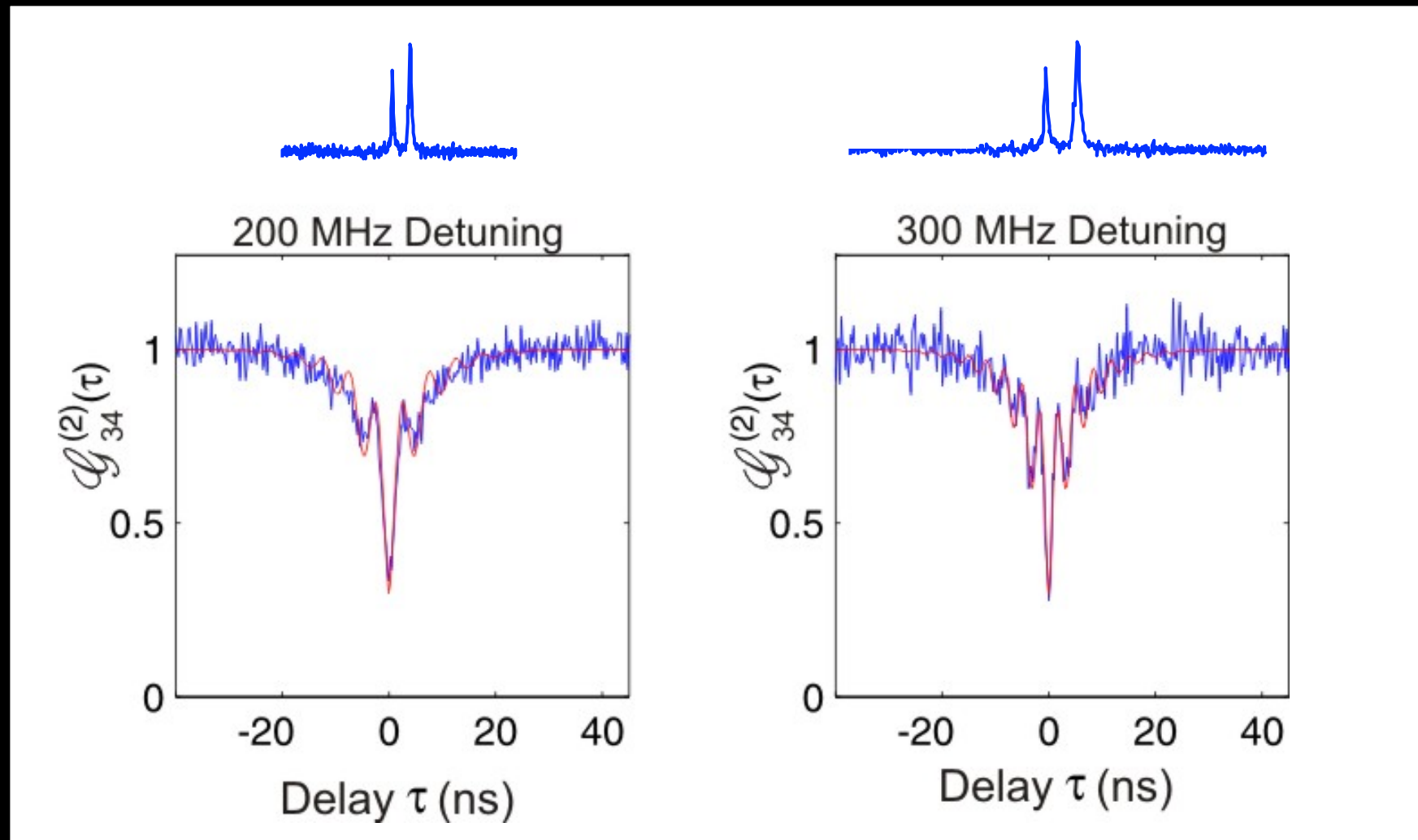


Proof of two-photon interference



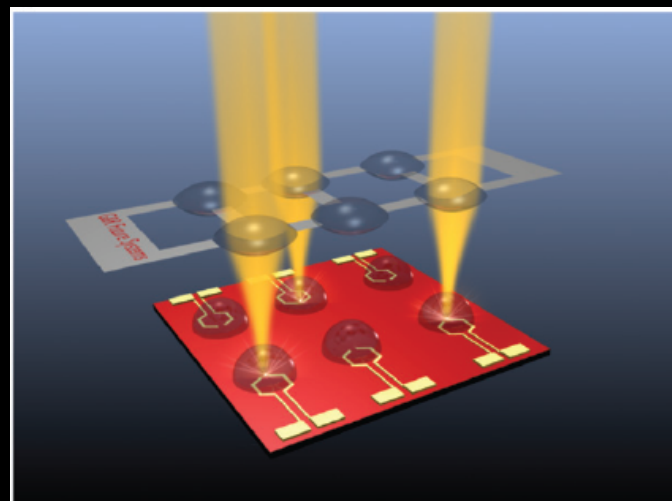
R. Lettow, Y. Rezus, G. Zumofen, A. Renn, E. Ikonen, S. Götzinger, V. Sandoghdar, *Phys. Rev. Lett.* **104**, 123605 (2010).

Controlled tuning of the photons

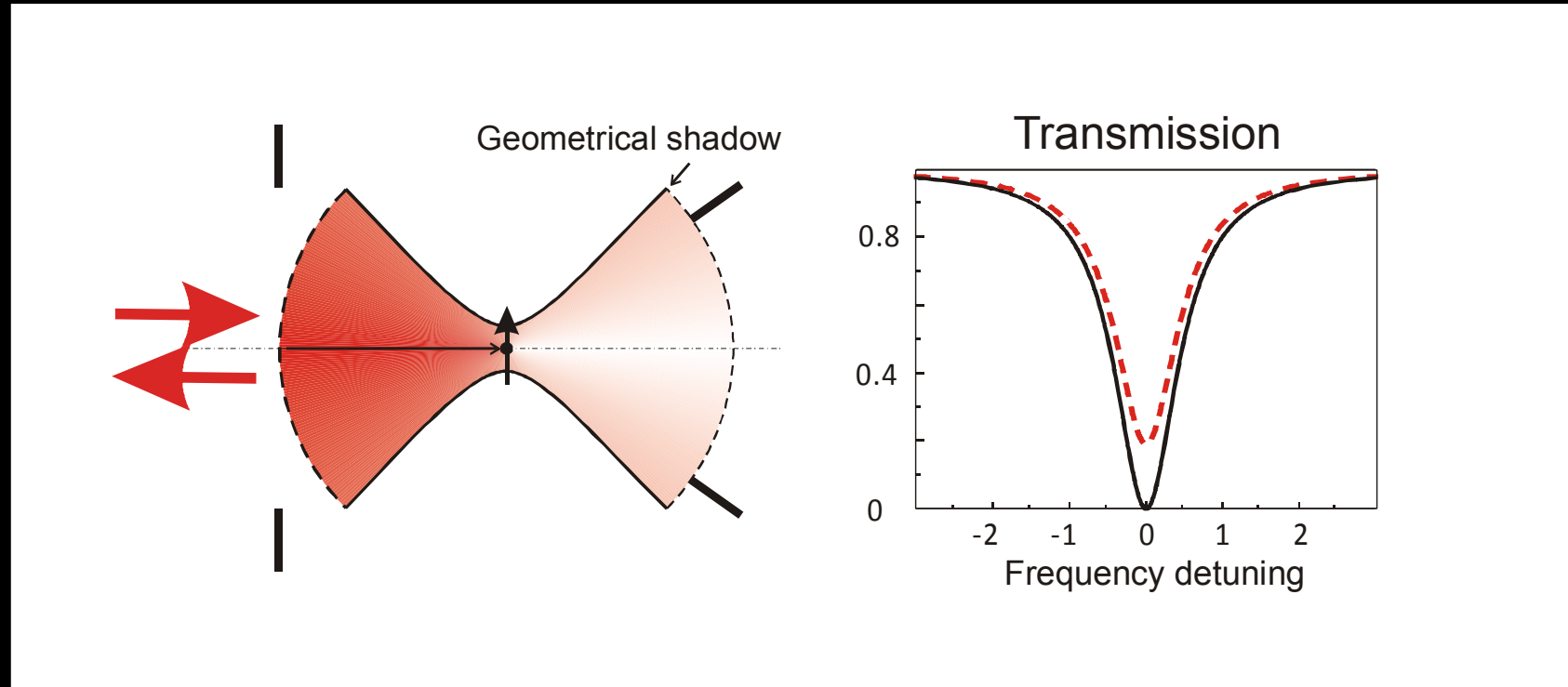


R. Lettow, Y. Rezus, G. Zumofen, A. Renn, E. Ikonen, S. Götzinger, V. Sandoghdar, *Phys. Rev. Lett.* **104**, 123605 (2010).

From two photons to
few photons



Perfect reflection of light by a single oscillating dipole



PHYSICAL REVIEW LETTERS **101**, 180404 (2008)



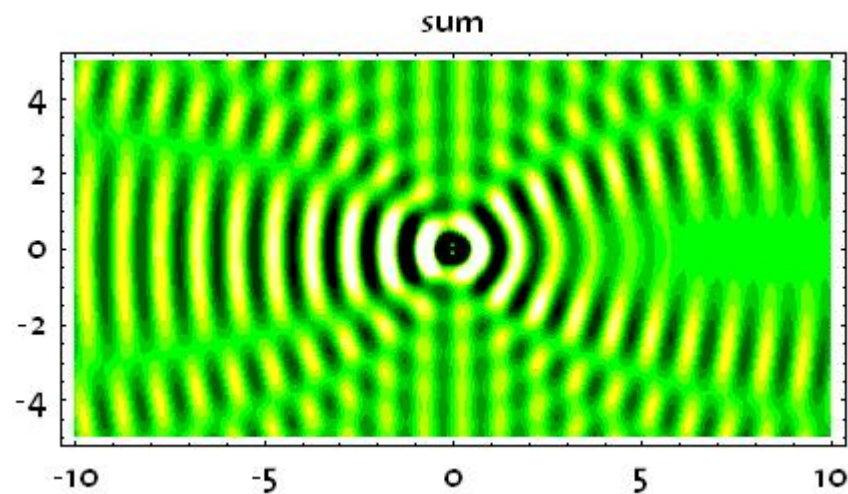
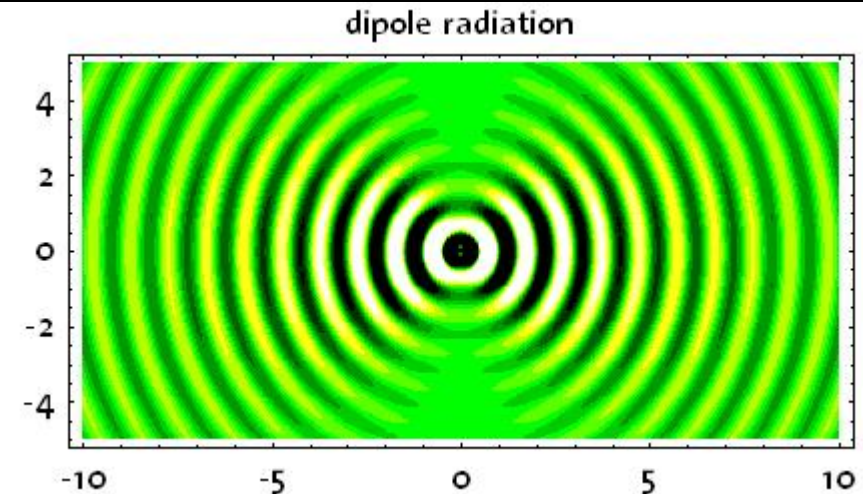
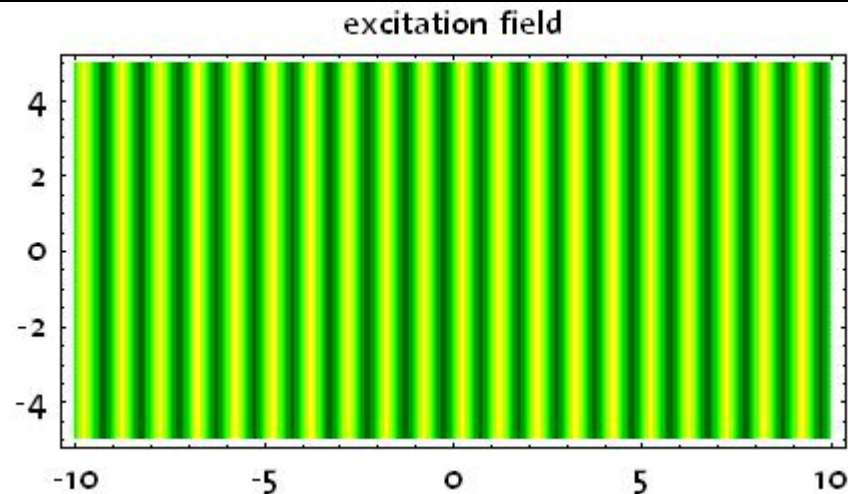
Perfect Reflection of Light by an Oscillating Dipole

G. Zumofen, N. M. Mojarad, V. Sandoghdar, and M. Agio

Nano-Optics Group, Laboratory of Physical Chemistry, ETH Zurich, CH-8093 Zurich, Switzerland

(Received 19 May 2008)

The extinction signal is due to interference

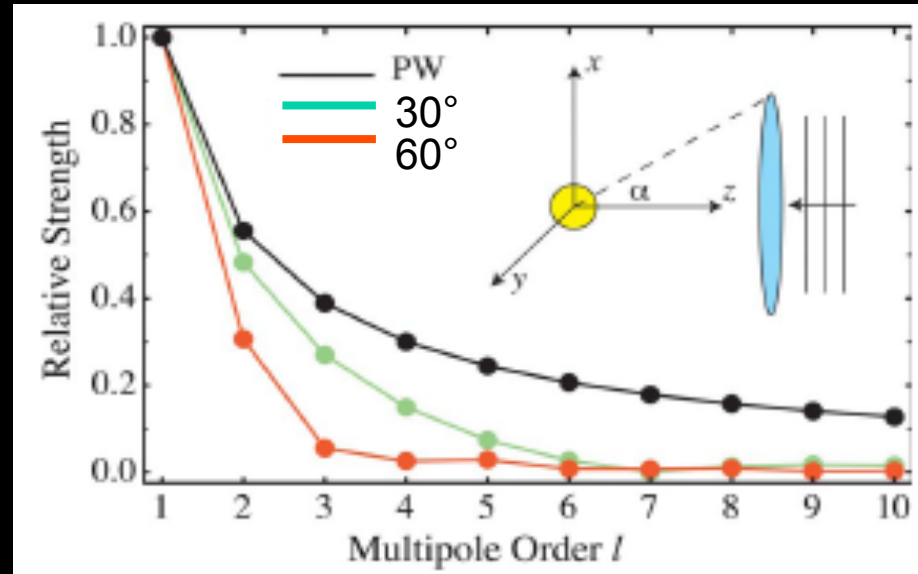


$$I_{total} = |E_{inc}|^2 + |E_{sca}|^2 + 2\text{Re}(E_{inc}E_{sca}^*)$$

↓
extinction

Complete extinction in the forward direction, hence the definition of the absorption cross section (optical theorem)

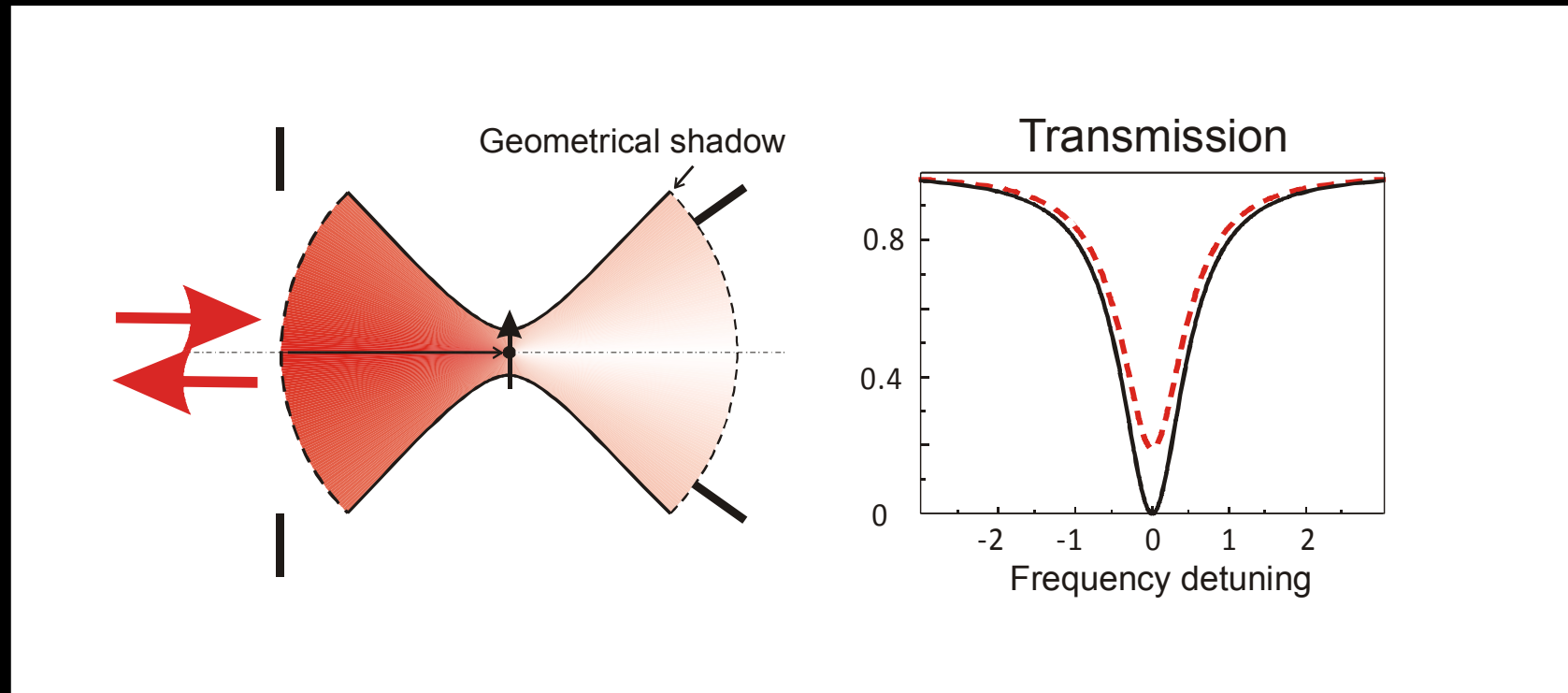
Multipolar content of a focused beam



A tightly focused beam matches the dipolar radiation well

N. M. Mojarad, V. Sandoghdar, M. Agio, *JOSA B* 25, 651 (2008).

Perfect reflection of light by a single oscillating dipole



PHYSICAL REVIEW LETTERS **101**, 180404 (2008)



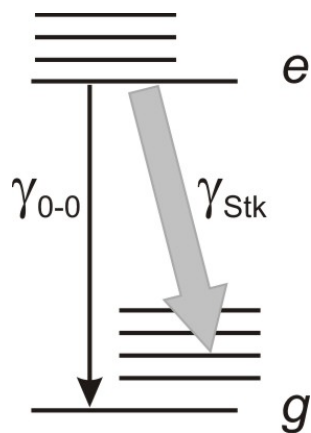
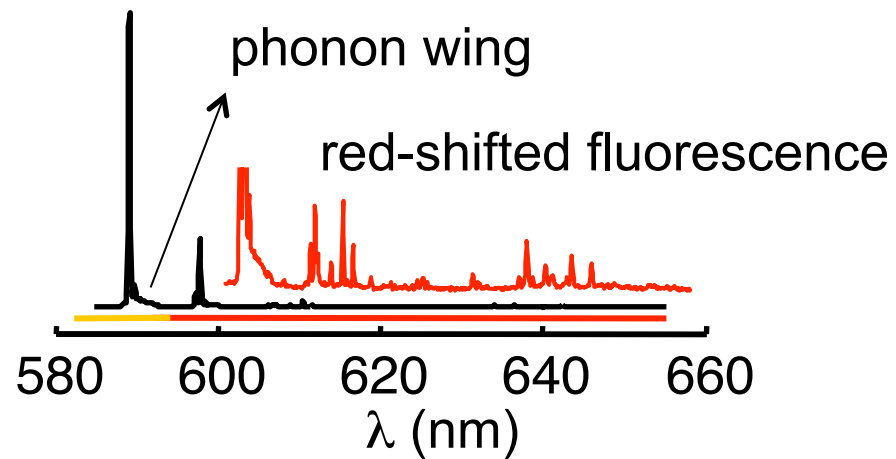
Perfect Reflection of Light by an Oscillating Dipole

G. Zumofen, N. M. Mojarad, V. Sandoghdar, and M. Agio

Nano-Optics Group, Laboratory of Physical Chemistry, ETH Zurich, CH-8093 Zurich, Switzerland

(Received 19 May 2008)

Zero-phonon line



$$\sigma_{ext} = \frac{3\lambda^2}{2\pi} \left(\frac{\gamma_0}{\gamma_{hom}} \right) \alpha$$

Mode matching:

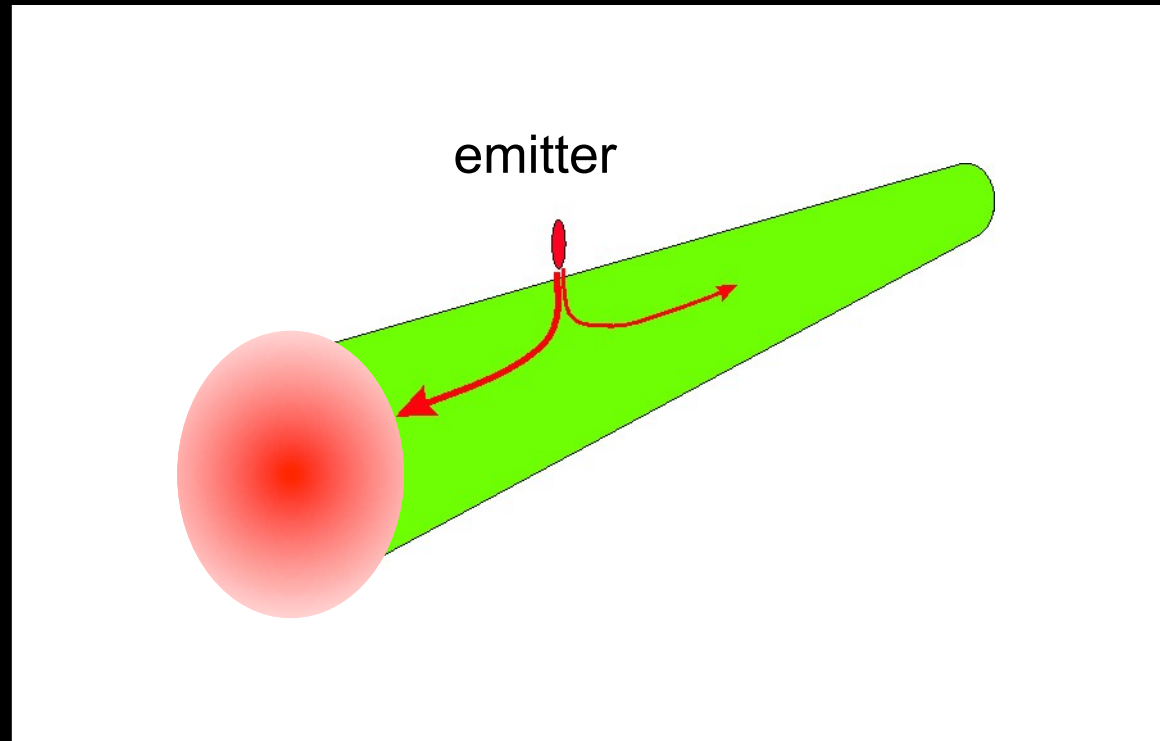
Match the light to the emitter

Match the emitter to the light

Coupling to nanofibers

Mode conversion: dipolar radiation to something like a TEM mode

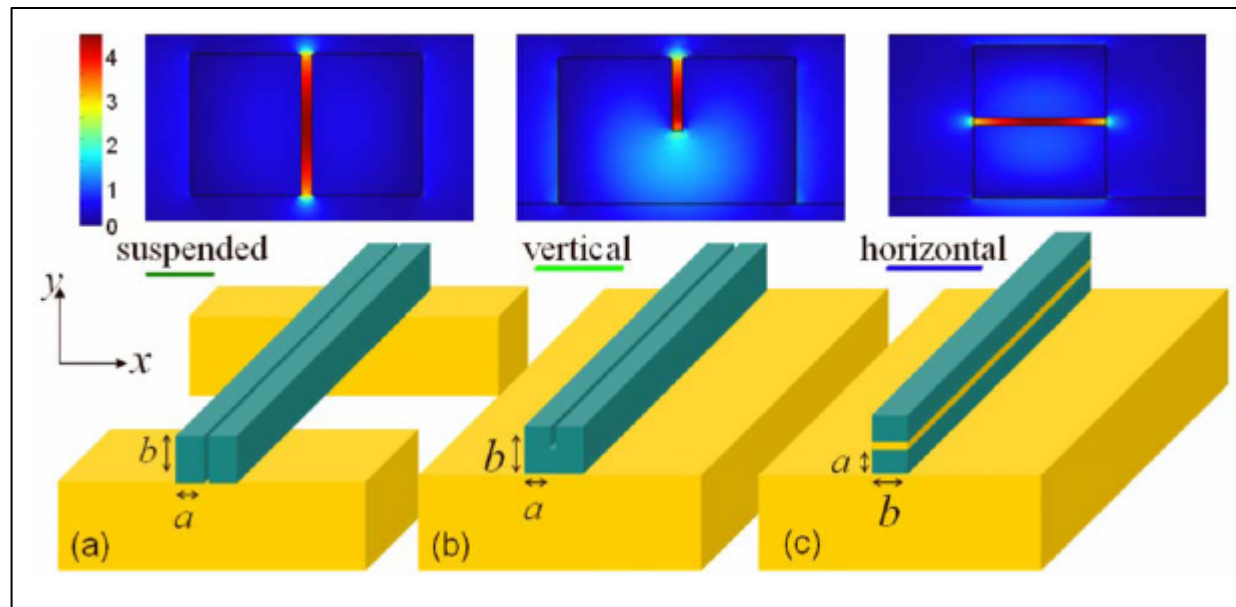
The higher the refractive index, the larger the coupling



PHYSICAL REVIEW A **80**, 011810(R) (2009)

Broadband waveguide QED system on a chip

Qimin Quan, Irfan Bulu, and Marko Lončar



Broadband enhancement of light emission in silicon slot waveguides

Young Chul Jun¹, Ryan M. Briggs², Harry A. Atwater², and Mark L. Brongersma^{1*}

27 April 2009 / Vol. 17, No. 9 / OPTICS EXPRESS 7479

Generation of single optical plasmons in metallic nanowires coupled to quantum dots

A. V. Akimov^{1,4*}, A. Mukherjee^{1*}, C. L. Yu^{2*}, D. E. Chang¹, A. S. Zibrov^{1,4}, P. R. Hemmer³, H. Park^{1,2} & M. D. Lukin¹

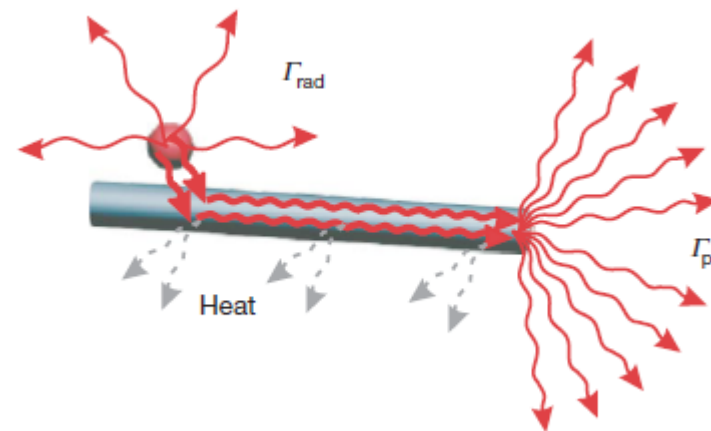
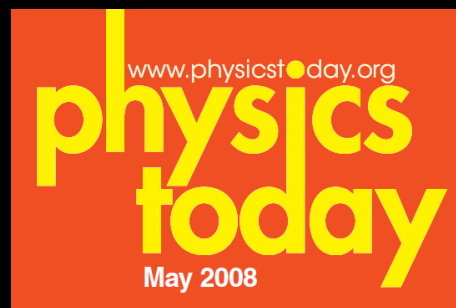
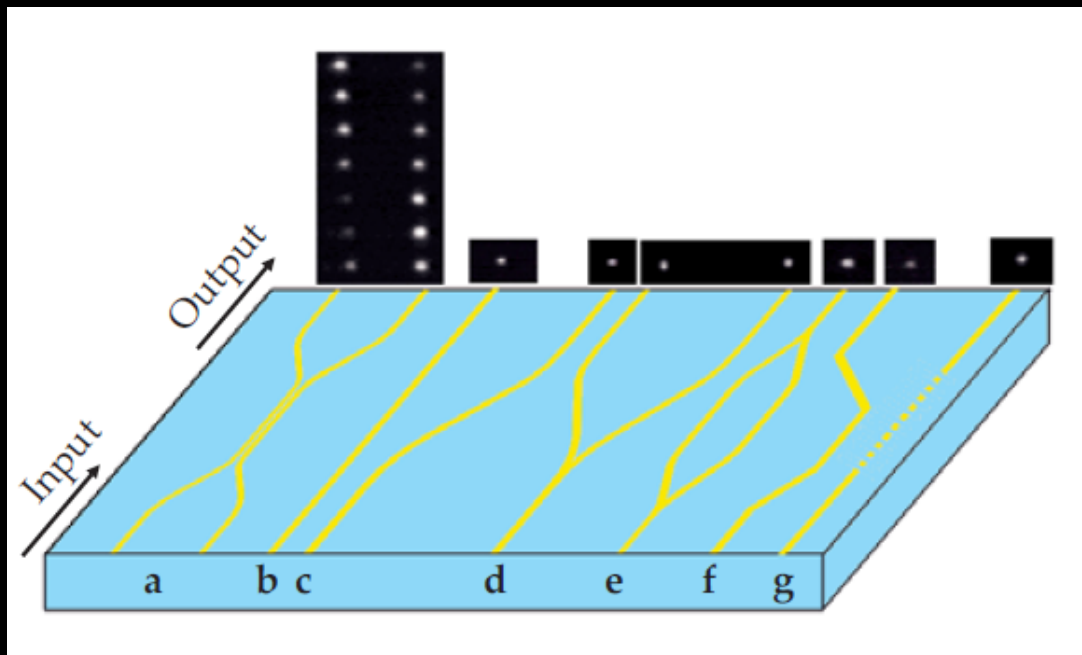


Figure 3 | Demonstration of single surface plasmon generation.

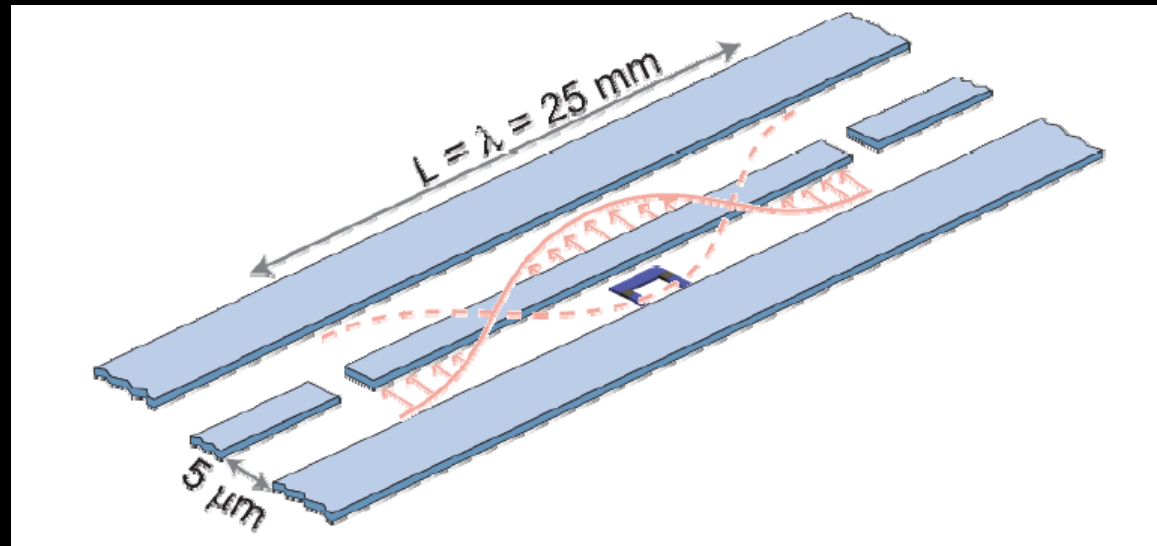
Surface-plasmon circuitry

Thomas W. Ebbesen, Cyriaque Genet, and Sergey I. Bozhevolnyi



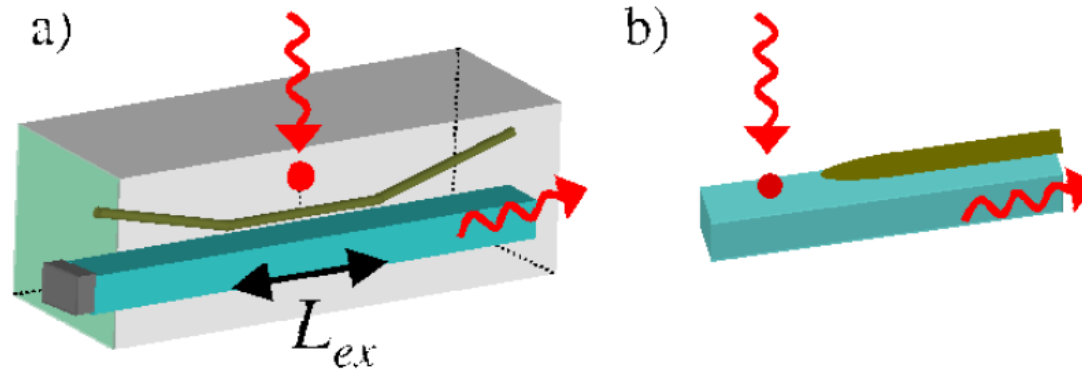
**feature
article**

Similarities to circuit QED



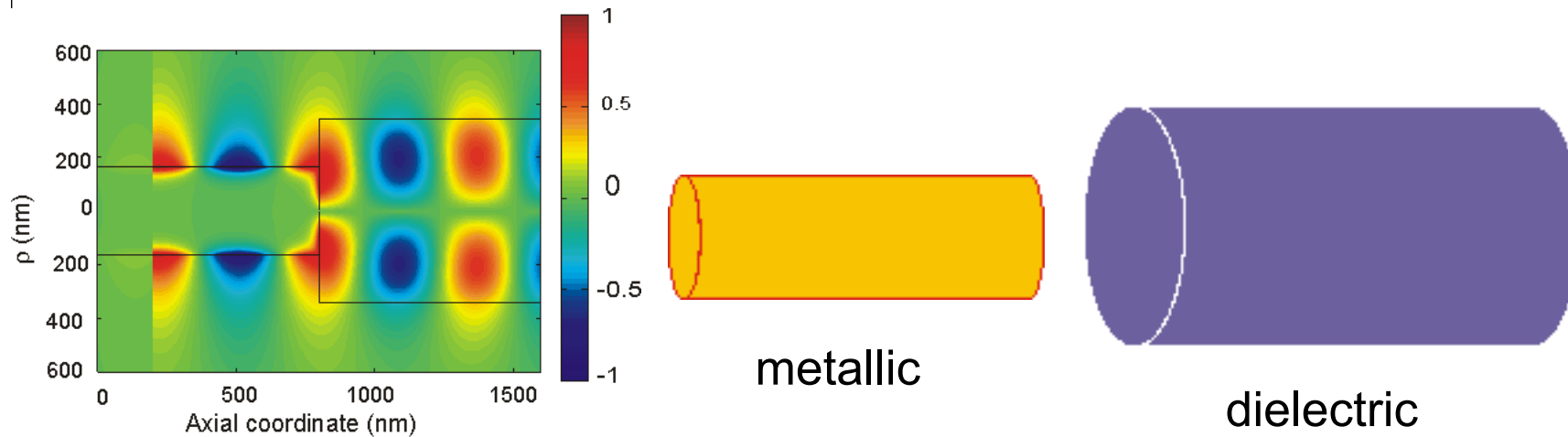
Quantum Optics with Surface Plasmons

D. E. Chang,¹ A. S. Sørensen,² P. R. Hemmer,^{1,3} and M. D. Lukin¹



>95% coupling of plasmonic and dielectric channels

broadband

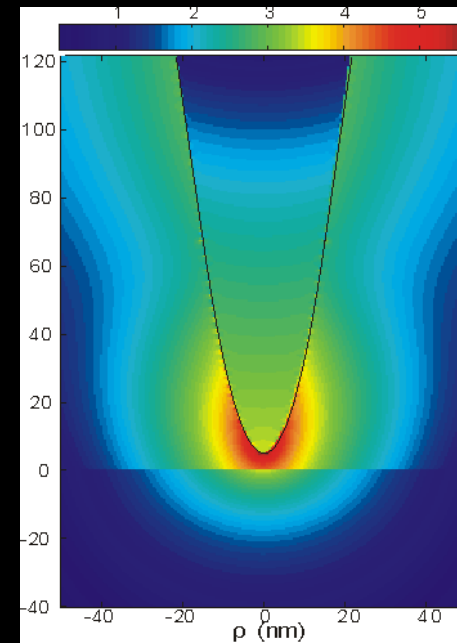
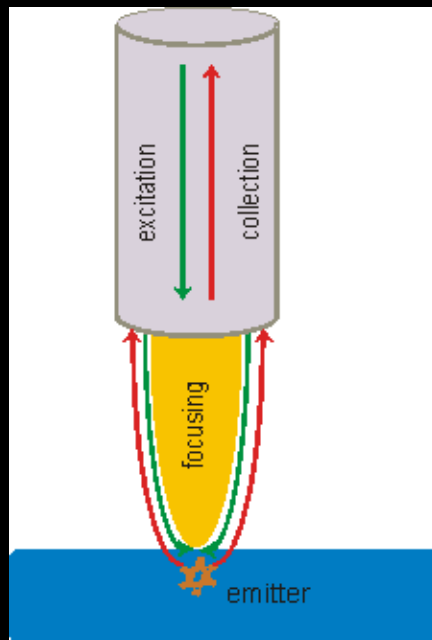


X. Chen, V. Sandoghdar, M. Agio, *Nano Lett.*, **9**, 3756 (2009).

... and super efficient couplig to a single emitter

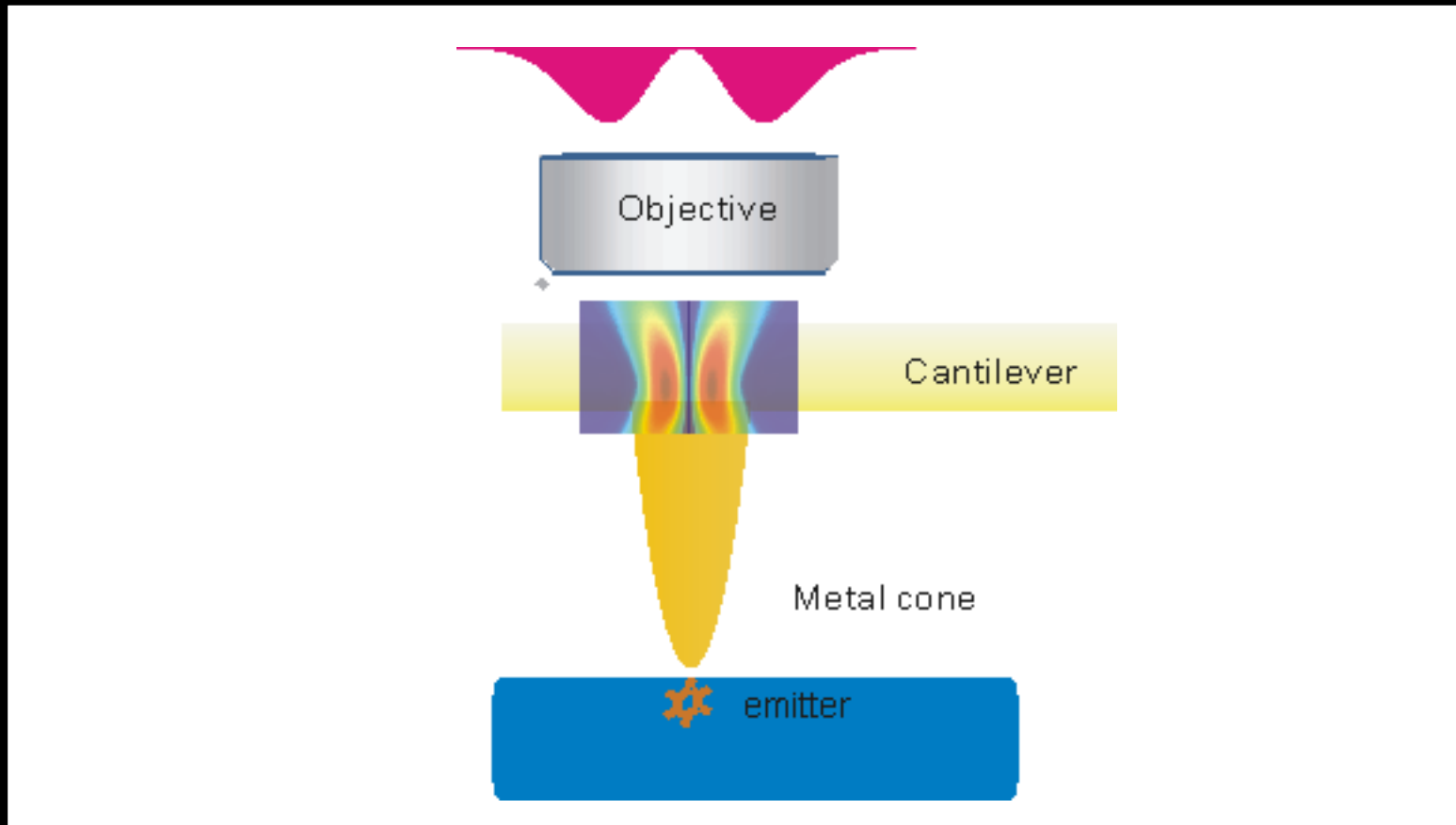


X.-W. Chen, V. Sandoghdar, M. Agio, Opt. Express, Focus issue: Unconventional Polarization States of Light, 18, 10878 (2010)



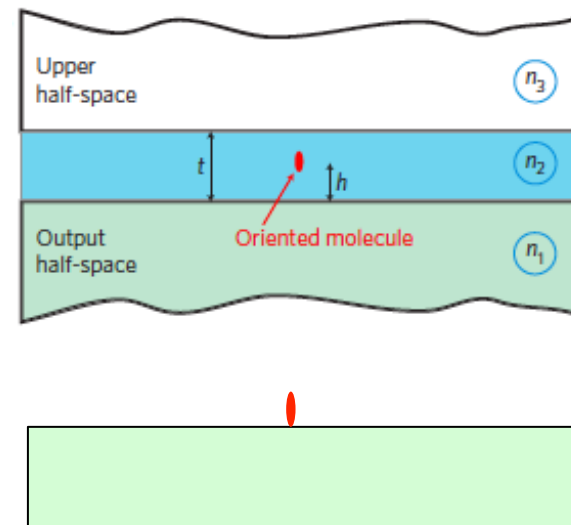
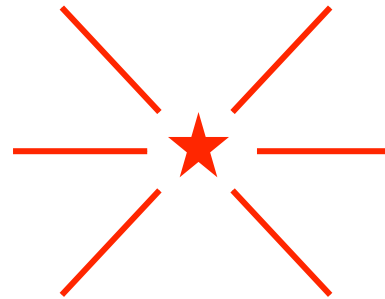
M.I. Stockman, Phys. Rev. Lett. 93, 137404 (2004), *and many others*

High-throughput high resolution SNOM



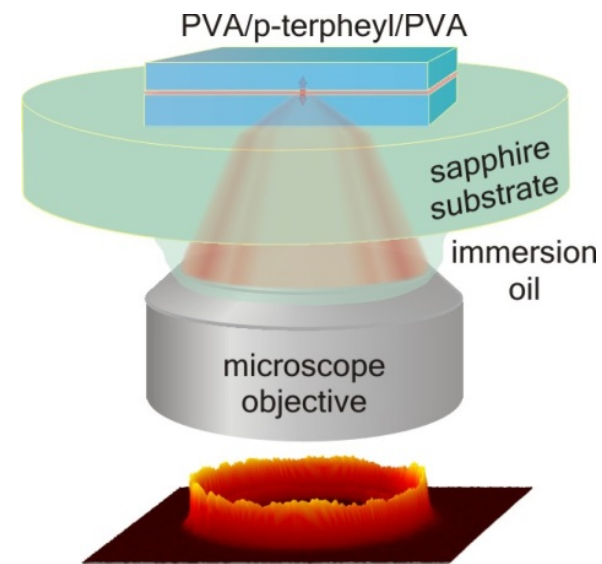
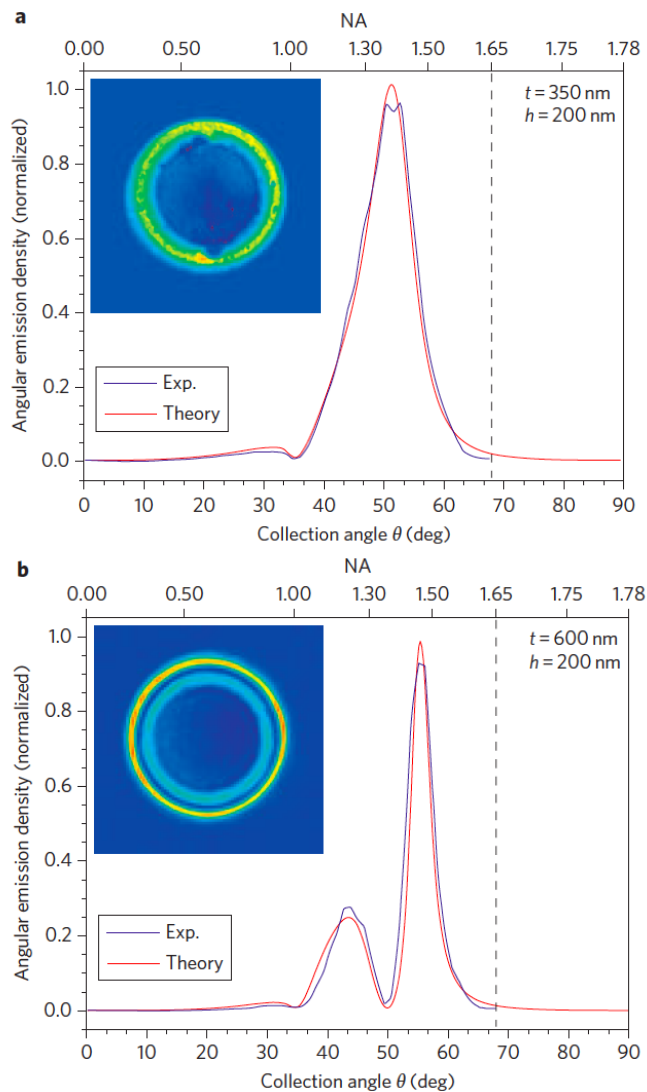
X.-W. Chen, V. Sandoghdar, M. Agio, *Opt. Express*, **18**, 10878 (2010);
Focus issue: Unconventional Polarization States of Light.

A planar dielectric antenna: 96% collection efficiency of single photons

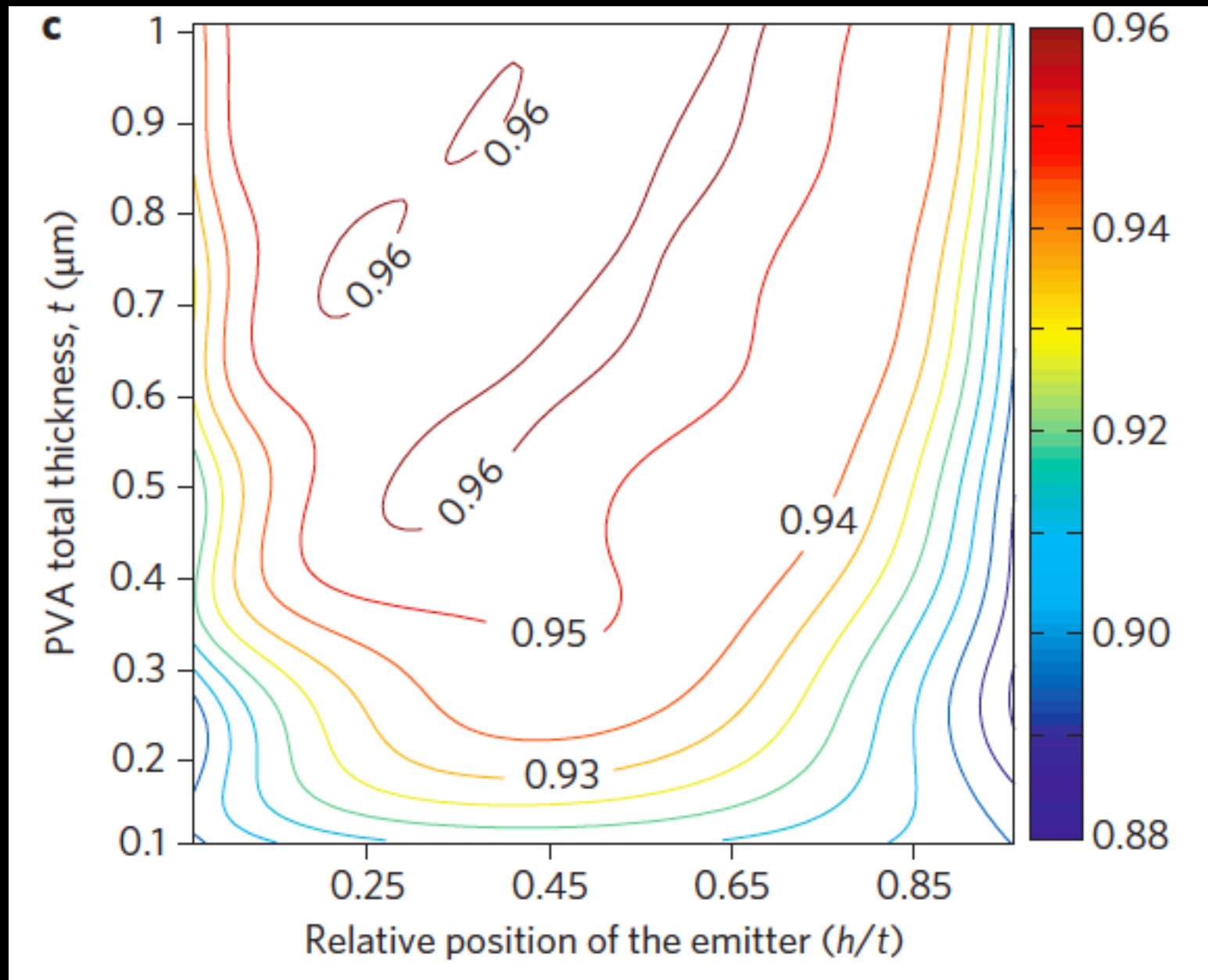


K-G. Lee, X-W. Chen, H. Eghlidi, P. Kukura, R. Lettow, A. Renn, V. Sandoghdar, S. Götzinger, *Nature Photonics*, **5**, 166 (2011).

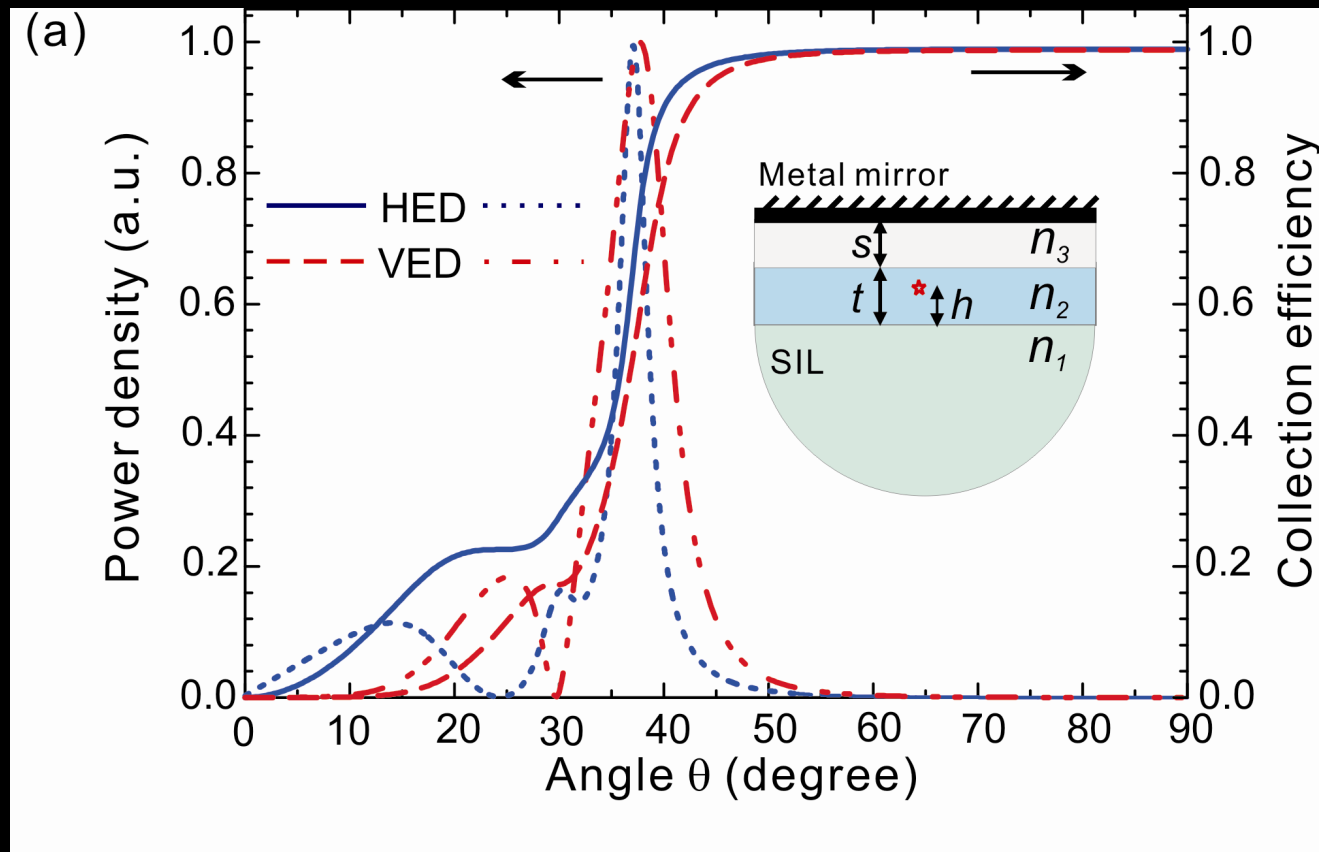
96% collection efficiency of single photons experimental results



K-G. Lee, X-W. Chen, H. Eghlidi, P. Kukura, R. Lettow, A. Renn, V. Sandoghdar, S. Göttinger, *Nature Photonics*, **5**, 166 (2011).



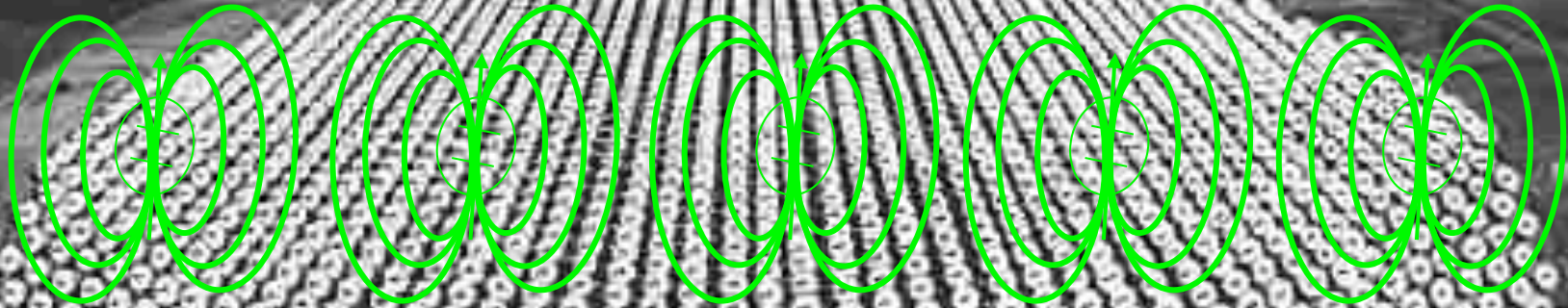
A metallo-dielectric antenna: 99% collection efficiency



X-W. Chen, S. Göttinger, V. Sandoghdar, *Opt. Lett.* **36**, 3545 (2011).

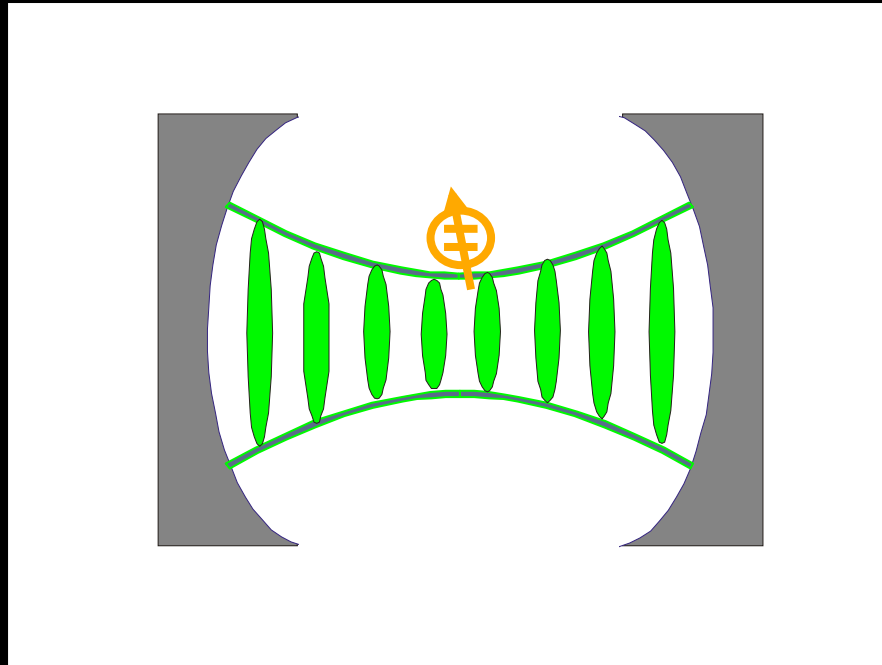
Connection to the efforts in atomic physics,
where collective effects are used

Spatial cooperative effect:
antenna theory



Connection to cavity QED

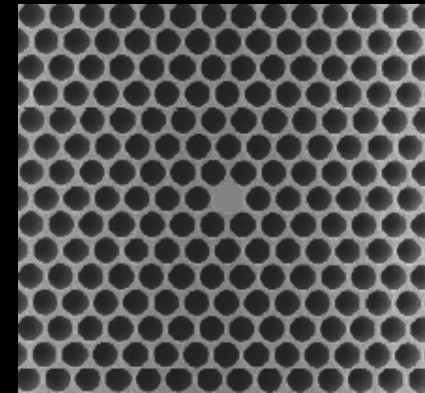
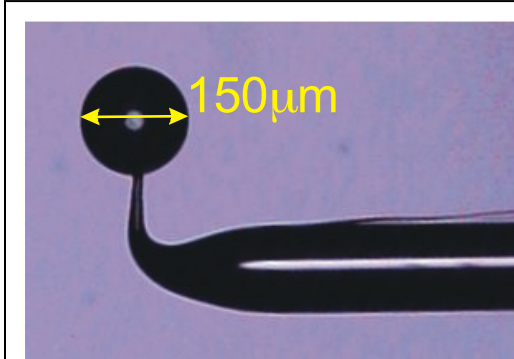
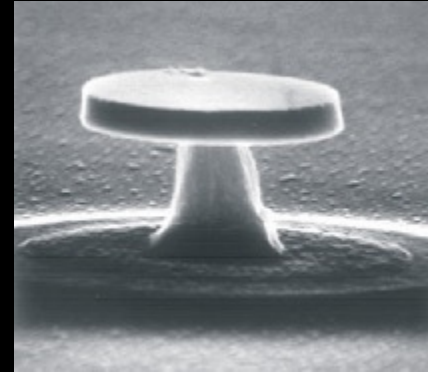
Confining light in a microcavity



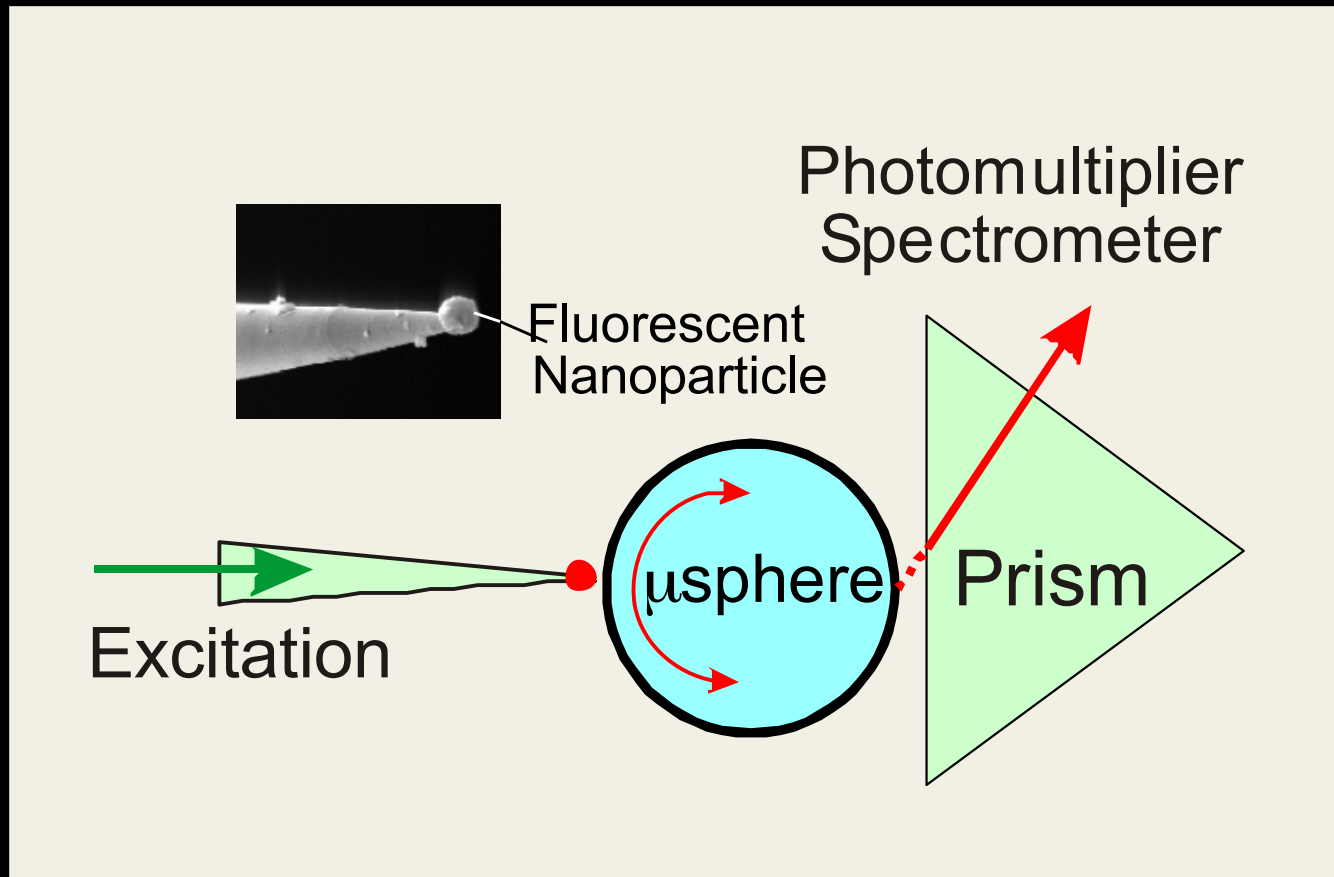
Multiple reflection \rightarrow interference \rightarrow resonance

Long but finite storage time, Q

Cavity finesse: mode volume, density of states



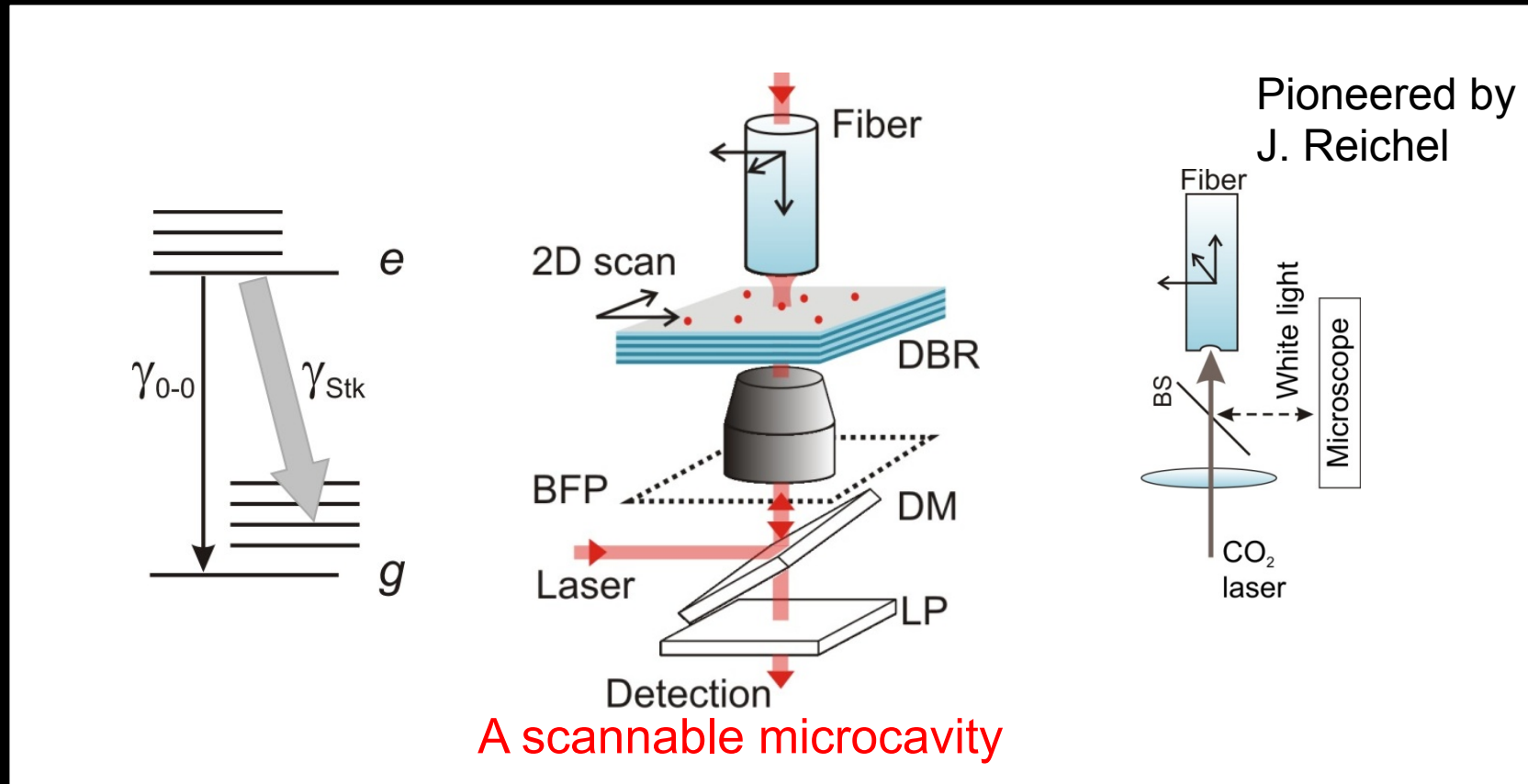
Coupling a nanoscopic emitter to the microsphere



S. Göttinger, L. de S. Menezes, A. Mazzei, S. Kühn, V. Sandoghdar & O. Benson, *Nano Letters* **6**, 1151 (2006).

S. Göttinger, L. de S. Menezes, O. Benson, D. V. Talapin, N. Gaponik, H. Weller, A. L. Rogach, V. Sandoghdar, *J. Opt. B* **6**, 154 (2004).

Enhancing the narrow-band emission Eliminating the broad Stokes-shifted photons

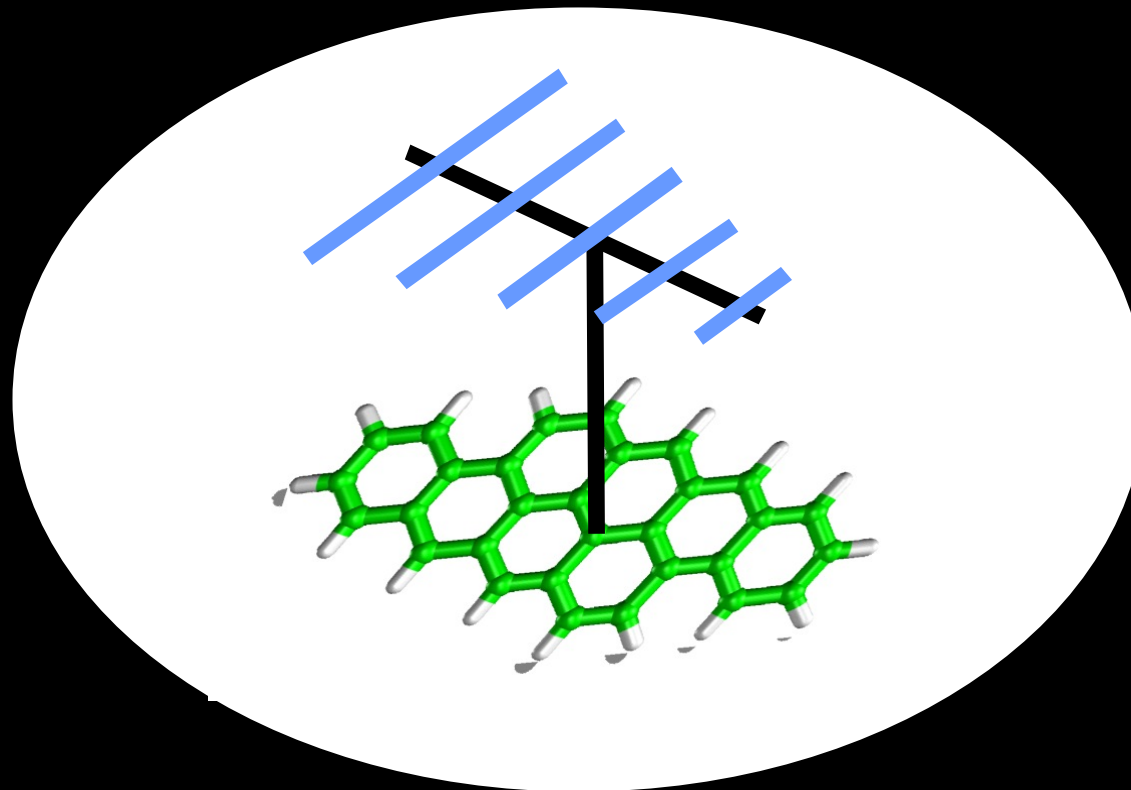


C. Toninelli, Y. Delley, T. Stöferle, A. Renn, S. Götzinger, V. Sandoghdar,
Appl. Phys. Lett. **97**, 021107 (2010).

Engineering the excitation and emission of an emitter by using *nano-antennas*

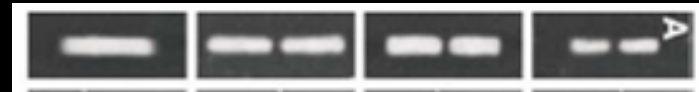
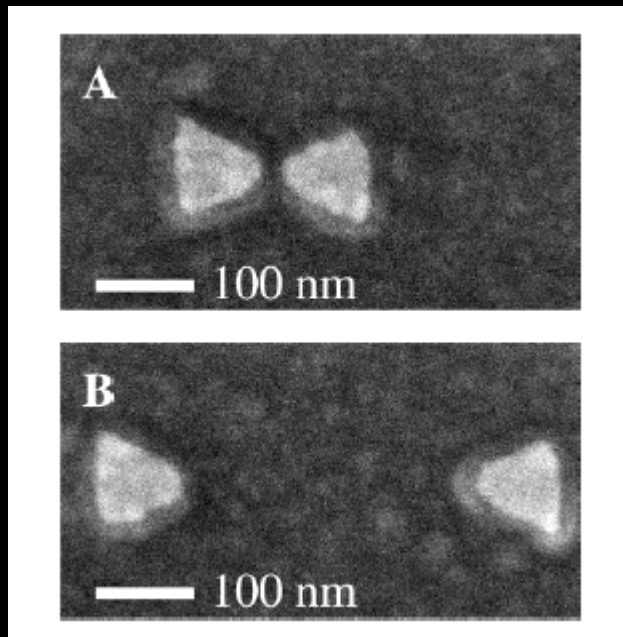
Cavity QED concepts have been already implemented to modify spontaneous emission and to reach strong coupling between an atom and a photon

Goal: achieve the same in the near field (without any quality factor!)



Optical nanoantennas

Inspired by radio antennae



Mühlschlegel, et al, Science (2005)

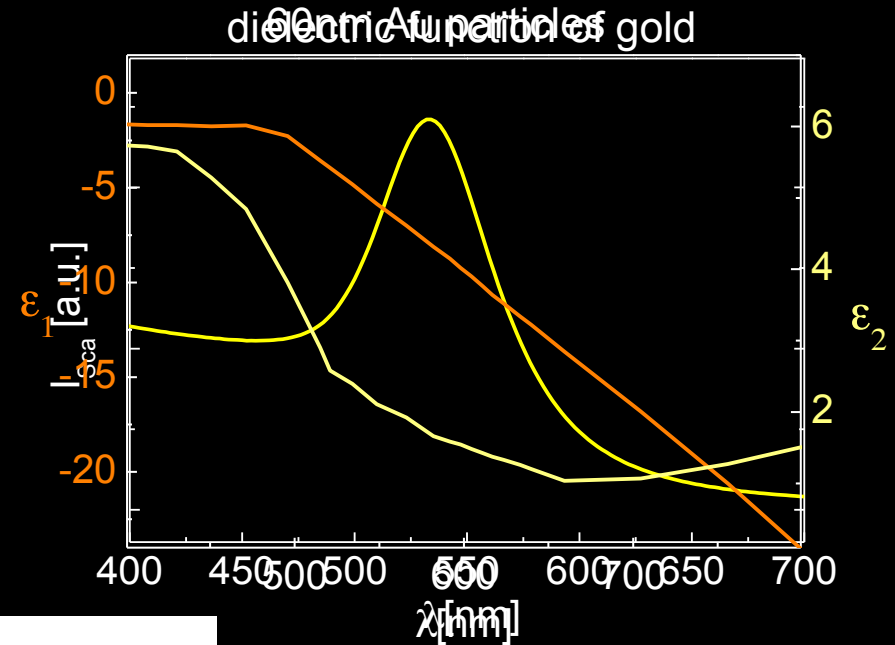
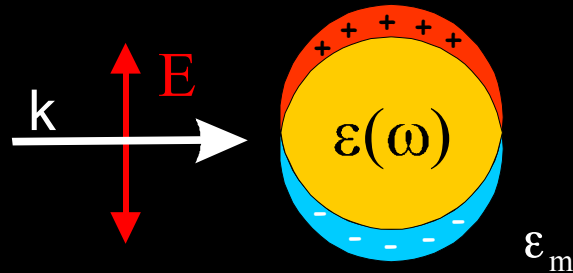
Fromm, et al, NanoLett (2004)

Antennas in space

Royal antennas



Optical Properties of Small Metal Particles



$$s(\lambda) = \eta\alpha(\lambda) = \eta\epsilon_{\text{med}}(\lambda) \frac{\pi D^3}{2} \frac{\epsilon_{\text{part}}(\lambda) - \epsilon_{\text{med}}(\lambda)}{\epsilon_{\text{part}}(\lambda) + 2\epsilon_{\text{med}}(\lambda)}$$

for $\epsilon(\omega) = -2\epsilon_m$

→ Resonance

full electrodynamic
calculation: Mie theory

G. Mie, *Ann. Phys.* **25** (1908)

The color of the particle depends on its material, size and shape

The First Nanotechnologists

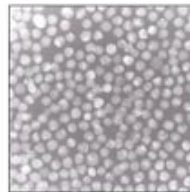
Ancient stained-glass makers knew that by putting varying, tiny amounts of gold and silver in the glass, they could produce the red and yellow found in stained-glass windows. Similarly, today's scientists and engineers have found that it takes only small amounts of a nanoparticle, precisely placed, to change a material's physical properties.

Gold particles in glass

Size*: 25 nm
Shape: sphere
Color reflected:

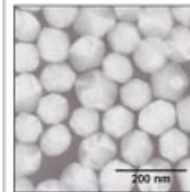


100 nanometers = 0.0001 millimeter



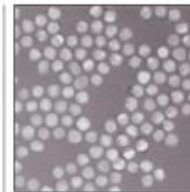
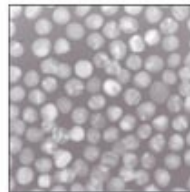
Silver particles in glass

Size*: 100 nm
Shape: sphere
Color reflected:



Had medieval artists been able to control the size and shape of the nanoparticles, they would have been able to use the two metals to produce other colors. Examples:

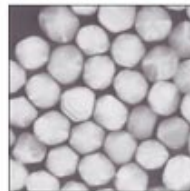
Size*: 50 nm
Shape: sphere
Color reflected:



Size*: 40 nm
Shape: sphere
Color reflected:



Size*: 100 nm
Shape: sphere
Color reflected:



Size*: 100 nm
Shape: prism
Color reflected:

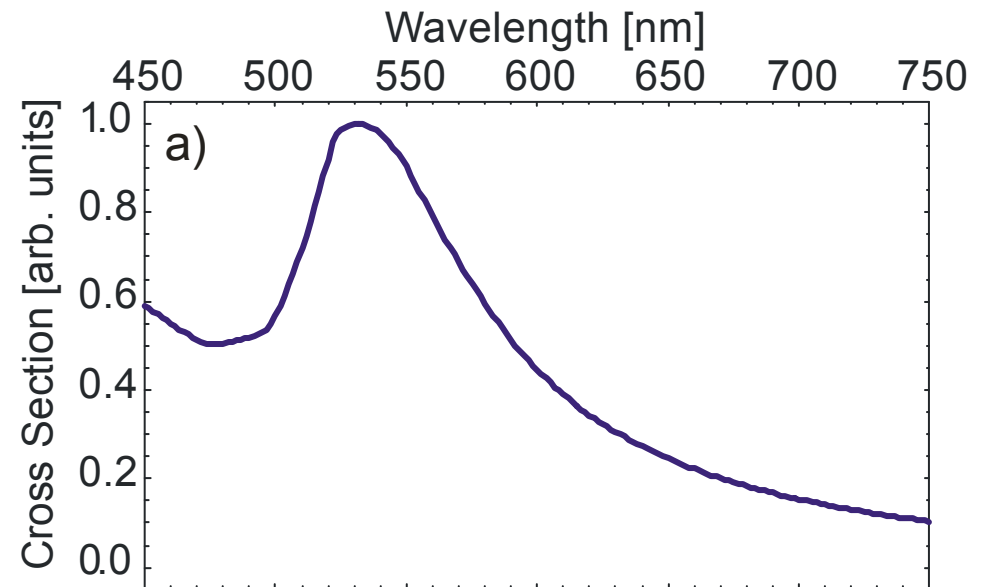
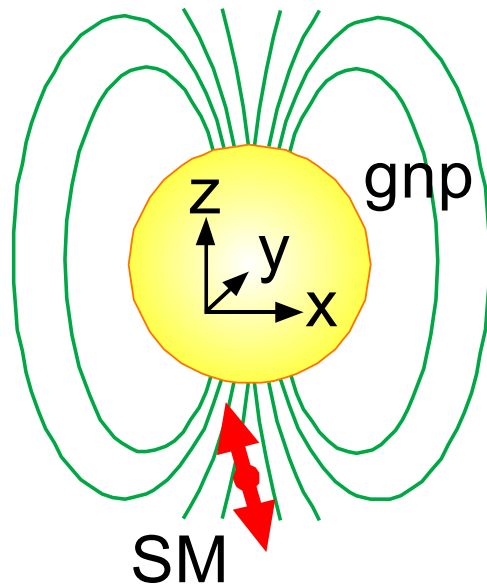


Source: Dr. Chad A. Mirkin, Institute of Nanotechnology, Northwestern University

*Approximate

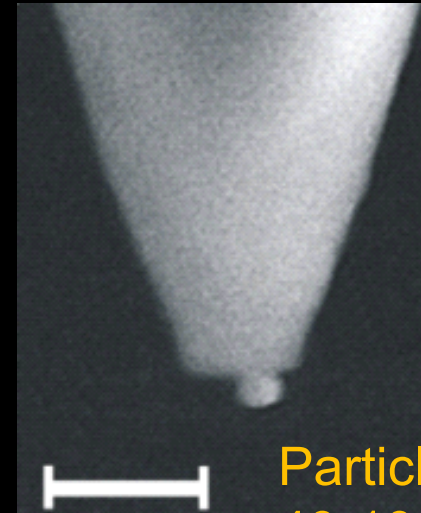
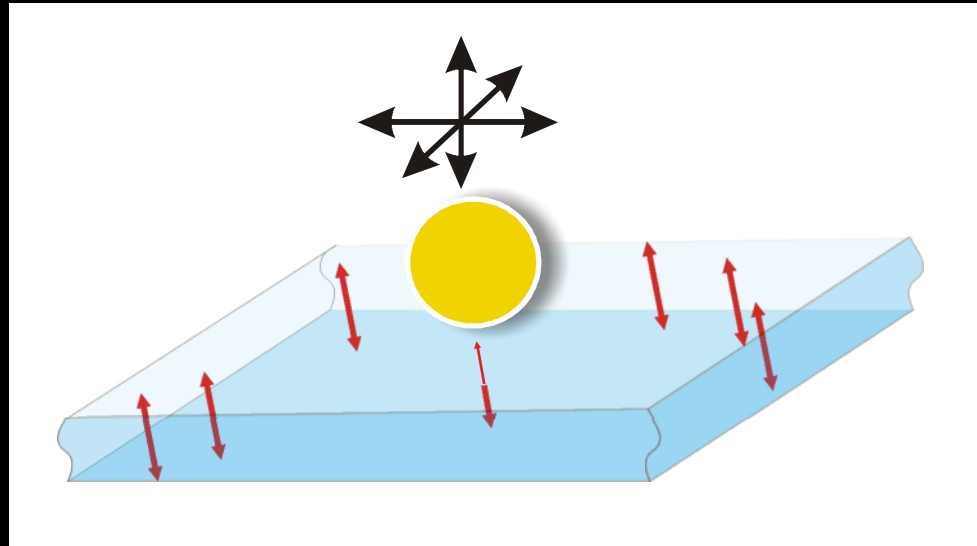


A single gold nanoparticle as a Nano-Antenna



The subwavelength gold particle acts as the extension/magnification of the molecular dipole moment
→ Faster emission

Controlled positioning of the nanoantenna



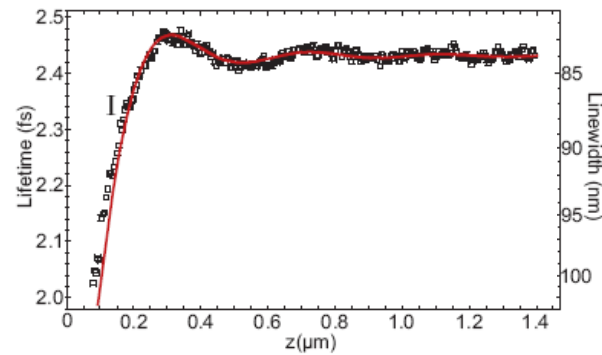
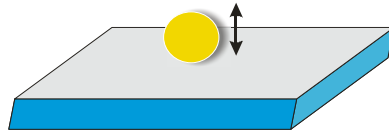
200 nm

Particle diameter
40-100 nm

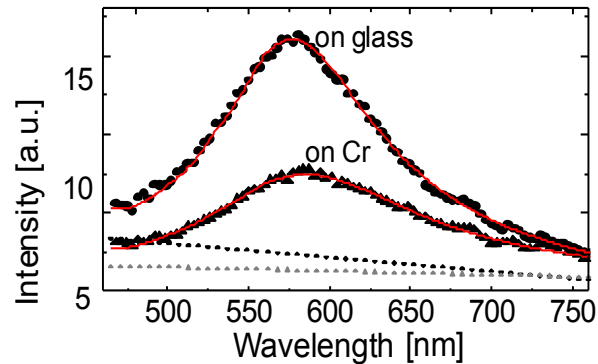
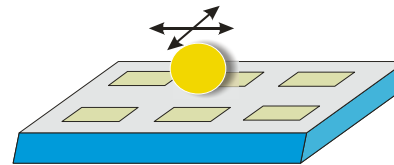
T. Kalkbrenner, M. Ramstein, J. Mlynek, V. Sandoghdar
J. Microscopy **202**, 72 (2001).

T. Kalkbrenner, U. Hakanson, & V. Sandoghdar,
Nano Lett. **4**, 2309 (2004).

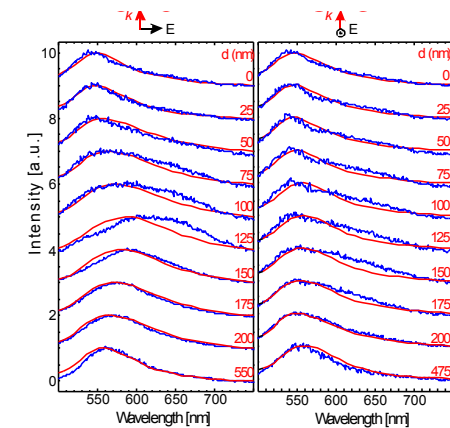
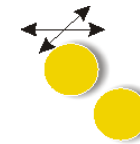
Local modification/control of antenna resonances



B. C. Buchler, et al,
Phys. Rev. Lett. **95**,
063003 (2005).



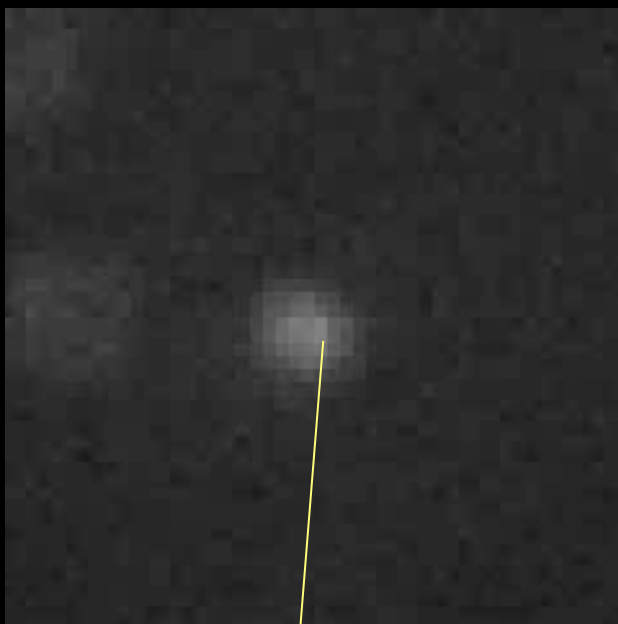
T. Kalkbrenner, et al,
Phys. Rev. Lett. **95**,
200801 (2005).



U. Hakanson, et al,
Phys. Rev. B **77**,
155408 (2008).

Controlled Interaction of a Single Gold Nanoparticle with a Single Molecule

Scanning the molecule

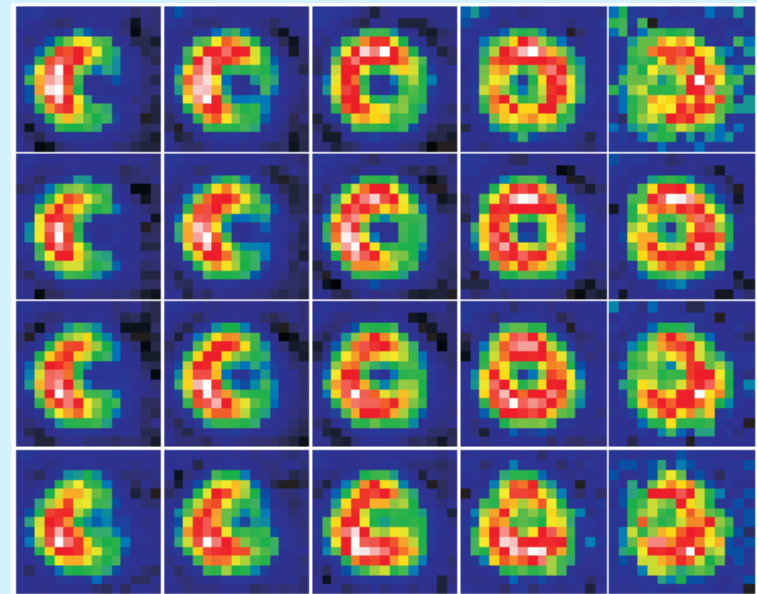
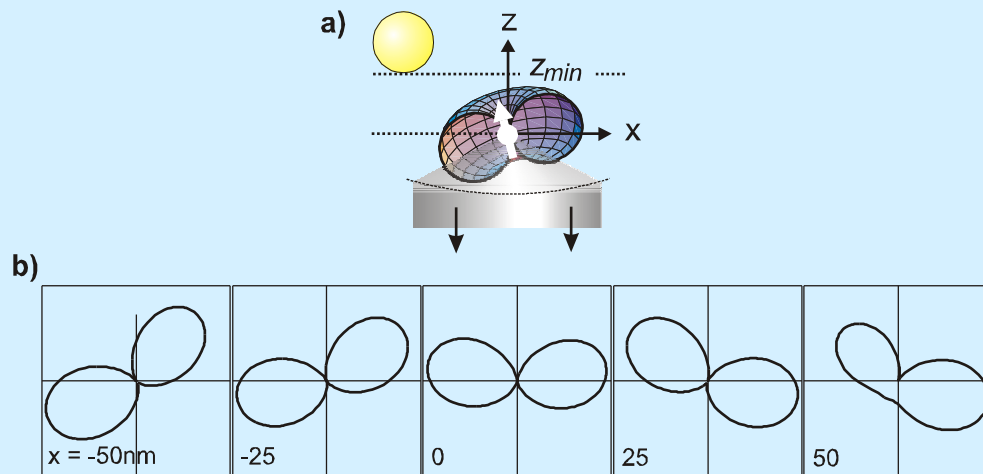


gold
nanoparticle

Scanning the particle



Modification of the molecular emission pattern



S. Kühn, et al, *PRL*. **97**, 017402 (2006).

S. Kühn, et al., *Mol. Phys.* **106**, 893 (2008).

enhancement of the excitation field
modification of the radiative lifetime
modification of the nonradiative lifetime

Overall collection efficiency

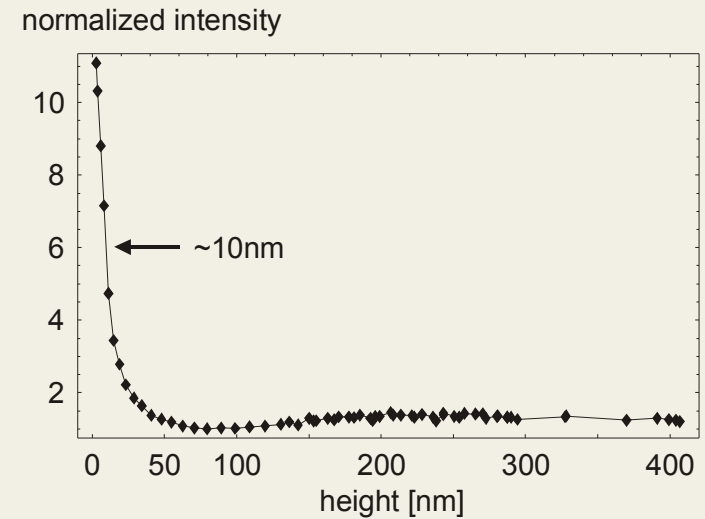
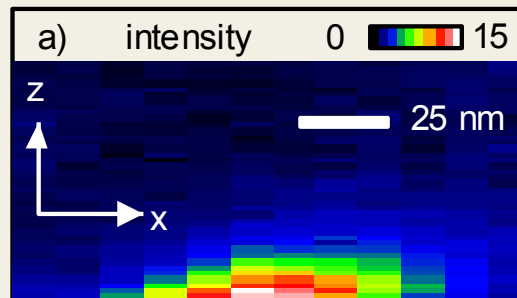
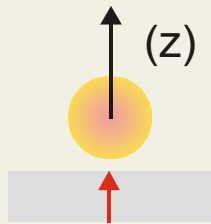
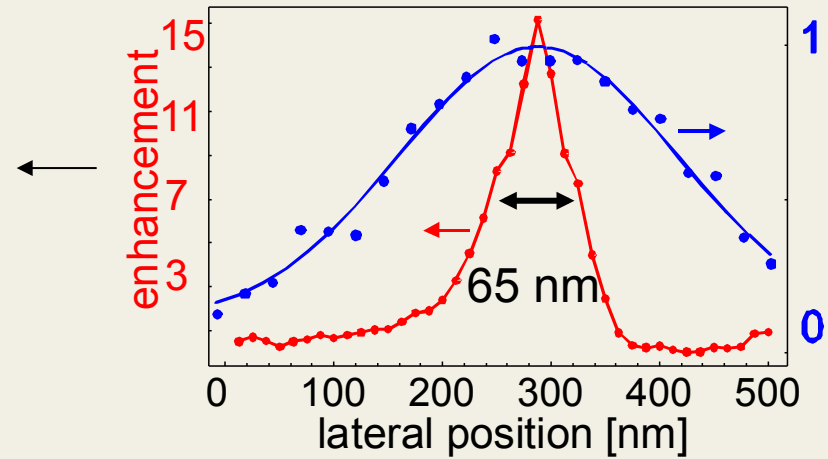
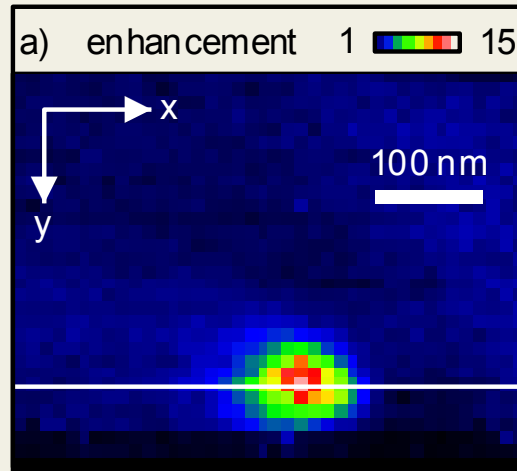
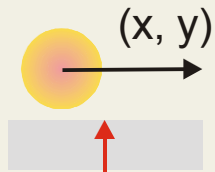
| Field enhancement |²

$$S_{fluor,enhanced} = \xi_{collection} \cdot K_{excitation} \cdot \eta_{system} \cdot S_{fluor,0}$$

Quantum efficiency
(of complete system)

$$\eta = \frac{\Gamma_{rad}}{\Gamma_{rad} + \Gamma_{nr}}$$

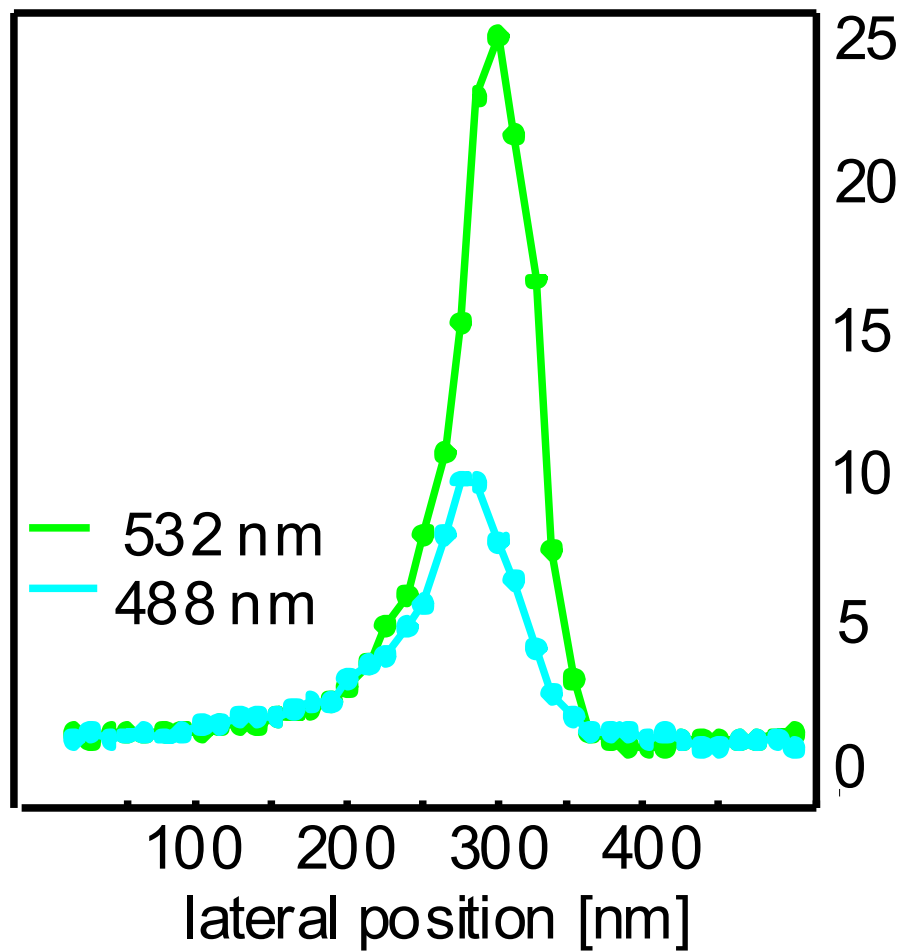
Influence on the Fluorescence Intensity



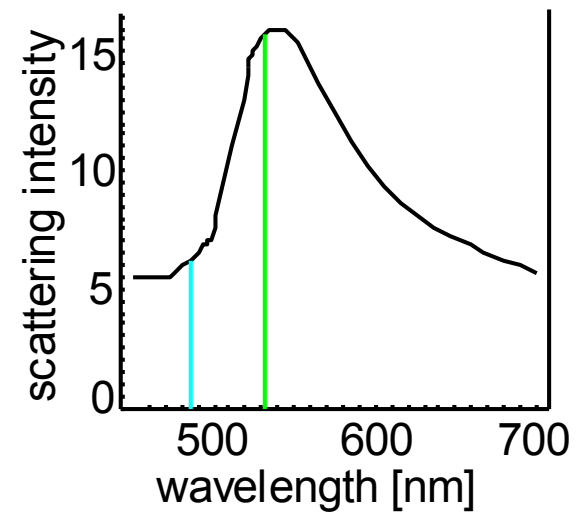
S. Kühn, U. Hakanson, L. Rogobete & V. Sandoghdar,
Phys. Rev. Lett. **97**, 017402 (2006) .

Direct demonstration of the antenna resonance effect

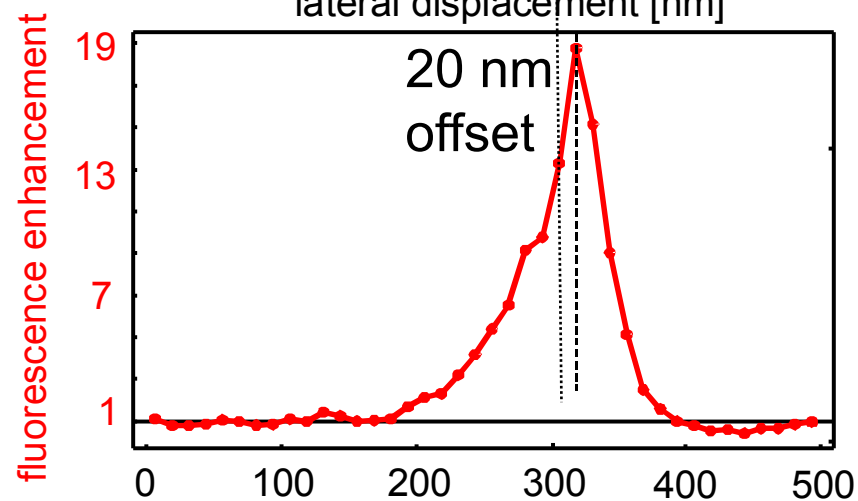
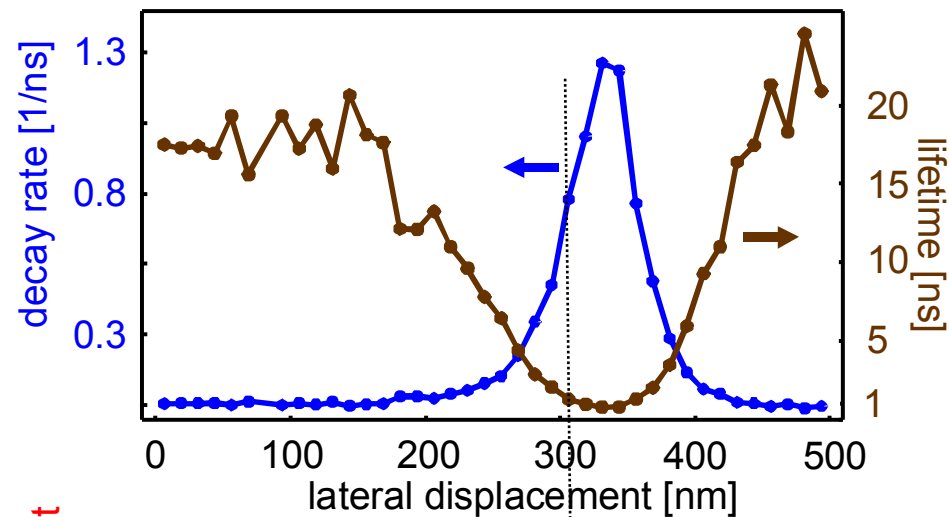
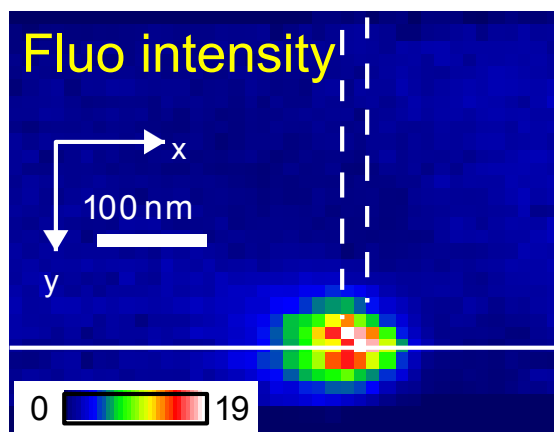
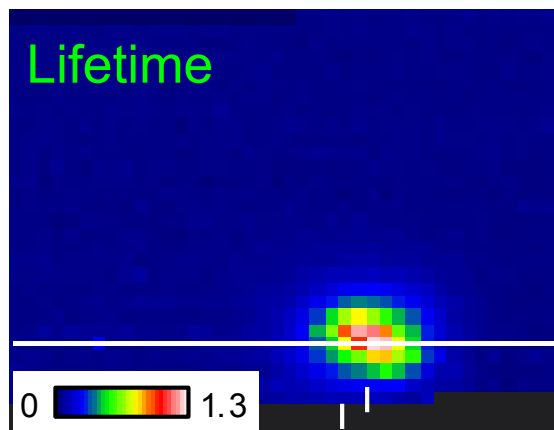
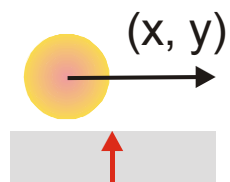
intensity enhancement



Plasmon resonance

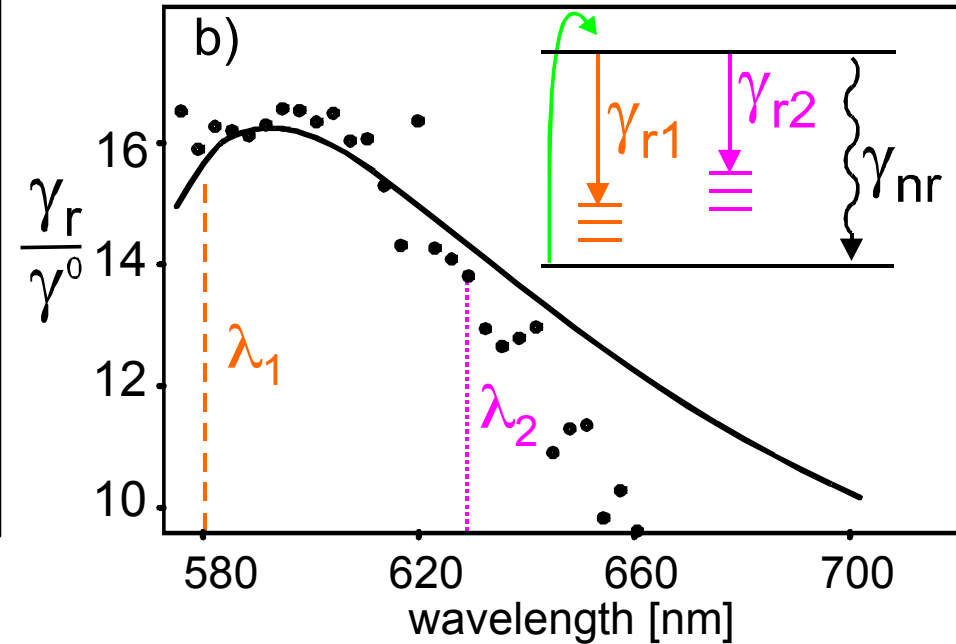
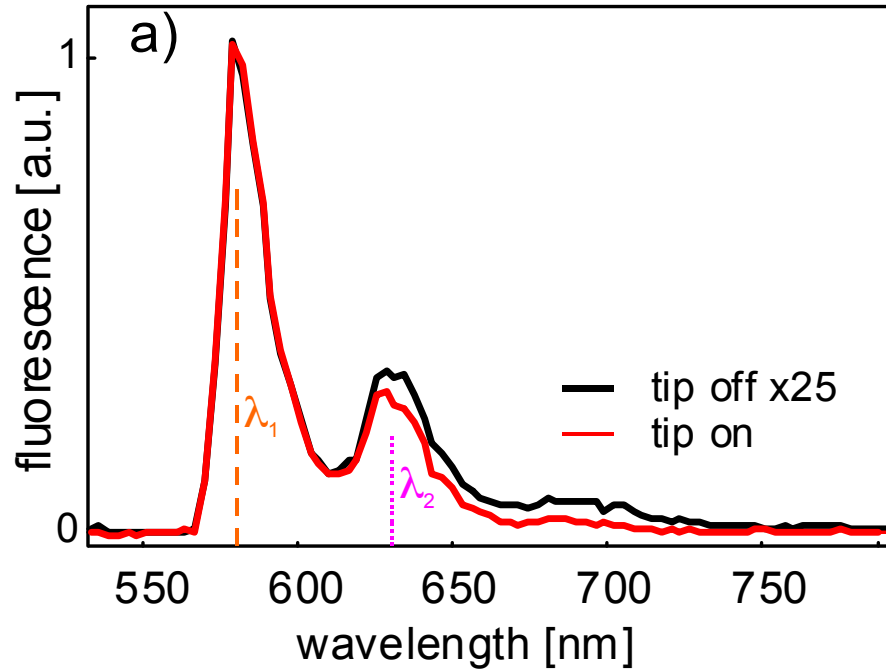


Influence on the Emission Lifetime



S. Kühn, U. Hakanson, L. Rogobete & V. Sandoghdar,
Phys. Rev. Lett. **97**, 017402 (2006) . AND the **Supplementary**

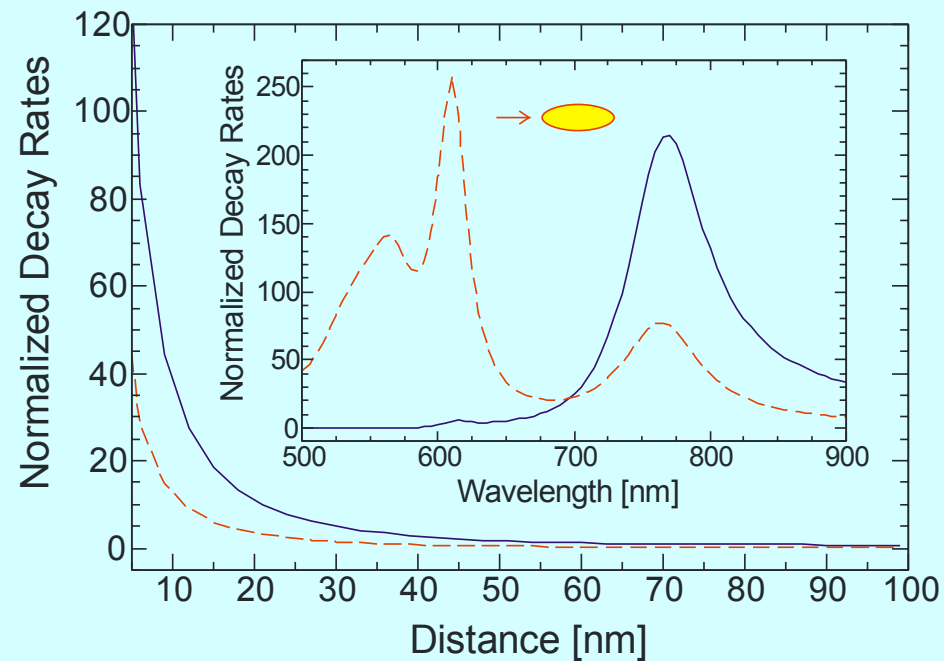
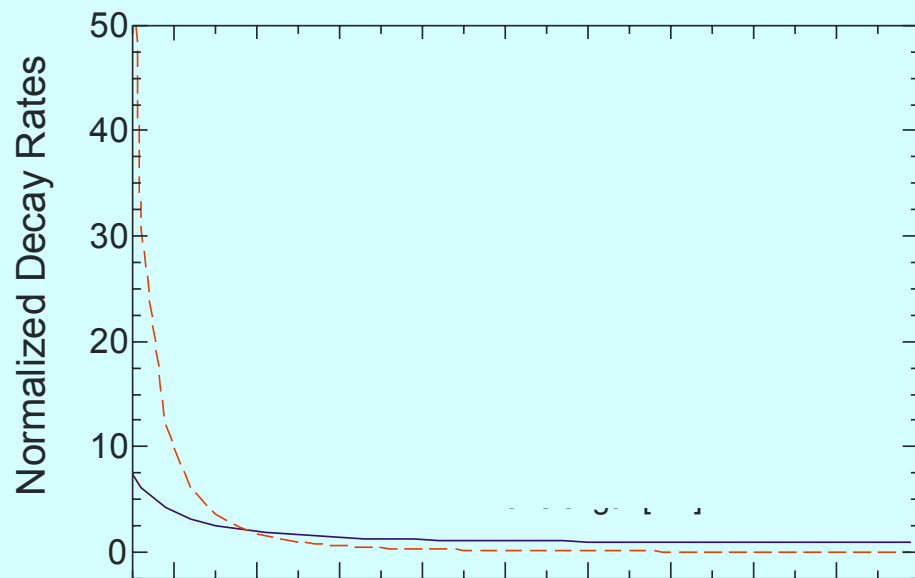
Modification of the emission spectrum



S. Kühn, U. Hakanson, L. Rogobete & V. Sandoghdar,
Phys. Rev. Lett. **97**, 017402 (2006) ; **Supplementary Material**

S. Kühn, et al., *Mol. Phys.*, **106**, 893 (2008).

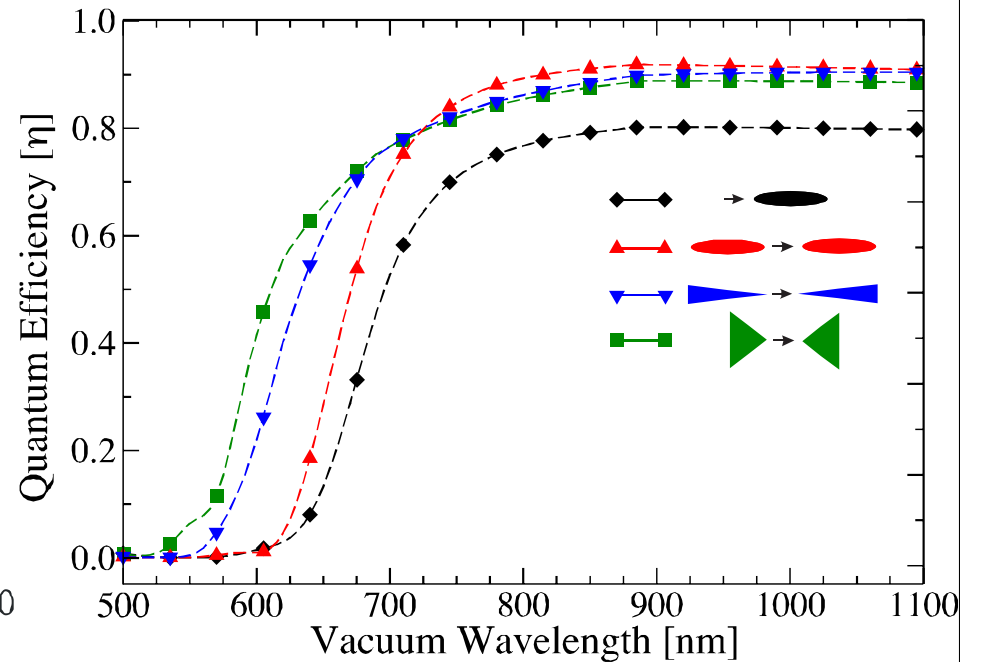
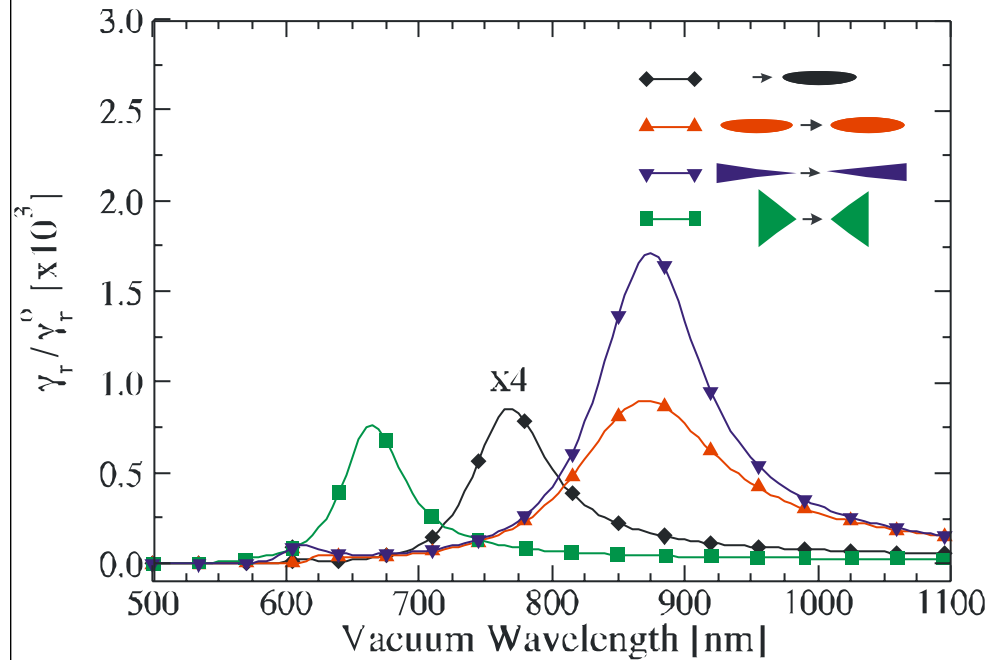
Engineering the quantum efficiency



L. Rogobete, F. Kaminski, M. Agio, V. Sandoghdar, *Opt. Lett.* **32**, 1623 (2007)

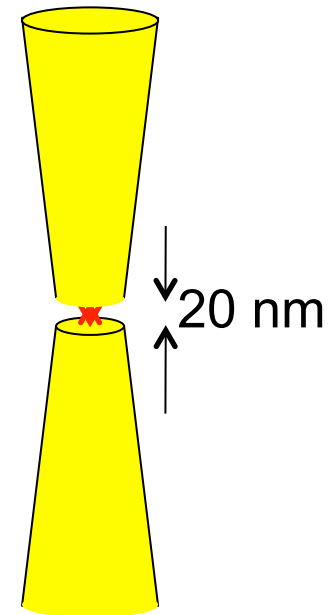
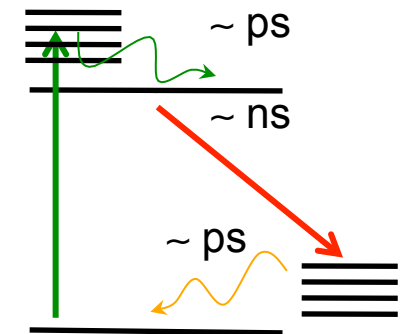
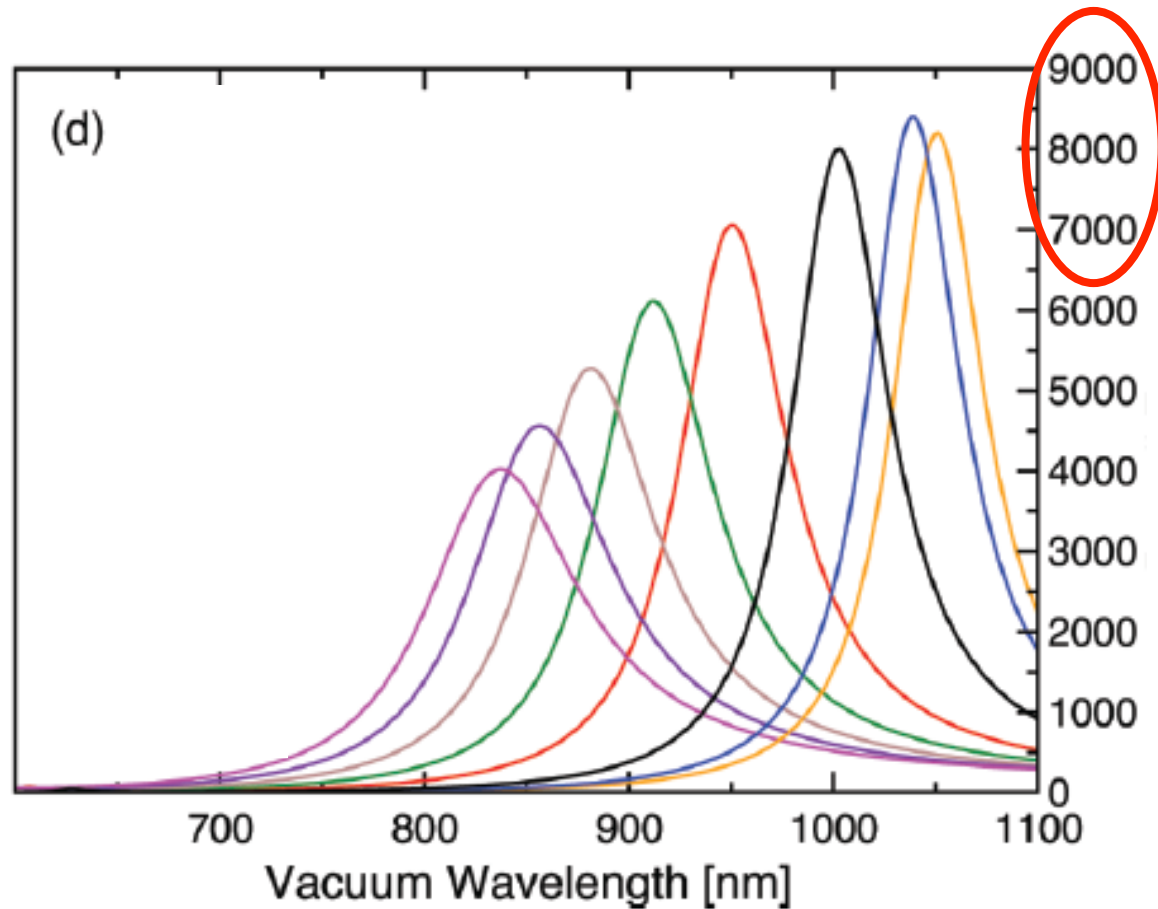
Optimizing the radiative emission rate

Key: mode matching



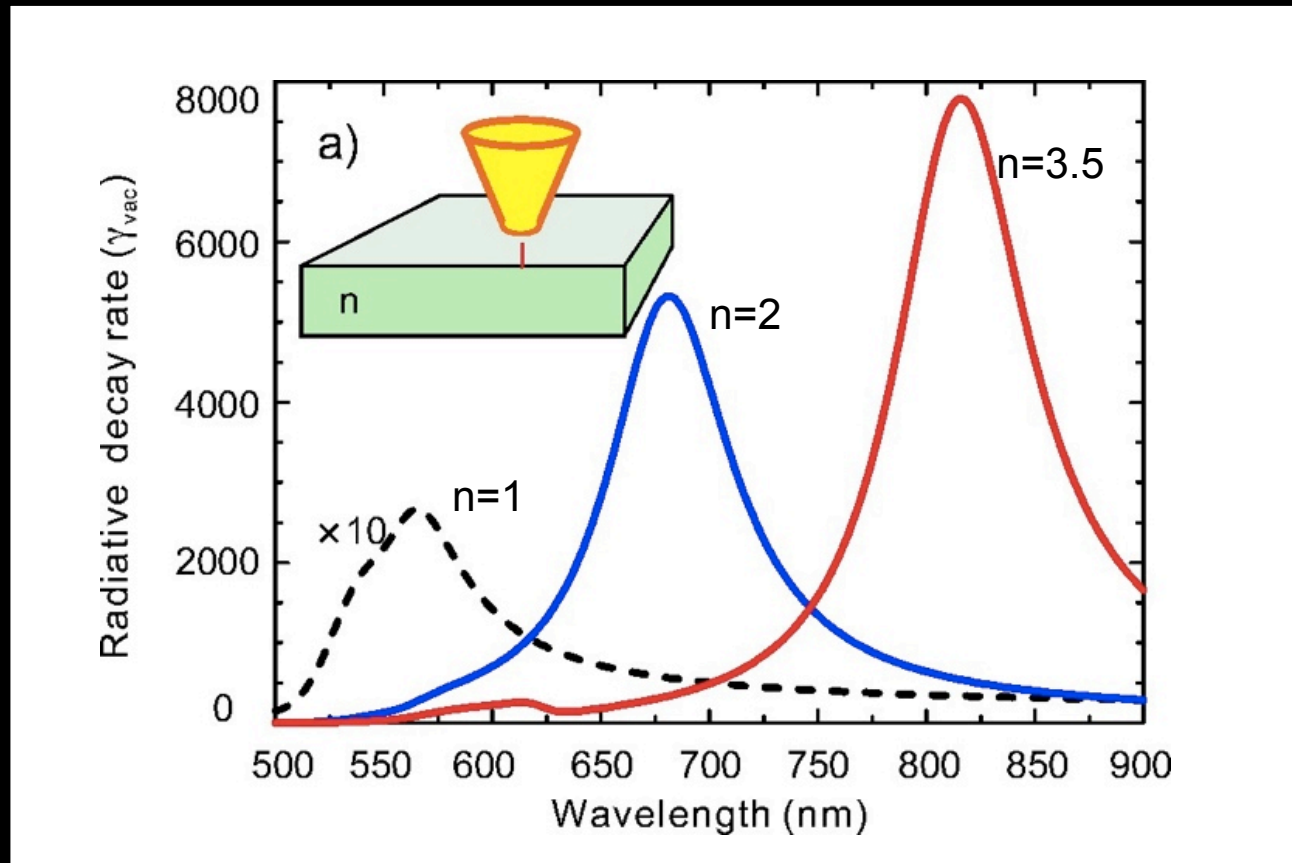
L. Rogobete, F. Kaminski, M. Agio, V. Sandoghdar,
Opt. Lett. **32**, 1623 (2007).

Antenna enhancement of spontaneous emission rate



A. Mohammadi, F. Kaminski, V. Sandoghdar, M. Agio,
J. Phys. Chem. C 114, 7372 (2010).

Ultrastrong enhancement of spontaneous emission



X-W. Chen, M. Agio, V. Sandoghdar, *Phys. Rev. Lett.* to appear (2012).

Room temperature experiments

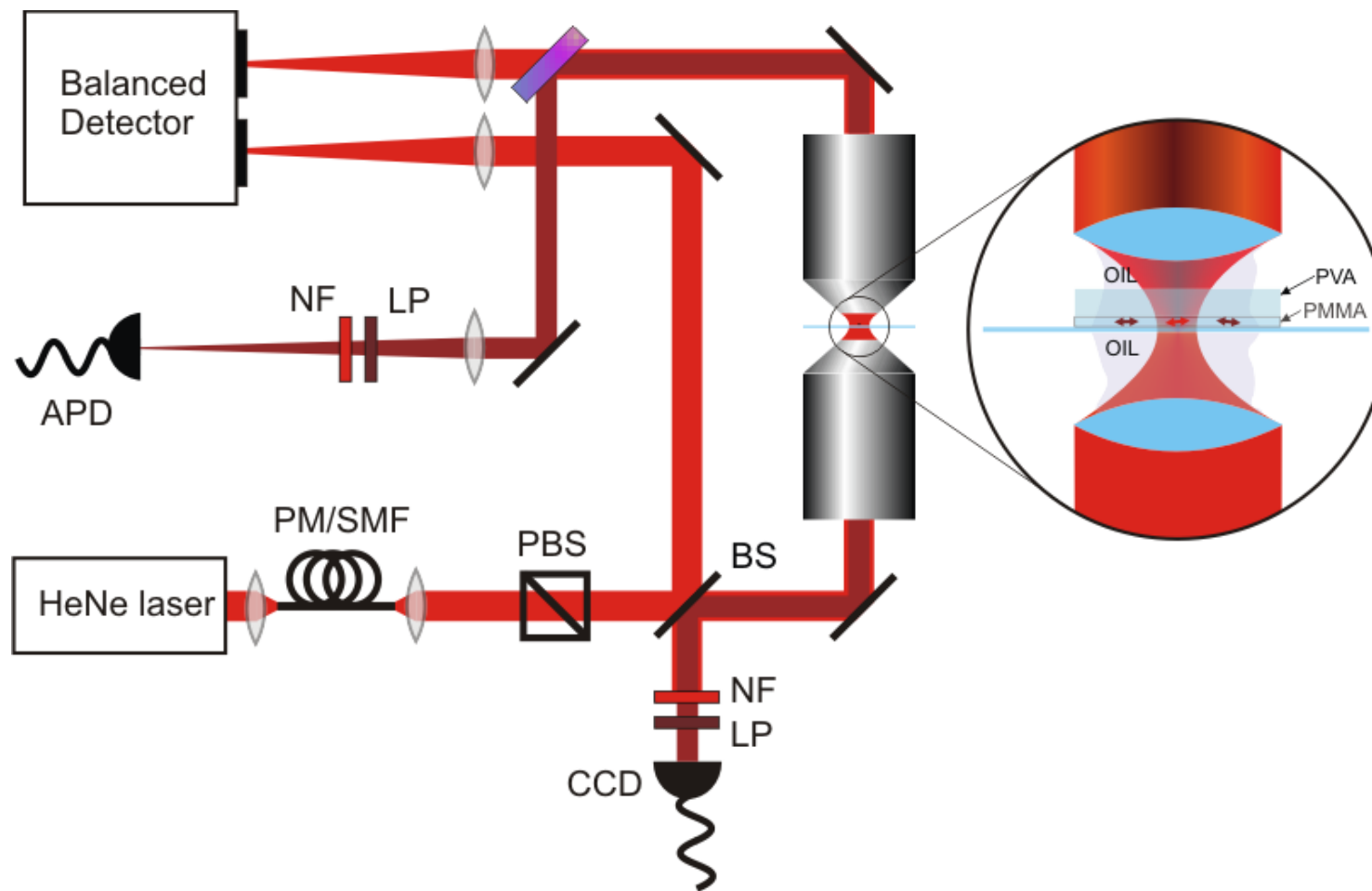
$$\sigma_{abs} = \frac{3\lambda^2}{2\pi} \left(\frac{\gamma_0}{\gamma} \right) 10^{-6}$$

Limit of A ~ (200 nm)²

Room-temperature in the solid state $\sigma \sim 10^{-15} \text{ cm}^2 \sim 0.1 \text{ (nm)}^2$

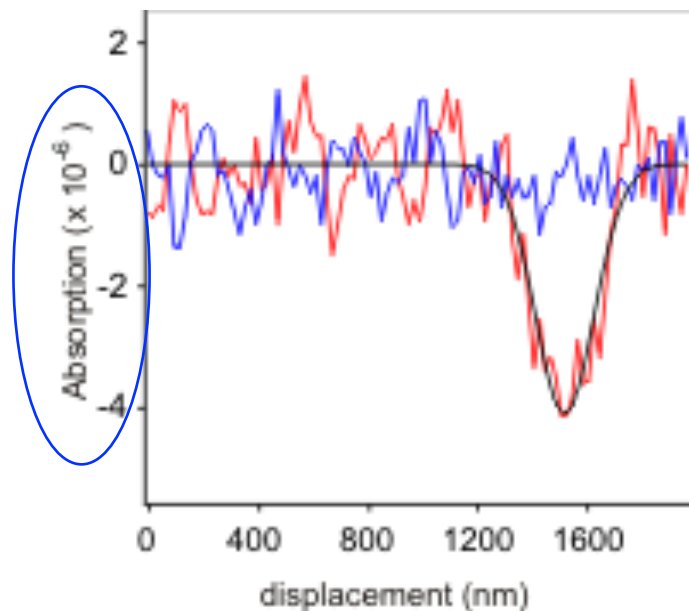
Need a signal-to-noise ratio of $10^{-6} \rightarrow$ laser noise suppression

Transmission measurement of *single molecules*



P. Kukura, M. Celebrano, A. Renn, V. Sandoghdar,
J. Phys. Chem. Lett. **1**, 3323 (2010).

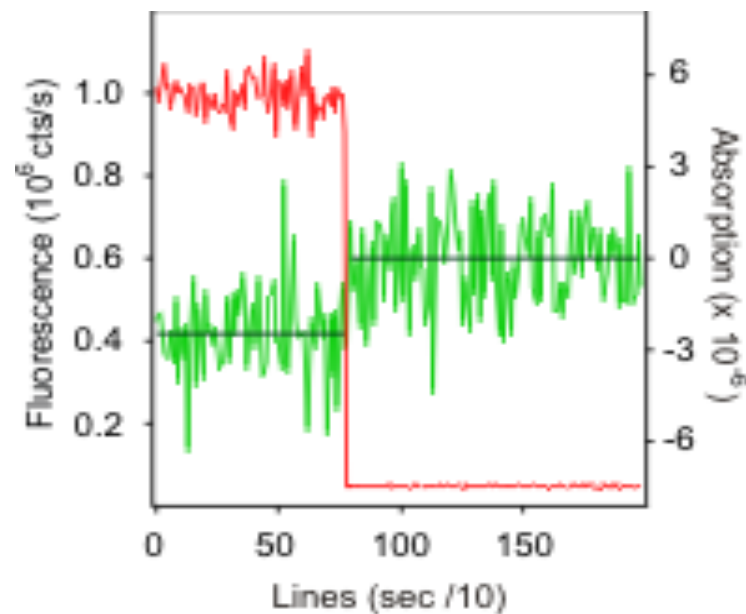
The first demonstration of single-molecule sensitivity in absorption @ room temperature



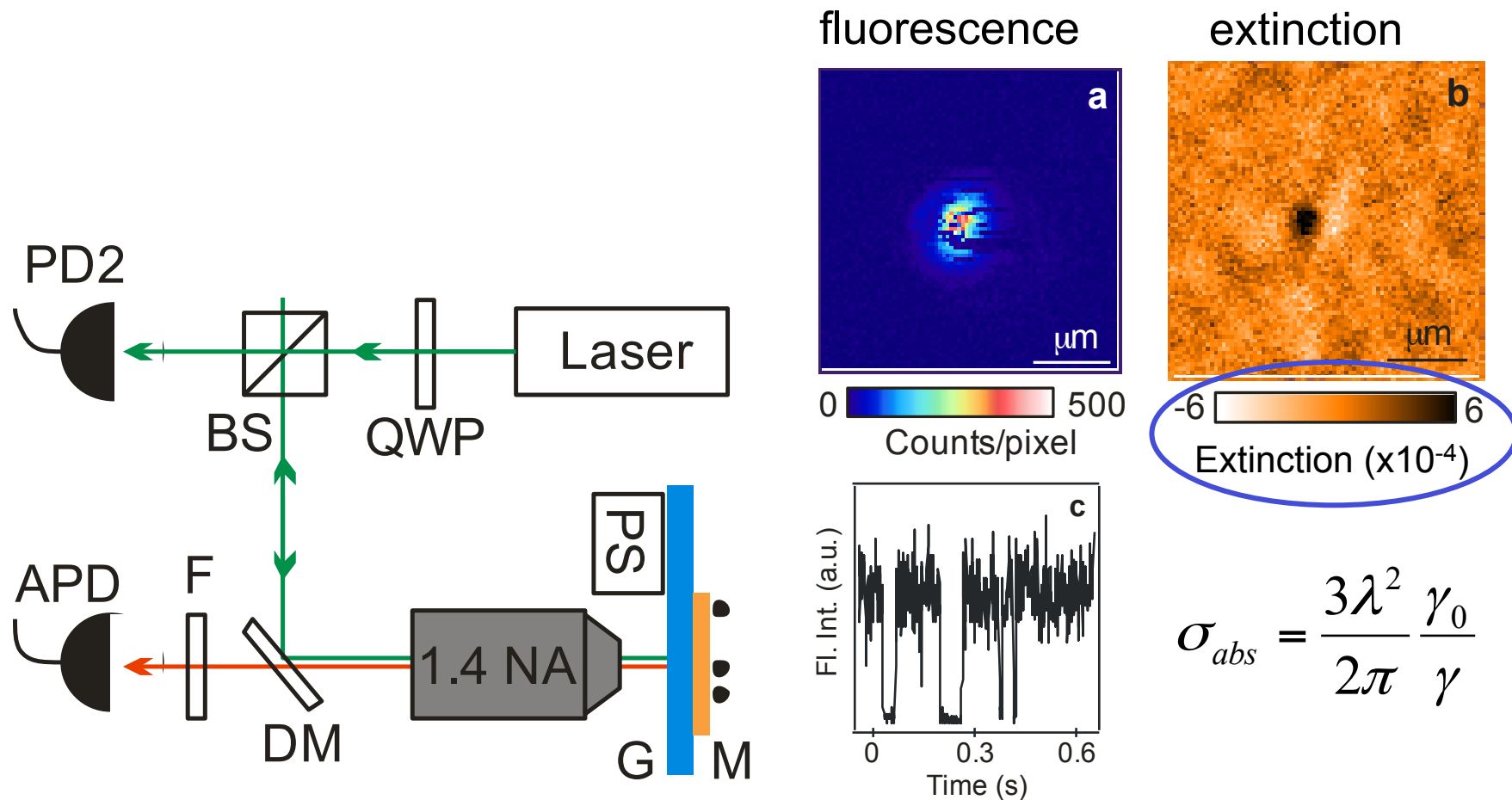
after bleaching

before bleaching

P. Kukura, M. Celebrano,
A. Renn, V. Sandoghdar,
J. Phys. Chem. Lett. **1**,
3323 (2010).



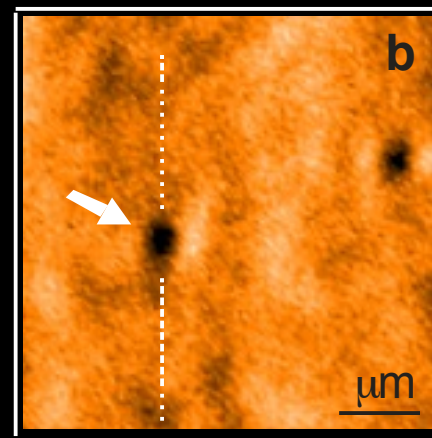
Room-Temperature Imaging of a Single (nonfluorescent) Quantum Dot



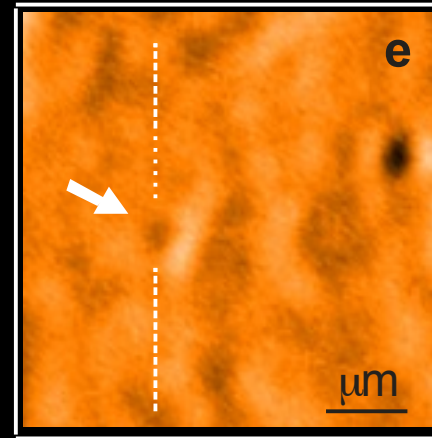
$$\sigma_{abs} = \frac{3\lambda^2}{2\pi} \frac{\gamma_0}{\gamma}$$

P. Kukura, M. Celebrano, A. Renn, V. Sandoghdar,
Nano Lett. 9, 926 (2009).

Main limitation: background scattering

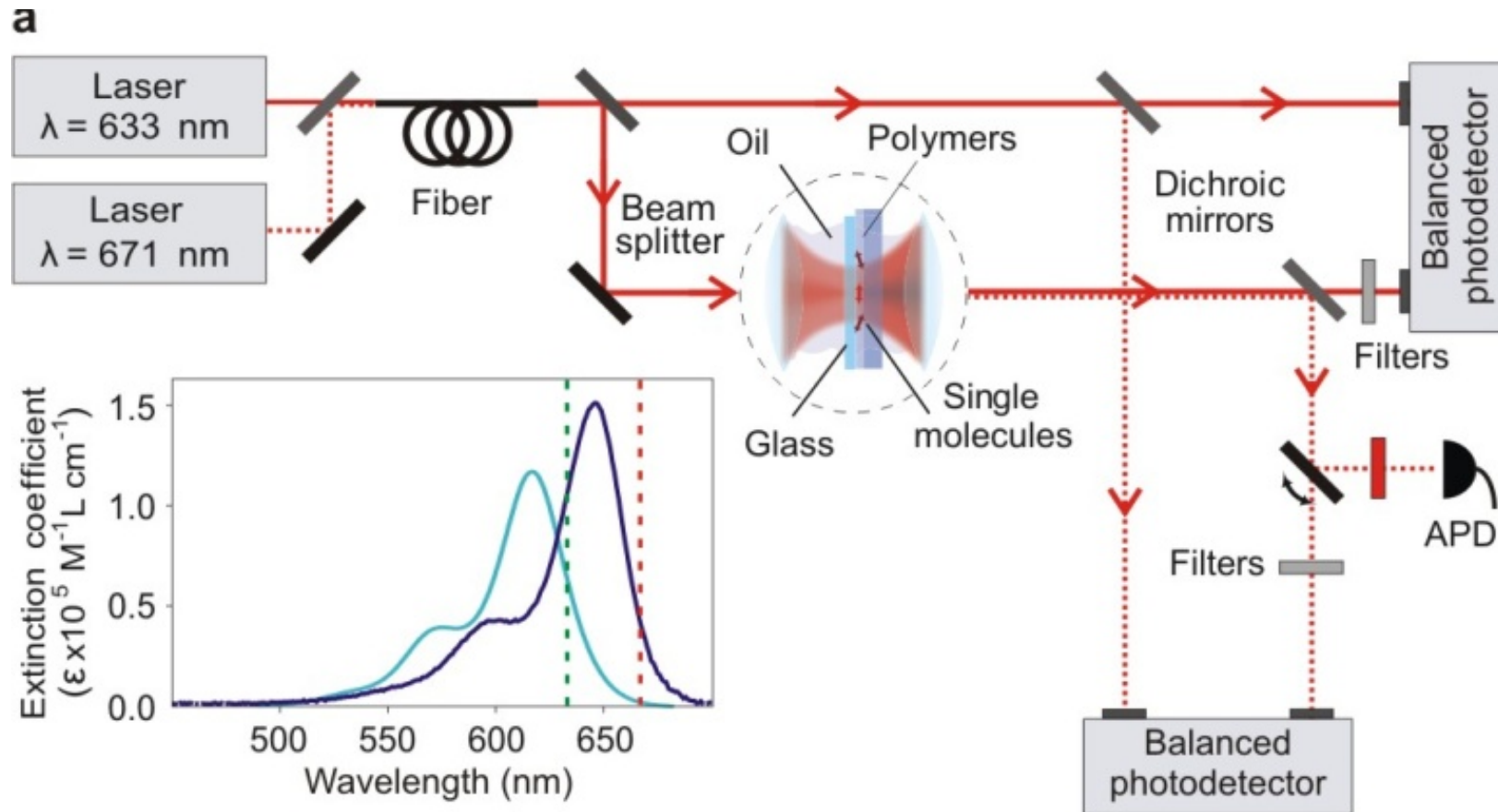


-6  6
Extinction ($\times 10^4$)



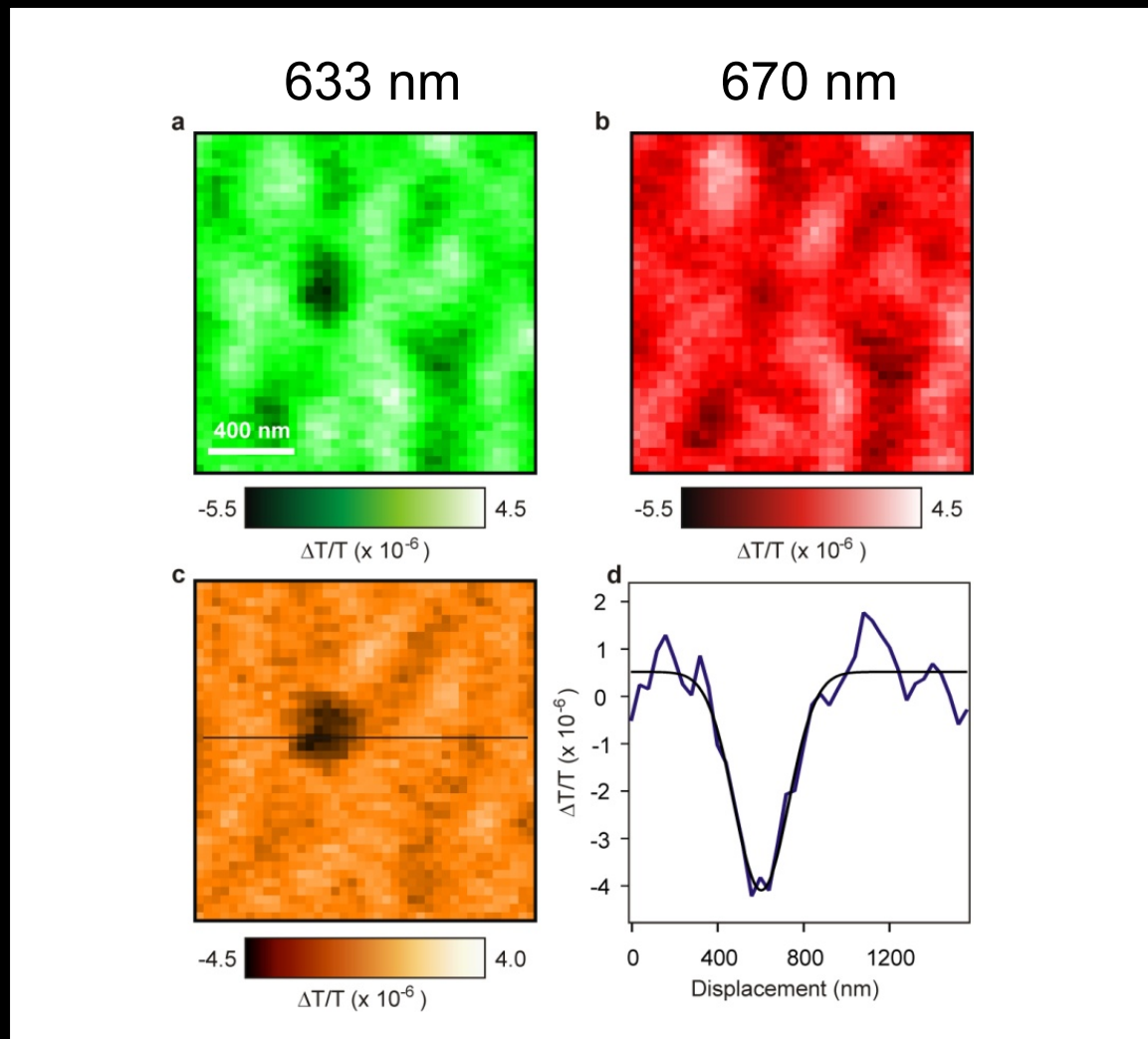
-6  6
Extinction ($\times 10^4$)

Single-molecule absorption *imaging* and *spectroscopy*



M. Celebrano, P. Kukura, A. Renn, V. Sandoghdar
Nature Photonics 5, 95 (2011).

Imaging a single molecule in absorption



M. Celebrano, P. Kukura, A. Renn, V. Sandoghdar
Nature Photonics 5, 95 (2011).

Extinction spectroscopy/microscopy lets us see strongly quenched systems

NANO LETTERS

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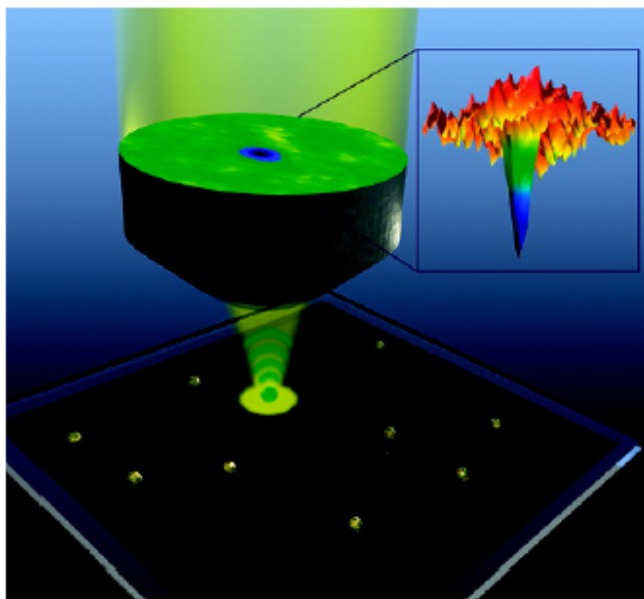
Letter

Imaging a Single Quantum Dot When It Is Dark

P. Kukura, M. Celebrano, A. Renn, and V. Sandoghdar

Nano Lett., 2009, 9 (3), 926-929 • DOI: 10.1021/nl801735y • Publication Date (Web): 01 August 2008

Downloaded from <http://pubs.acs.org> on March 16, 2009



Fluorescence

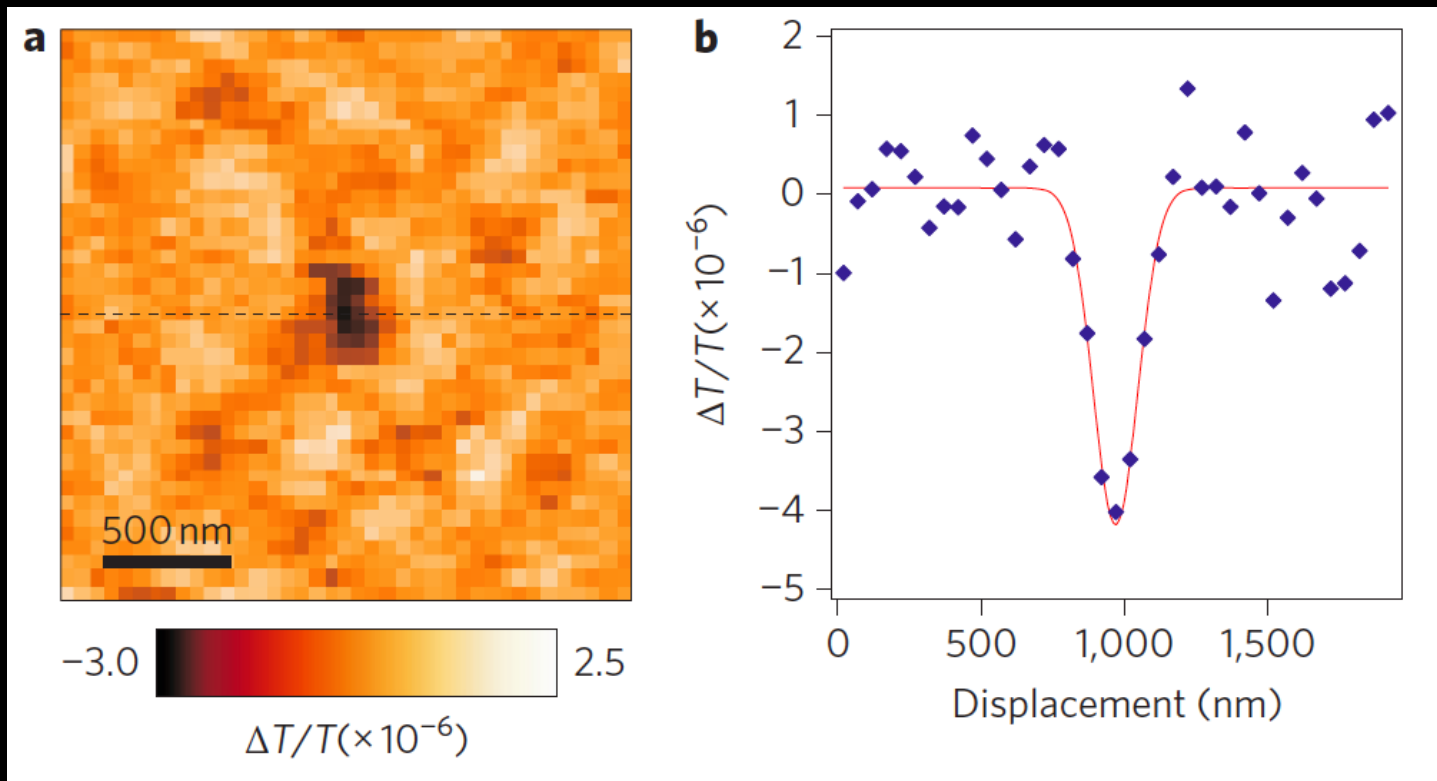
$$\eta \propto \frac{\gamma_r}{\gamma_r + \gamma_{nr}}$$

Extinction

$$\sigma \propto \frac{\gamma_r}{\gamma_r + \gamma_{nr} + \gamma_{deph}}$$

$$\gamma_{deph} \approx 10^{4-5} \gamma_r$$

Imaging a quenched molecule



During the same summer !

Room-Temperature Detection of a Single Molecule's Absorption by Photothermal Contrast

A. Gaiduk, M. Yorulmaz, P. V. Ruijgrok, M. Orrit*

SCIENCE VOL 330 15 OCTOBER 2010

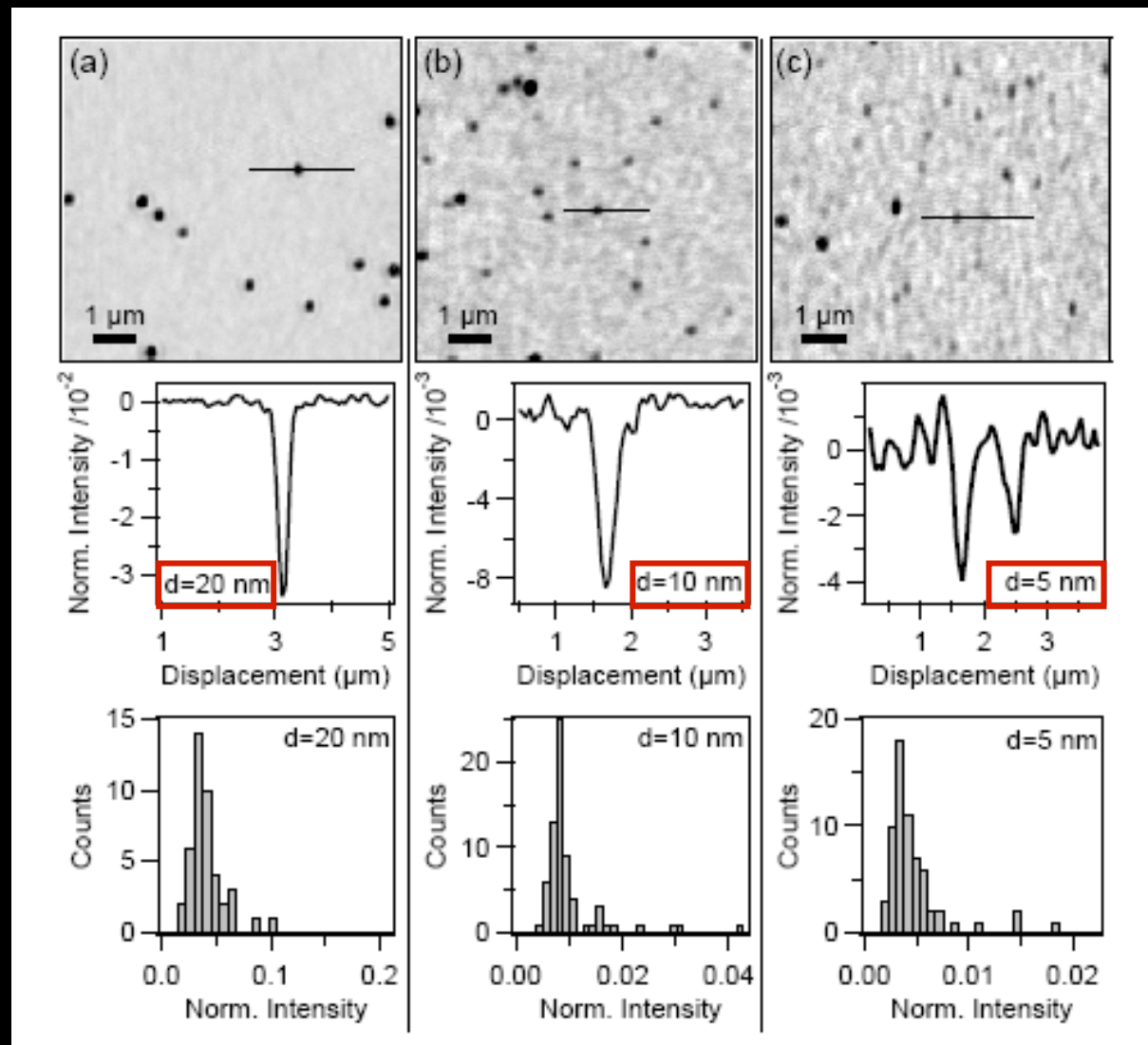
Ground-State Depletion Microscopy: Detection Sensitivity of Single-Molecule Optical Absorption at Room Temperature

Shasha Chong,[†] Wei Min,^{†,‡} and X. Sunney Xie*

J. Phys. Chem. Lett. 2010, 1, 3316–3322

Scattering without absorption

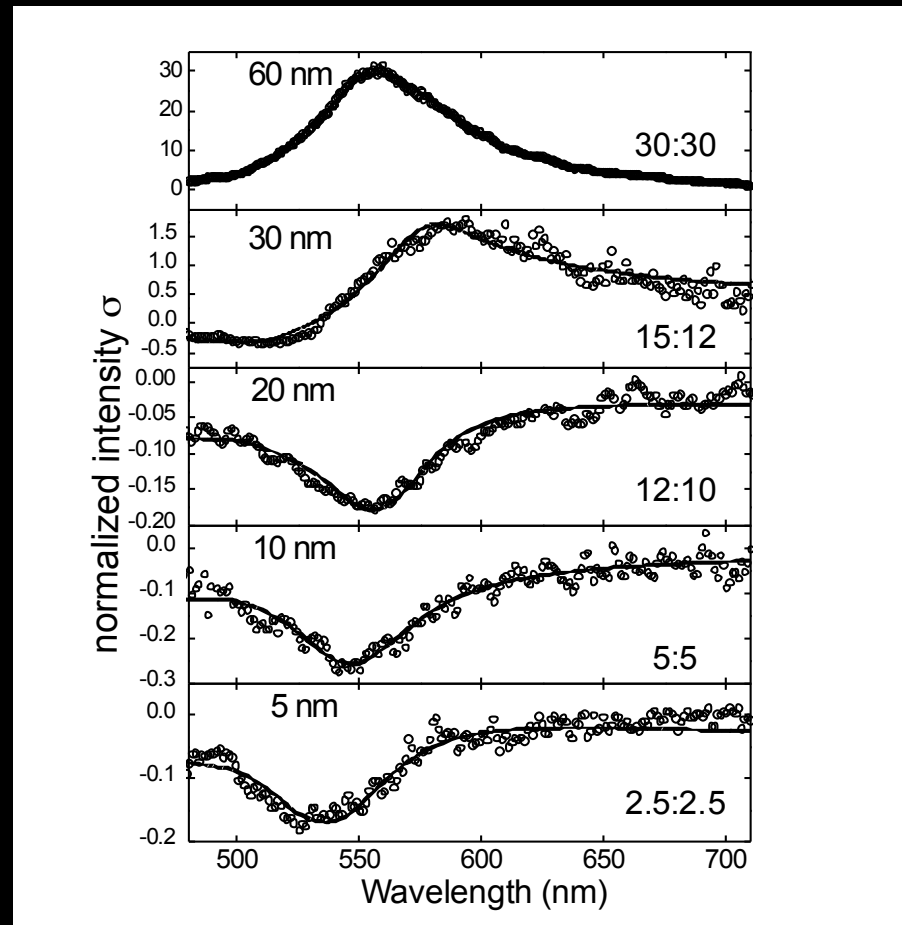
Detection of gold nanoparticles as small as 5 nm



V. Jacobsen, P. Stoller, C. Brunner, V. Vogel, V. Sandoghdar,
Opt. Exp. **14**, 405 (2006).

Plasmon Spectra of Single Gold Nanoparticles

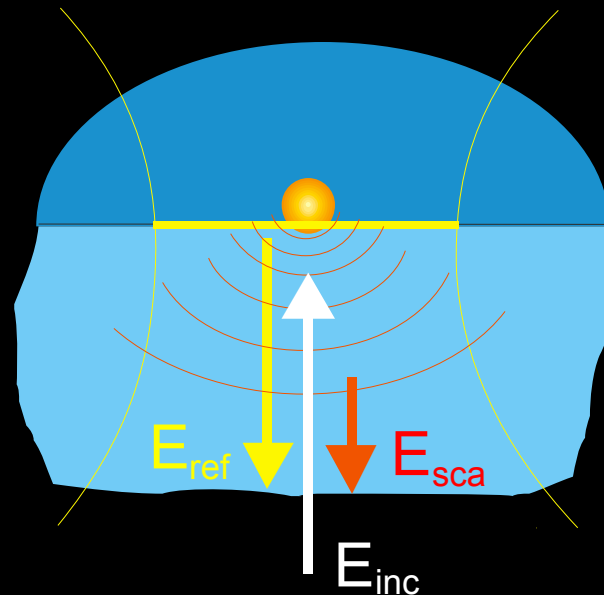
Excitation by quasi-white light from photonic crystal fibers



K. Lindfors, T. Kalkbrenner, P. Stoller, & V. Sandoghdar, *PRL* **93**, 037401 (2004).

P. Stoller, V. Jacobsen, V. Sandoghdar, *Opt. Lett.* **31**, 2474 (2006).

Interferometric detection of *nonfluorescent* nanoparticles



$$E_{\text{ref}} \sim r e^{-i\pi/2} E_{\text{inc}}$$

polarizability

$$E_{\text{sca}} = s E_{\text{inc}} \sim D^3 \frac{\varepsilon - \varepsilon_m}{\varepsilon + 2\varepsilon_m} \frac{e^{ikR}}{R} E_{\text{inc}}$$

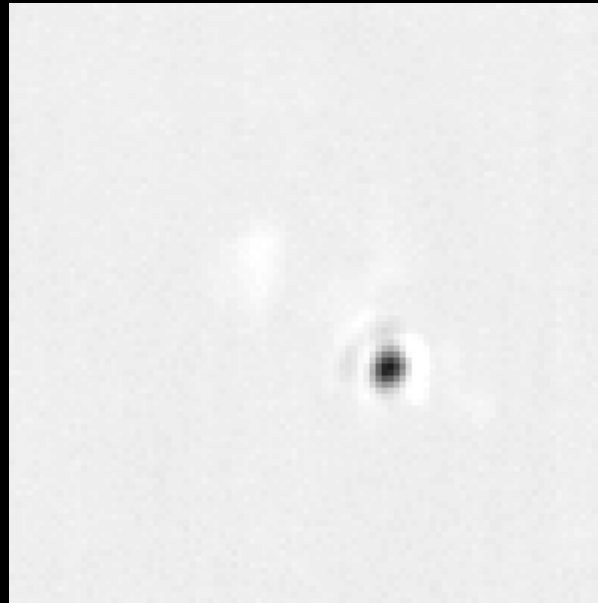
φ : scattering phase: $s = |s| e^{i\varphi}$

scales as D^3

$$I_{\text{det}} = |E_{\text{ref}} + E_{\text{sca}}|^2 = |E_{\text{inc}}|^2 \left\{ r^2 + |s|^2 - 2r|s|\sin\varphi \right\}$$

K. Lindfors, T. Kalkbrenner, P. Stoller, & V. Sandoghdar, *PRL* **93**, 037401 (2004).

Free diffusion of gold nanoparticles on supported membranes



40nm AuNPs diffusing on a DOPC supported membrane (0.8% GM1)

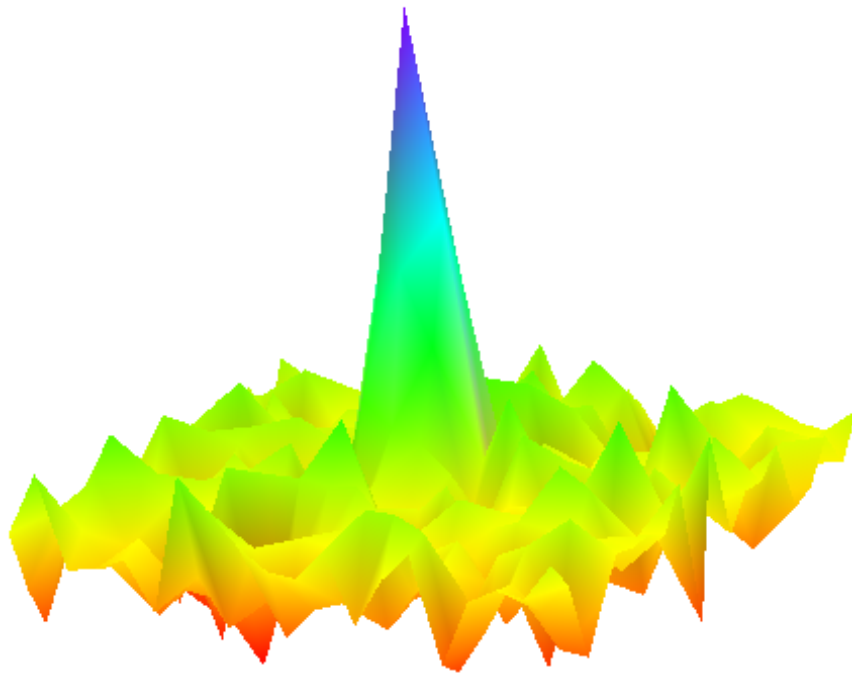
Labeling scheme: AuNP – CTxB – GM1 lipid molecule

Frame rate: 200 Hz (video is 10x slow)

Field of view: $6.3 \times 6.3 \mu\text{m}^2$

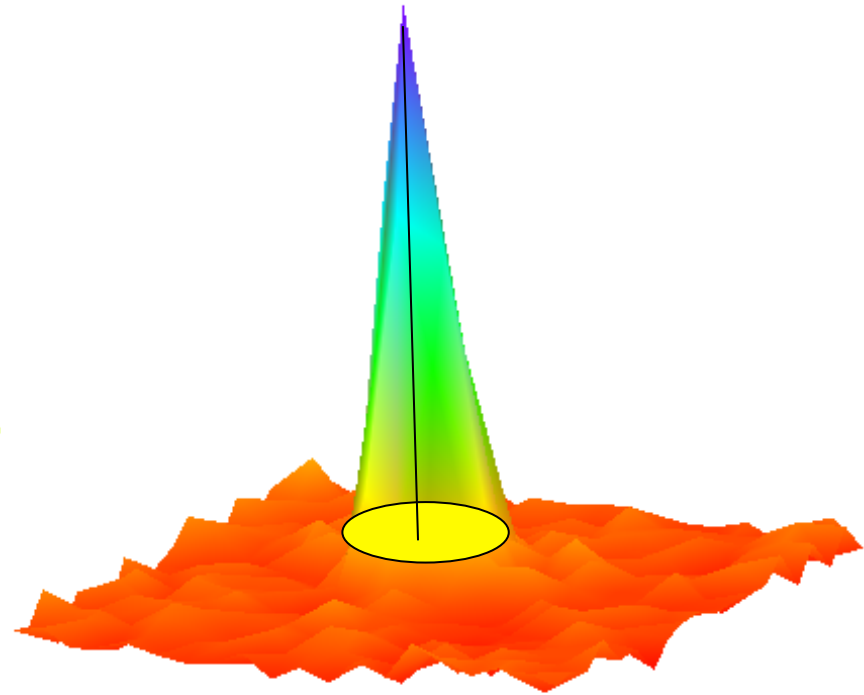
Accuracy in localizing a single nano-object

SNR ~ 10

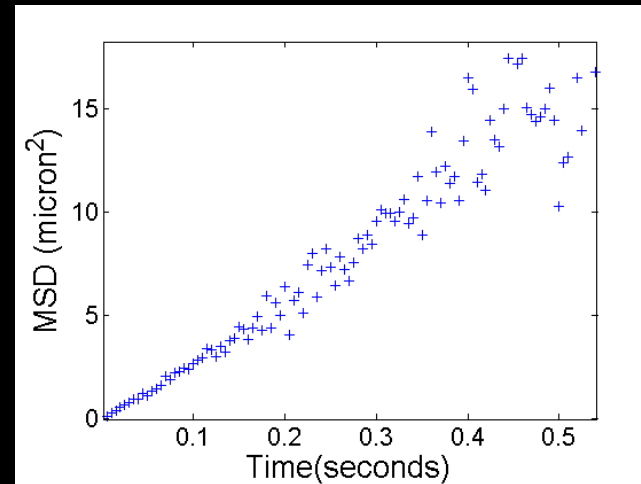
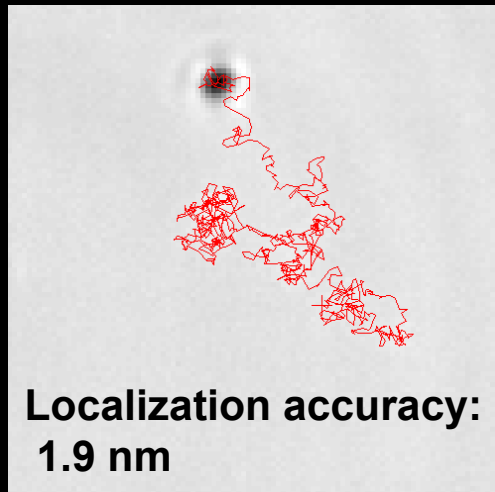


Localization accuracy: 10 nm

SNR ~ 50

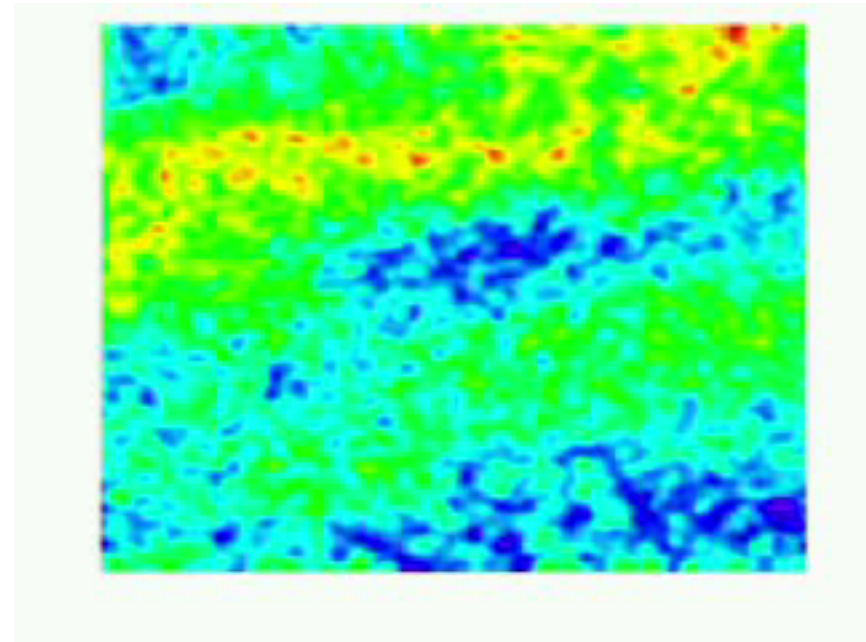
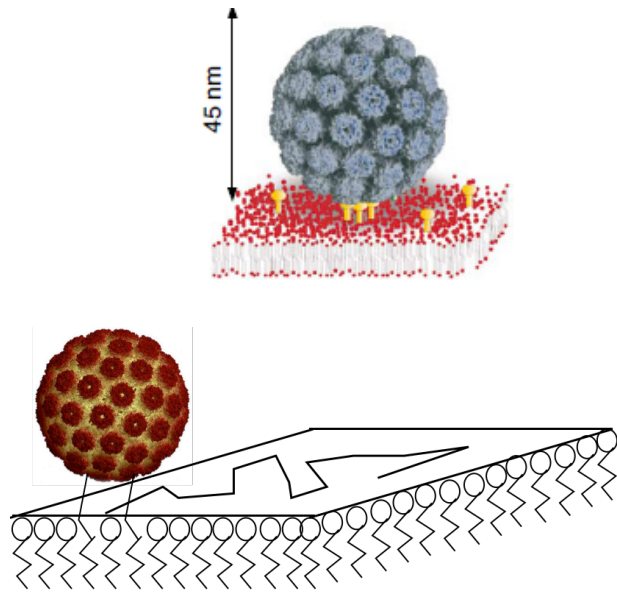


Localization accuracy: 2 nm



Diffusion coefficient: $8.32 \mu\text{m}^2/\text{sec}$

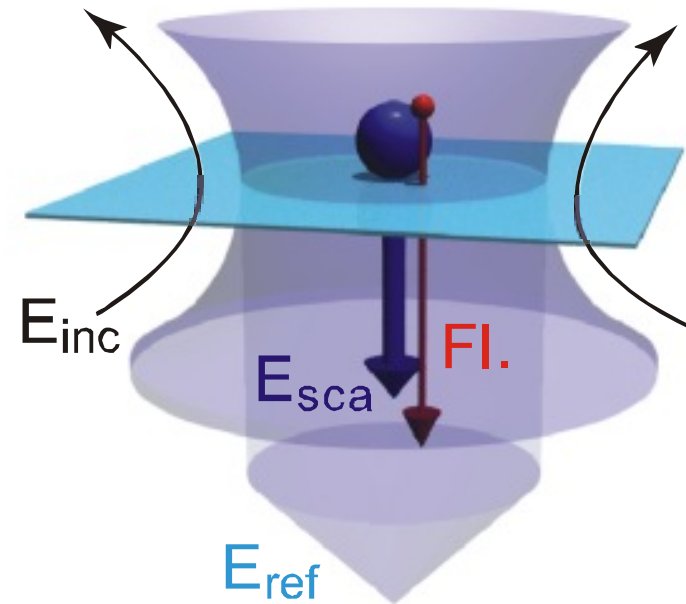
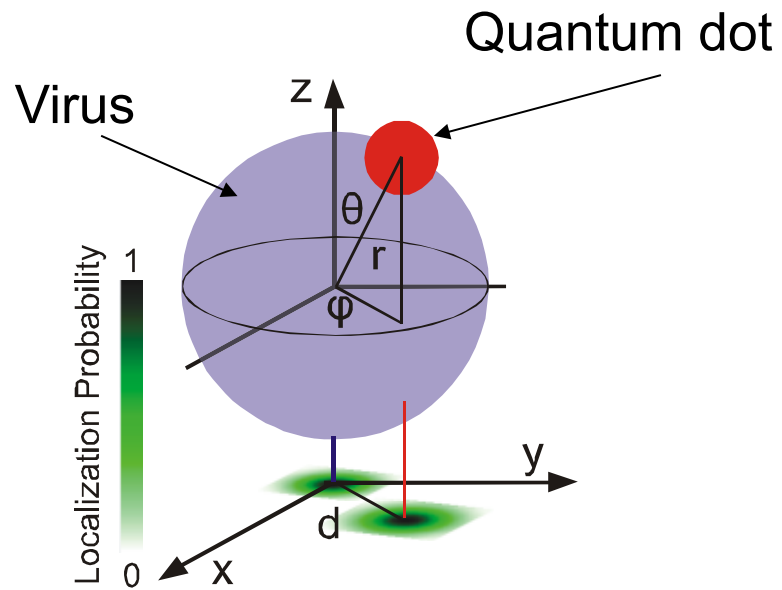
Interferometric Tracking of a single *naked* virion on a supported bilayer membrane containing GM1 receptors



H. Ewers, V. Jacobsen, E. Klotzsch, A. Smith, A. Helenius, V. Sandoghdar,
Nano Lett. **7**, 2263 (2007).

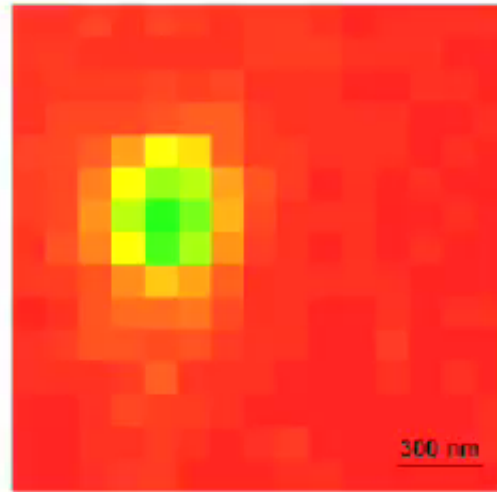
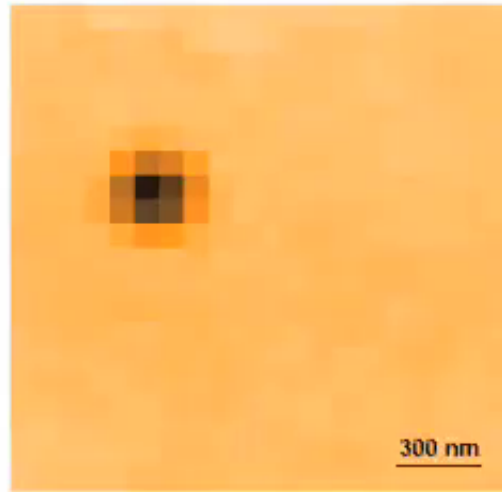
P. Kukura, H. Ewers, C. Müller, A. Renn, A. Helenius, V. Sandoghdar
Nature Methods, **6**, 923 (2009);

How does a virus move on the membrane: rolls or slides?
Translational & orientational nano-motion of a virus

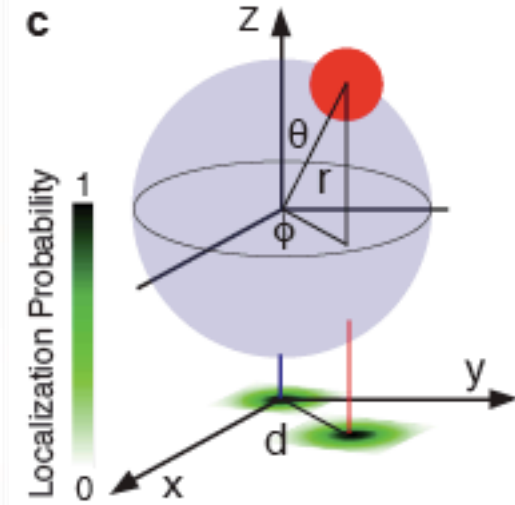


Colocalization of scattering and fluorescence

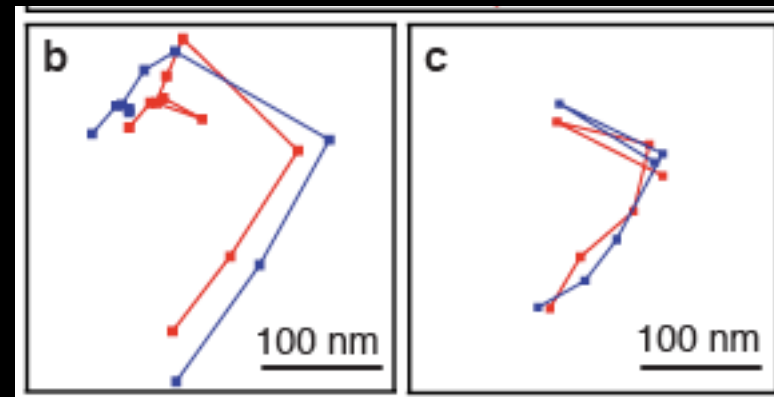
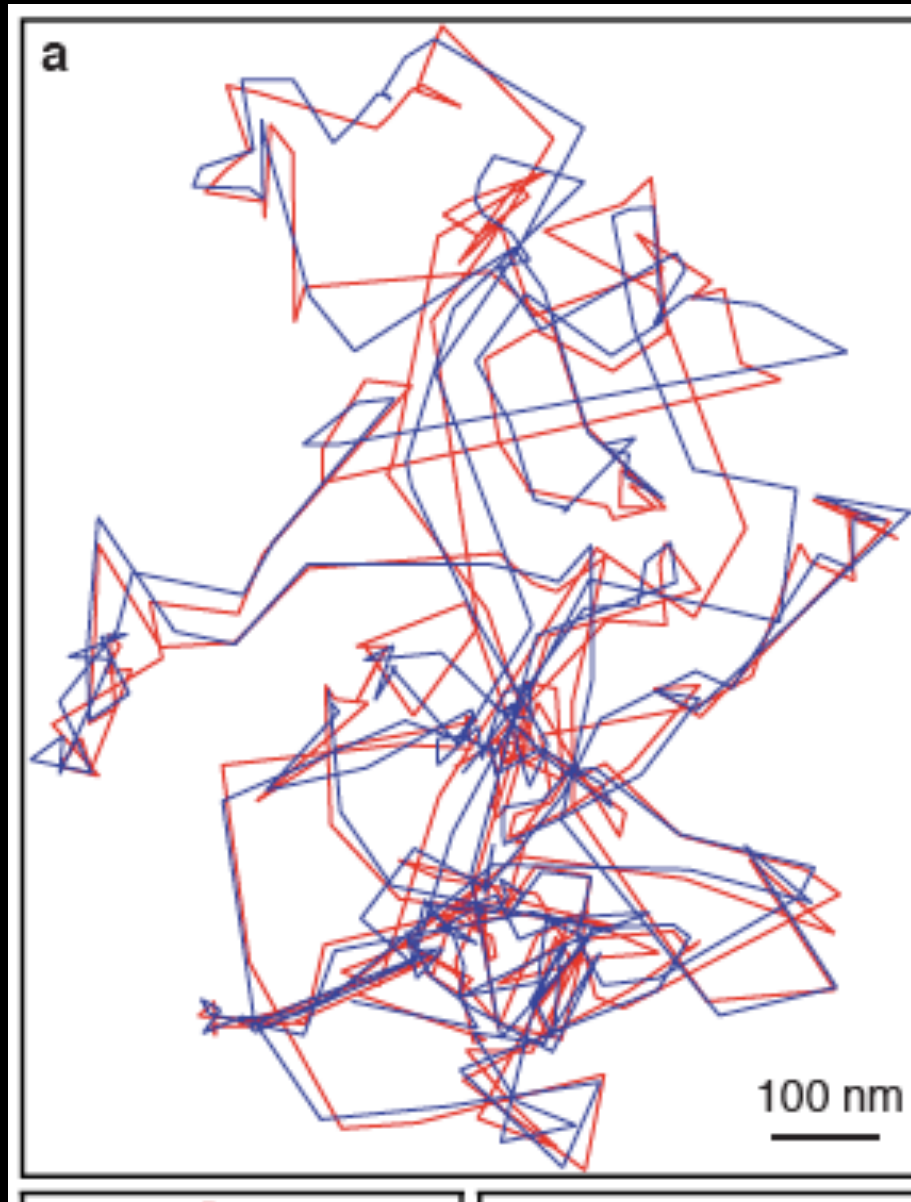
a
Scattering (%)
1
0



c



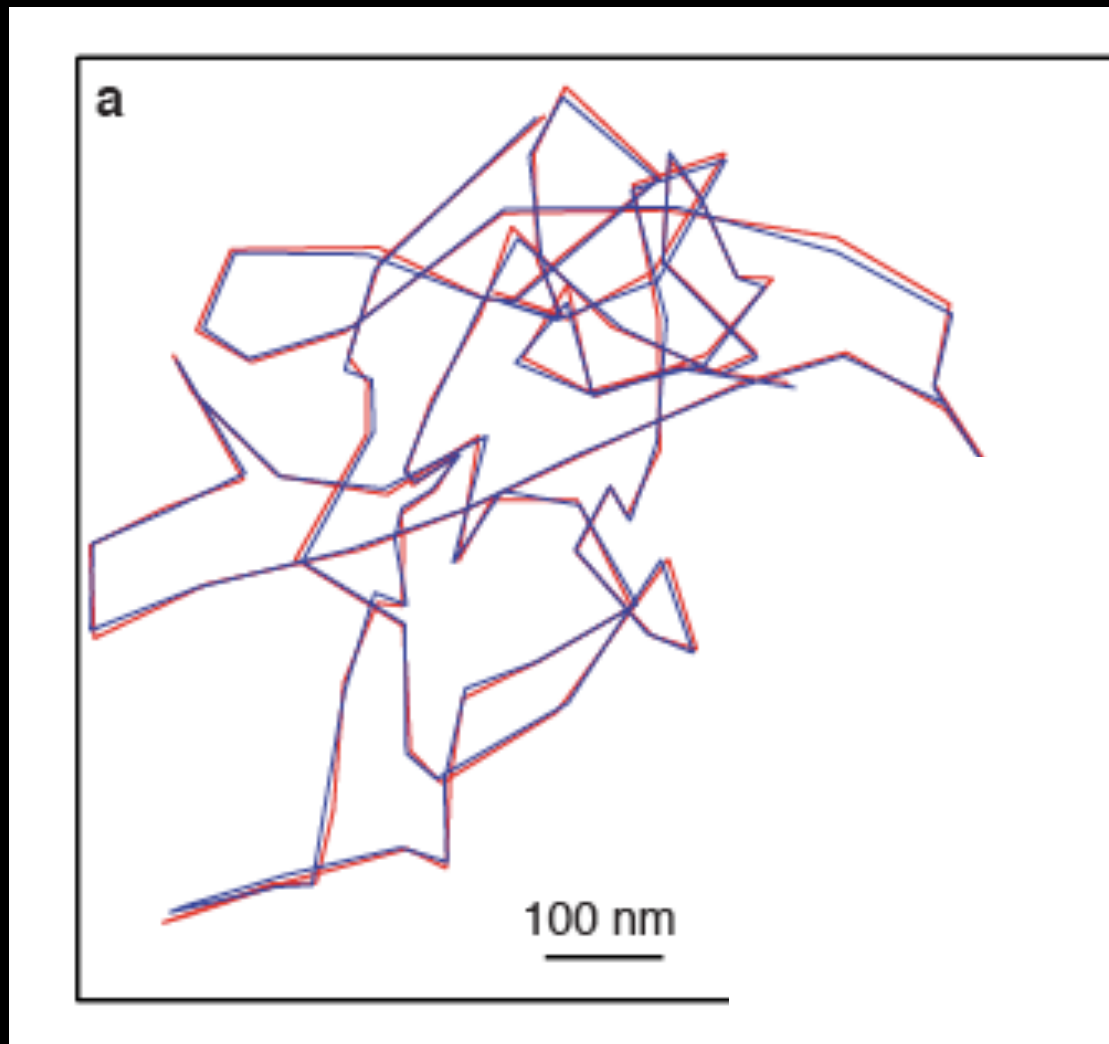
Simultaneous trajectories of the virus center of mass
and of the quantum dot: *0.05% receptor concentration*



The virus dance

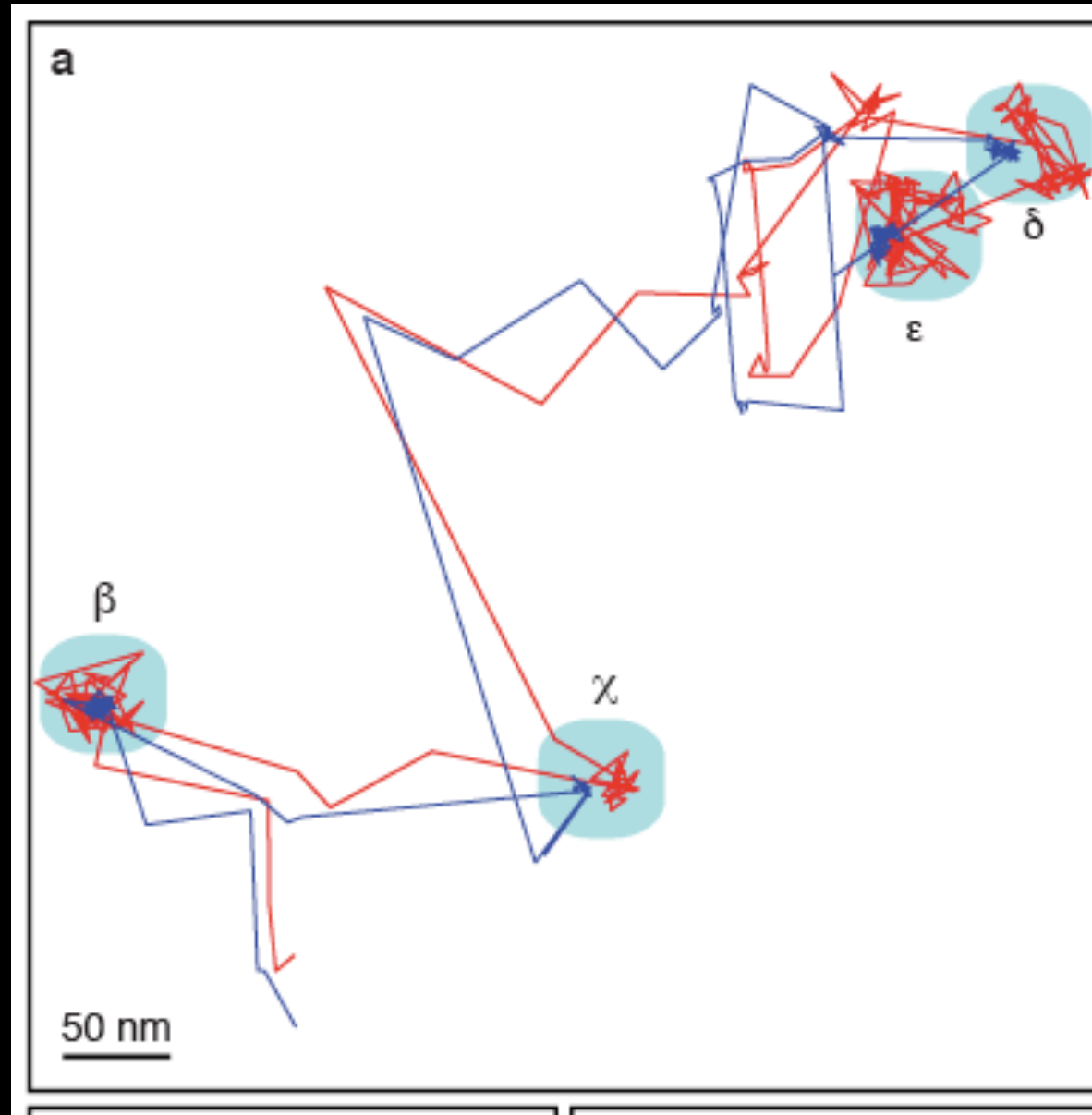
Checking

Fluo and scattering trajectories of a homogeneously doped bead

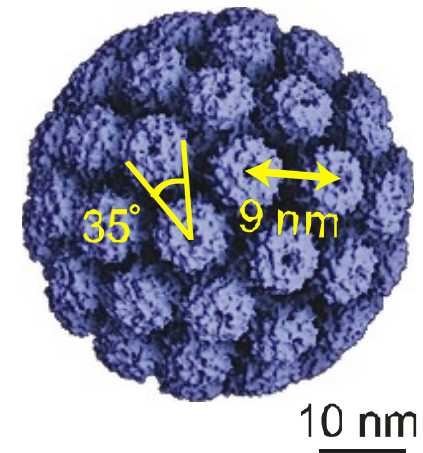
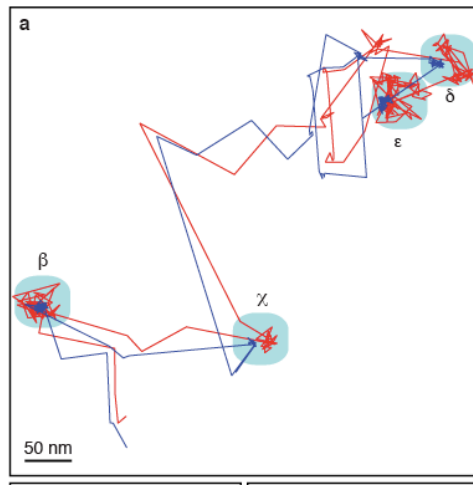
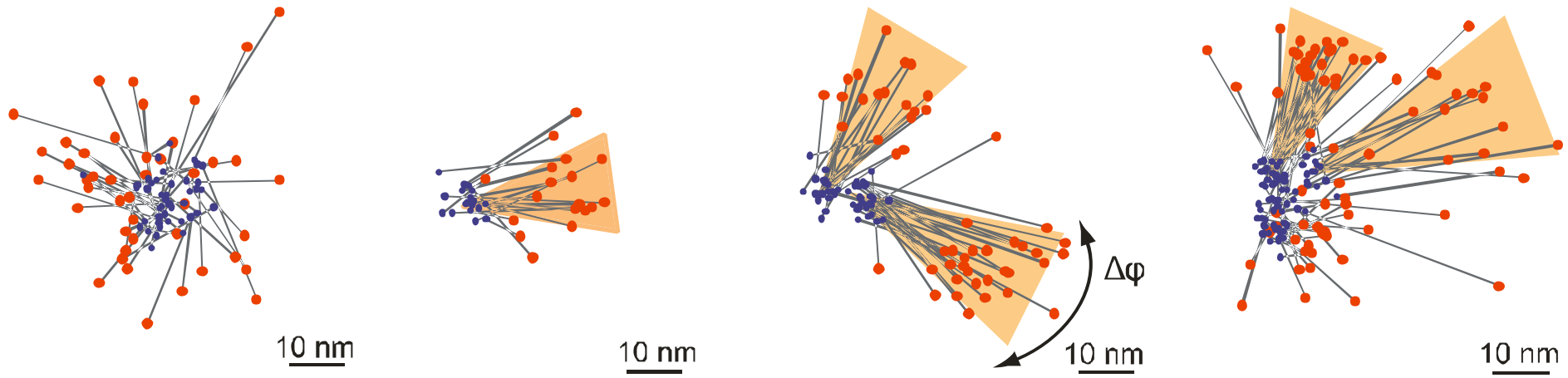


Fluo and scattering „trajectories“ of an immobilized Qdots-labeled virus

Simultaneous trajectories of the virus center of mass and of the quantum dot: 1% receptor concentration

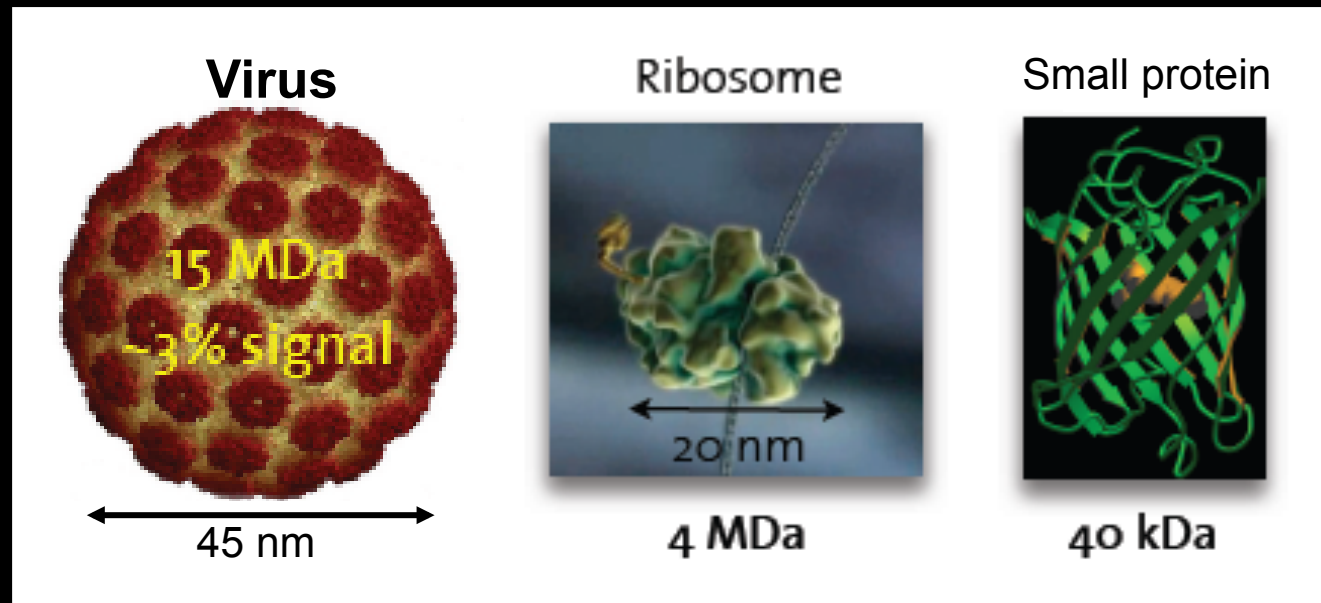


Resolving the nanomotion of the virus *Sliding, tumbling, rocking of a single virus*



P. Kukura, H. Ewers, C. Müller, A. Renn, A. Helenius, V. Sandoghdar
Nature Methods, 6, 923 (2009);

Direct detection of smaller biological nano-objects:
no absorption, no fluorescence



We are close to the detection of a single protein.

Goal: *Detect and image single nonfluorescent atoms on a surface*

Conclusions

Direct detection of single nanoparticles and molecules using LIGHT and *without fluorescence* is now possible

Focused light is one of the most sensitive sensors of matter